

The spectral action for Dirac operators with skew-symmetric torsion

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joint work with

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- 1 Connections with torsion
- 2 Spectral action for pure gravity with torsion
- 3 Lagrangian for SM in presence of torsion
- 4 Questions

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set-up

Consider 4-dim. Riemannian manifolds M , closed and spin.

General connection for vector fields has the form

$$\nabla_X Y = \nabla_X^{LC} Y + A(X, Y)$$

where ∇^{LC} is Levi-Civita connection
and A is a $(2, 1)$ -tensor field.

Torsion 3-form of ∇ is

$$T(X, Y, Z) = \langle \nabla_X Y - \nabla_Y X - [X, Y], Z \rangle$$



∇ is compatible with the Riemannian metric $\langle \cdot, \cdot \rangle$ and has same geodesics as ∇^{LC}

if and only if $A(X, Y, Z) = \langle A(X, Y), Z \rangle$ is totally anti-symmetric. (“Torsion is skew-symmetric.”)

In this case: $T = 2A$, and hence

$$\nabla_X Y = \nabla_X^{LC} Y + \frac{1}{2} T(X, Y, \cdot)^\#$$

Induced connection for spinor fields is then

$$\nabla_X \psi = \nabla_X^{LC} \psi + \frac{1}{4} (X \lrcorner T) \cdot \psi,$$

where $(X \lrcorner T) \cdot$ is Clifford multiplication by the 2-form $T(X, \cdot, \cdot)$.



Dirac operator D is given by $D\psi = \sum_i e_i \cdot \nabla_{e_i} \psi$, for any orthonormal frame e_i .

Bochner formula (Agricola-Friedrich)

$$D^2 = \Delta + \frac{3}{4}dT + \frac{1}{4}R - \frac{9}{8}T_0^2,$$

where $\Delta = \tilde{\nabla}^* \tilde{\nabla}$ is the Laplacian associated to spin connection

$$\tilde{\nabla}_X \psi = \nabla_X^{LC} \psi + \frac{3}{4}(X \lrcorner T) \cdot \psi,$$

dT is the exterior differential of the 3-form T ,

R is the scalar curvature and

$$T_0^2 = \frac{1}{6} \sum_{i,j=1}^n \|T(e_i, e_j, \cdot)\# \|^2.$$

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Chamseddine-Connes:

Spectral action (bosonic part) of Dirac operator D is the number of eigenvalues of D in the interval $[-\Lambda, \Lambda]$ (with $\Lambda \in \mathbb{R}^+$):

$$I = \text{Tr} f \left(\frac{D^2}{\Lambda^2} \right),$$

where Tr is the L^2 -trace over the space of spinor fields, and f is cut-off function with support in $[-1, +1]$ which is constant near 0.

From the heat trace asymptotics for $t \rightarrow 0$

$$\mathrm{Tr} \left(e^{-tD^2} \right) \sim \sum_{n \geq 0} t^{n-2} a_{2n}(D^2)$$

(with Seeley-deWitt coefficients $a_{2n}(D^2)$)

one gets (by $t = \Lambda^{-2}$) an asymptotics for the spectral action

$$I = \mathrm{Tr} f \left(\frac{D^2}{\Lambda^2} \right) \sim \Lambda^4 f_4 a_0(D^2) + \Lambda^2 f_2 a_2(D^2) + \Lambda^0 f_0 a_4(D^2)$$

as $\Lambda \rightarrow \infty$.

Here f_4, f_2, f_0 are moments of the cut-off function f .

Bochner formula is $D^2 = \Delta - E$ with $E = -\frac{3}{4}dT - \frac{1}{4}R + \frac{9}{8}T_0^2$.

Insert this into formulas for Seeley-deWitt coefficients:

$$a_0(D^2) = \frac{1}{4\pi^2} \int_M d\text{vol}$$

$$a_2(D^2) = \frac{1}{96\pi^2} \int_M (6 \text{tr}(E) + 4R) d\text{vol}$$

$$a_4(D^2) = \frac{1}{5760\pi^2} \int_M \left(\text{tr} \left(60 \Delta E + 60RE + 180E^2 + 30 \Omega_{ij} \Omega_{ij} \right) \right. \\ \left. + 48\Delta^{LC}R + 20R^2 - 8\|Ric\|^2 + 8\|Riem\|^2 \right) d\text{vol}$$

where Ric and $Riem$ are Ricci/Riemannian curvature tensors of the metric,

and $\Omega_{ij} = \tilde{\nabla}_{e_i} \tilde{\nabla}_{e_j} - \tilde{\nabla}_{e_j} \tilde{\nabla}_{e_i} - \tilde{\nabla}_{[e_i, e_j]}$ is the curvature of $\tilde{\nabla}$.

