

Indefinite Theta Series and Applications

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Based on

C.N. arXiv:1609.01224 and K. Bringmann, C.N. arXiv:1803.09270

Theta Functions for Positive Definite Lattices

Let Λ be a positive-definite, even lattice. Theta functions are defined as:

$$\Theta_{\mu}^Q(\tau, \mathfrak{z}) := \sum_{n \in \Lambda + \mu} q^{\frac{1}{2}n^2} e^{2\pi i \mathfrak{z} \cdot n}, \quad \mu \in \Lambda^* / \Lambda$$

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- Elliptic transformations:

$$\Theta_{\mu}^Q(\tau, \mathfrak{z} + m) = \Theta_{\mu}^Q(\tau, \mathfrak{z}) \quad \text{for } m \in \Lambda,$$

$$\Theta_{\mu}^Q(\tau, \mathfrak{z} + m\tau) = q^{-\frac{1}{2}m^2} e^{-2\pi i \mathfrak{z} \cdot m} \Theta_{\mu}^Q(\tau, \mathfrak{z}) \quad \text{for } m \in \Lambda.$$

- Modular Transformations:

$$\Theta_{\mu}^Q(\tau + 1, \mathfrak{z}) = e^{\pi i \mu^2} \Theta_{\mu}^Q(\tau, \mathfrak{z}),$$

$$\Theta_{\mu}^Q\left(-\frac{1}{\tau}, \frac{\mathfrak{z}}{\tau}\right) = \frac{(-i\tau)^{N/2}}{\sqrt{|\Lambda^*/\Lambda|}} e^{\pi i \mathfrak{z}^2 / \tau} \sum_{\nu \in \Lambda^* / \Lambda} e^{-2\pi i \mu \cdot \nu} \Theta_{\nu}^Q(\tau, \mathfrak{z}).$$

Theta Functions and Indefinite Signature Lattices?

Theta functions we defined do not converge for indefinite signature lattices, signature (N_+, N_-) , since $q^{n^2/2}$ factors grow unbounded along negative-definite directions.

- However, we can form real analytic modular forms for indefinite signature lattices.

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- Preserving holomorphy?

Example: $\sum_{n \in (\Lambda + \mu) \cap \mathcal{C}} q^{n^2/2} e^{2\pi i n \cdot \mathfrak{z}}$ where \mathcal{C} is a non-negative cone.

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An even more specific example:

$$\sum_{n \in \Lambda + \mu} \left[\frac{1}{2^{N_-}} \prod_{j=1}^{N_-} \left(\operatorname{sgn}(c_j \cdot n) - \operatorname{sgn}(c'_j \cdot n) \right) \right] q^{\frac{1}{2}n^2}.$$

Lorentzian Lattices and Mock Modularity

For signature $(n - 1, 1)$ (Lorentzian) lattices,

$$\sum_{n \in \Lambda + \mu} \frac{1}{2} [\operatorname{sgn}(c \cdot n) - \operatorname{sgn}(c' \cdot n)] q^{n^2/2},$$

the modular completion for signature $(n - 1, 1)$ (Lorentzian) lattices is given by [Zwegers 2002]:

$$\operatorname{sgn}(c \cdot n) - \operatorname{sgn}(c' \cdot n) \rightarrow \operatorname{erf}(c \cdot n \sqrt{2\pi\tau_2}) - \operatorname{erf}(c' \cdot n \sqrt{2\pi\tau_2}),$$

where $c^2 = c'^2 = -1$ and $c \cdot c' < 0$.

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- In physics [Akhoury, Comtet]: $H = \frac{p^2}{2} + \frac{1}{2} (W'(x))^2 + \frac{1}{2} W''(x) \sigma_3$.
With $\phi_{\pm} = W'(\pm\infty)$, the supersymmetric partition function is:

$$\frac{1}{2} \left(\frac{\phi_+}{|\phi_+|} - \frac{\phi_-}{|\phi_-|} \right) - \frac{1}{2} \left[\operatorname{erfc} \left(\phi_+ \sqrt{\beta/2} \right) - (+ \leftrightarrow -) \right].$$

Indefinite Theta Series for All Signatures

[Alexandrov, Banerjee, Manschot, Pioline 2016], [N 2016], [Westerholt-Raum 2016], [Kudla 2016], [Funke, Kudla 2017], [Zagier, Zwegers]

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$$\sum_{n \in \Lambda + \mu} \left[\frac{1}{2^{N_-}} \prod_{j=1}^{N_-} (\operatorname{sgn}(c_j \cdot n) - \operatorname{sgn}(c'_j \cdot n)) \right] q^{\frac{1}{2}n^2}$$

Modular completion is achieved by replacing

$$\prod_{j=1}^r \operatorname{sgn}(-f^{(j)} \cdot n) \rightarrow E_r((f^{(1)}, \dots, f^{(r)}), \sqrt{2\tau_2}x)$$

where the 'generalized error function' is defined by (in an orthonormal basis):

$$E_r(\mathcal{M}; u) := \int_{\mathbb{R}^r} d^r u' e^{-\pi(u-u')^T(u-u')} \operatorname{sgn}(\mathcal{M}^T u').$$

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[Imbimbo, Mukhi]: $H = \frac{p_i^2}{2} + \frac{1}{2} (\partial_i W(x))^2 - \frac{i}{2} \chi_1^i \chi_2^j \partial_i \partial_j W(x)$. Defining $y^j = \partial_j W(x)$ and B the range of this mapping, the susy partition function is:

$$N \int_B d^r y \left(\frac{\beta}{2\pi} \right)^{r/2} e^{-\beta(y^i)^2/2}.$$

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ex: Elliptic genera of squashed toric sigma models [Gupta, Murthy, N 2018].