Clifford relations and cyclicity of trace $\implies \text{tr}(dT) = 0$

Therefore

$$a_2(D^2) = \frac{1}{16\pi^2} \int_M \left(\frac{9}{2} T_0^2 - \frac{1}{3} R \right) d\text{vol}.$$

Neglecting the a_4 -term in the spectral action, one obtains the classical Einstein-Cartan-action: Only torsion free critical points upon variation of metric and torsion 3-forms.

We compute for the curvature Ω_{ij} of the connection $\tilde{\nabla}$:

$$\begin{aligned} \Omega_{ij} = & \sum_{a,b} \left(\frac{1}{4} \langle R(\mathbf{e}_i, \mathbf{e}_j) \mathbf{e}_a, \mathbf{e}_b \rangle + \frac{3}{8} a(\nabla T)(\mathbf{e}_i, \mathbf{e}_j, \mathbf{e}_a, \mathbf{e}_b) \right. \\ & \left. + \frac{9}{16} \sum_c (T(\mathbf{e}_i, \mathbf{e}_c, \mathbf{e}_a) T(\mathbf{e}_j, \mathbf{e}_c, \mathbf{e}_b) - T(\mathbf{e}_j, \mathbf{e}_c, \mathbf{e}_a) T(\mathbf{e}_i, \mathbf{e}_c, \mathbf{e}_b)) \right) \mathbf{e}_a \mathbf{e}_b, \end{aligned}$$

where $a(\nabla T)(\mathbf{e}_i, \mathbf{e}_j, \mathbf{e}_a, \mathbf{e}_b) = \nabla_{\mathbf{e}_i} T(\mathbf{e}_j, \mathbf{e}_a, \mathbf{e}_b) - \nabla_{\mathbf{e}_j} T(\mathbf{e}_i, \mathbf{e}_a, \mathbf{e}_b)$.

Then using $\text{tr}(\mathbf{e}_k \mathbf{e}_l \mathbf{e}_s \mathbf{e}_t) = 4(\delta_{ls} \delta_{kt} - \delta_{lt} \delta_{ks})$:

$$\begin{aligned} \sum_{i,j} \text{tr}(\Omega_{ij} \Omega_{ij}) = & -8 \sum_{\substack{i \neq j \\ a,b}} \left(\frac{1}{4} \langle R(\mathbf{e}_i, \mathbf{e}_j) \mathbf{e}_a, \mathbf{e}_b \rangle \right. \\ & \left. + \frac{3}{8} a(\nabla T)(\mathbf{e}_i, \mathbf{e}_j, \mathbf{e}_a, \mathbf{e}_b) + \frac{9}{8} c(T)(\mathbf{e}_i, \mathbf{e}_j, \mathbf{e}_a, \mathbf{e}_b) \right)^2, \end{aligned}$$

where $c(T)(\mathbf{e}_i, \mathbf{e}_j, \mathbf{e}_a, \mathbf{e}_b) = \sum_c T(\mathbf{e}_i, \mathbf{e}_c, \mathbf{e}_a) T(\mathbf{e}_j, \mathbf{e}_c, \mathbf{e}_b)$.



Representation theory of $O(4) \rightsquigarrow$ evaluation of $\sum_{i,j} \text{tr}(\Omega_{ij}\Omega_{ij})$.

The fourth Seeley-deWitt coefficient is

$$a_4(D^2) = \frac{1}{16\pi^2} \int_M \left(\frac{1}{72} R^2 - \frac{1}{45} \|Ric\|^2 - \frac{7}{360} \|Riem\|^2 \right. \\ \left. + \frac{81}{32} (T_0^2)^2 - \frac{27}{32} \|c(T)\|^2 \right. \\ \left. + \frac{9}{16} \|dT\|^2 - \frac{3}{8} \|d^*T\|^2 - \frac{3}{8} \|sym_0^2(\nabla T)\|^2 \right. \\ \left. - \frac{3}{8} R(T) \right) dvol$$

with

$$R(T) = - \sum_{\substack{i \neq j \\ a, c}} Ric(e_j, e_a) T(e_i, e_c, e_a) T(e_j, e_c, e_i) \\ + \sum_{\substack{i \neq j \\ a, b, c}} \langle W(e_i, e_j) e_a, e_b \rangle T(e_i, e_c, e_a) T(e_j, e_c, e_b)$$

Observations

- $R(T)$ couples torsion and the trace free component of the curvature tensor. \rightsquigarrow Expect critical points of spectral action with non-zero torsion.
- Existence of the derivative terms of T . \rightsquigarrow Torsion gets dynamical

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Specific particle model

Ingredients of the particle model:

- internal Hilbert space: \mathcal{H}_f with connection $\nabla^{\mathcal{H}_f}$
- full Hilbert space: $\mathcal{H}_{SM} = L^2(M, \mathcal{S}) \otimes \mathcal{H}_f \ni \psi \otimes \chi$
- twisted connection (with torsion):

$$\widehat{\nabla}^{SM} = \nabla \otimes \text{id}_{\mathcal{H}_f} + \text{id}_{\mathcal{S}} \otimes \nabla^{\mathcal{H}_f}$$
- associated Dirac operator: $D^{\widehat{\nabla}^{SM}}$
- field Φ of endomorphisms of \mathcal{H}_f
 (encodes Higgs boson, Yukawa couplings, etc.)

The Chamseddine-Connes Dirac operator:

$$D_{\Phi}(\psi \otimes \chi) = D^{\widehat{\nabla}^{SM}}(\psi \otimes \chi) + \gamma_5 \psi \otimes \Phi \chi$$

Bochner formula for $D_\Phi(\psi \otimes \chi) = D^{\widehat{\nabla}^{SM}}(\psi \otimes \chi) + \gamma_5 \psi \otimes \Phi \chi$:

$$(D_\Phi)^2(\psi \otimes \chi) = \Delta^{\overline{\nabla}}(\psi \otimes \chi) - E_\Phi(\psi \otimes \chi)$$

with

$$\begin{aligned} E_\Phi(\psi \otimes \chi) &= \left(\left(-\frac{3}{4} dT - \frac{1}{4} R + \frac{9}{8} T_0^2 \right) \psi \right) \otimes \chi \\ &+ \frac{1}{2} \cdot \sum_{i \neq j} (\mathbf{e}_i \cdot \mathbf{e}_j \cdot \psi) \otimes \left(\Omega_{ij}^{\mathcal{H}} \chi \right) \\ &+ \sum_{i=1}^n \gamma_5 \mathbf{e}_i \cdot \psi \otimes [\nabla_{\mathbf{e}_i}^{\mathcal{H}_f}, \Phi] \chi - \psi \otimes (\Phi^2) \chi \end{aligned}$$

Bosonic Lagrangian of the SM coupled to gravity and torsion:

$$\begin{aligned}
 I = & \frac{24\Lambda^4 f_4}{\pi^2} \int_M d\text{vol} \\
 & + \frac{\Lambda^2 f_2}{\pi^2} \int_M \left\{ 27 T_0^2 - 2R - a|\varphi|^2 - \frac{1}{2}c \right\} d\text{vol} \\
 & + \frac{f_0}{2\pi^2} \int_M \left\{ \frac{1}{6}R^2 - \frac{4}{15}\|Ric\|^2 - \frac{7}{30}\|Riem\|^2 + \frac{243}{8}(T_0^2)^2 - \frac{27}{4}\|c(T)\|^2 \right. \\
 & \quad + \frac{27}{4}\|dT\|^2 - \frac{9}{2}\|d^*T\|^2 - \frac{9}{2}\|sym_0^2(\nabla T)\|^2 - \frac{9}{2}R(T) \\
 & \quad + g_3^2\|G\|^2 + g_2^2\|F\|^2 + \frac{5}{3}g_1^2\|B\|^2 \\
 & \quad + a|D_\nu\varphi|^2 + b|\varphi|^4 + 2e|\varphi|^2 + \frac{1}{2}d + \frac{1}{6}R \left(a|\varphi|^2 + \frac{1}{2}c \right) \\
 & \quad \left. - \frac{9}{4}T_0^2 \left(a|\varphi|^2 + \frac{1}{2}c \right) \right\} d\text{vol}.
 \end{aligned}$$

Observations

- Torsion couples to the Higgs field φ .
- $R(T)$ couples torsion and the trace free component of the curvature tensor. \rightsquigarrow Expect critical points of spectral action with non-zero torsion.
- Existence of the derivative terms of T . \rightsquigarrow Torsion gets dynamical

The full Standard Model action is given by

$$I_{SM} = \text{Tr} f \left(\frac{D_\Phi^2}{\Lambda^2} \right) + \langle \Psi, D_\Phi \Psi \rangle \quad \text{with} \quad \Psi \in \mathcal{H}_{SM}$$

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Some Questions

Experimental signatures:

- Torsion field of the Earth (Weyl tensor $\neq 0$)?
- Effects on freely falling particles with different spin?

The cosmological constant:

- Spectral action “predicts” huge cosm. constant:
 $\Lambda \sim 10^{17} \text{GeV}$
- Torsion might induce Fermi condensate (Perez, Rovelli)
- Dynamical cancelation of cosm. constant + inflation?

Further implications:

- Connections to LQG (Barbero-Immirzi parameter)?
- Torsion in QFT?