Shadows or 'What are these functions?'

The $\bar{\tau}$ derivative of $E_r(F; \sqrt{2\tau_2}x)$ is:

$$-\frac{i}{\sqrt{2\tau_2}} \sum_{j=1}^r f^{(j)} \cdot x e^{2\pi\tau_2(f^{(j)} \cdot x)^2} E_{r-1}(F_{[r]/\{j\} \perp \{j\}}; \sqrt{2\tau_2}x)$$

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$$\partial_{\bar{\tau}} [\text{Indefinite theta series of signature } (N_+, N_-)] \\ = \tau_2^{-1/2} \sum [\text{Weight } \frac{3}{2} \text{ unary theta function}]^* \times$$

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- Mixed mock modular forms are defined as real analytic modular forms satisfying $\frac{\partial}{\partial \bar{\tau}} \hat{h}(\tau, \bar{\tau}) \in \bigoplus_j (\tau_2^{r_j} M_{k+r_j} \otimes \overline{M}_{2+r_j})$.

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- Indefinite theta functions lead us to a more general space where $\bar{\tau}$ derivative could involve mock modular forms (of higher depth).

An Example: Vafa-Witten Invariants

$U(N)$ Vafa-Witten Invariants on \mathbb{P}^2 :

- $U(2)$ case: $\frac{h_{2,\alpha}(\tau)}{\eta(\tau)^6}$, $\alpha = 0, 1$ where $h_{2,\alpha}$ is three times the Hurwitz class number generator function. [Vafa, Witten 1995]

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- $U(3)$ case: $\frac{h_{3,\mu}(\tau)}{\eta(\tau)^9}$, $\mu = 0, 1, -1$ where [...]

$$h_{3,0}(\tau) = \frac{1}{9} - q + 3q^2 + 17q^3 + 41q^4 + 78q^5 + 120q^6 + 193q^7 + \dots,$$

$$h_{3,1}(\tau) = h_{3,-1}(\tau) = 3q^{\frac{5}{3}} + 15q^{\frac{8}{3}} + 36q^{\frac{11}{3}} + 69q^{\frac{14}{3}} + 114q^{\frac{17}{3}} + \dots$$

These functions can be written in terms of signature $(2, 2)$ indefinite theta series [Manschot 2014, 2017].

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- Modular completion can be written as:

$$\widehat{h}_{3,\mu}(\tau, \bar{\tau}) = h_{3,\mu}(\tau) - \frac{3\sqrt{3}i}{2\sqrt{2}\pi} \sum_{\alpha \bmod 2} \int_{-\bar{\tau}}^{i\infty} dw \frac{\widehat{h}_{2,\alpha}(\tau, -w) \vartheta_{\frac{2\mu+3\alpha}{6}}(3w)}{(-i(w+\tau))^{\frac{3}{2}}}.$$

Circle Method

Goal: A Rademacher series for $f_{3,\mu}(\tau) := \frac{h_{3,\mu}(\tau)}{\eta(\tau)^9}$. [Bringmann, N 2018]
extending results of [Bringmann, Manschot 2010] on mixed mock modular forms.

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Writing $f_{3,\mu}(\tau) = \sum_{n=0}^{\infty} \alpha_{3,\mu}(n) q^{n-\Delta_\mu}$ we have

$$\alpha_{3,\mu}(n) = \int_i^{i+1} f_{3,\mu}(\tau) e^{-2\pi i(n-\Delta_\mu)\tau} d\tau.$$

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If $f_{3,\mu}(\tau)$ were a weakly holomorphic modular form, this would be an easy task. We would use an integration path that gets closer and closer to the real line near the rationals and use $f_{3,\mu}$'s behavior at $i\infty$:

$$f_{3,\mu}(\tau) = \frac{1}{9} q^{-\frac{3}{8}} \delta_{\mu,0} + \text{non-polar terms.}$$

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That is not enough to determine the behavior near rationals, e.g.

$$f_{3,\mu}\left(-\frac{1}{\tau}\right) = \tau^{-3/2} \sum \psi_{\nu,\mu} \left[f_{3,\nu}(\tau) + c_1 \frac{h_{2,\alpha}(\tau)}{\eta(\tau)^9} \int_0^{i\infty} \frac{\vartheta_{\frac{2\nu+3\alpha}{6}}(3w)}{(-i(w+\tau))^{\frac{3}{2}}} dw \right. \\ \left. + c_2 \frac{1}{\eta(\tau)^9} \int_0^{i\infty} \int_{w_1}^{i\infty} \frac{\vartheta_{\frac{\alpha}{2}}(w_2) \vartheta_{\frac{2\nu+3\alpha}{6}}(3w_1)}{(-i(w_2+\tau))^{\frac{3}{2}} (-i(w_1+\tau))^{\frac{3}{2}}} dw_2 dw_1 \right].$$

Application to Vafa-Witten Invariants

Each of $f_{3,\nu}(\tau)$, $\frac{h_{2,\alpha}(\tau)}{\eta(\tau)^9}$ and $\frac{1}{\eta(\tau)^9}$ has $q^{-3/8}$ pieces that contribute to the Fourier coefficients. How about the integrals?

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For this purpose, we find Mordell-type representations, e.g. the two dimensional integral can be expressed as

$$\int_{\mathbb{R}^2} d^2w g(w) q^{\frac{1}{3}(w_1^2+w_2^2+w_1w_2)}.$$

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Now, the circle method can be used with these modified polar terms.

- Fourier coefficients of depth-two $f_{3,\mu}(\tau)$ are expressed as a series, taking as input only the polar terms of the involved holomorphic functions and some information on the shadow.
- The mock terms in the modular transformation contribute power-law suppressed terms to the leading asymptotics of Fourier coefficients.

Thank you!