The Role of Noise and Dissipation in the Hadronization of the QGP

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Introduction and Motivation

 Depending on the nature of the QCD transition, phase conversion can occur in a number of ways: nucleation vs. spinodal decomposition

- Time scales in HIC -> hadronization via explosive decomp.
- Most theoretical attempts: rapid change in effective potential (Polyakov loop model, effective chiral models, etc) -> very fast spinodal decomposition ("explosive phase conversion")

[•] Signatures: e.g., Scavenius, Dumitru & Jackson (2001); Dumitru & Pisarski (2002); Randrup (2004), Rafelski (talk here).

<u>Questions</u>:

- What are the effects of dissipation and noise on the evolution of the order parameter for the chiral transition?
- Can dissipation prevent explosive decomposition in high-energy heavy ion collisions?

<u>Method</u>: Langevin description inspired by microscopic nonequilibrium field theory results (real-time lattice simulations)

Main point: even if the system quickly reaches the unstable (spinodal) region there is still no guarantee that it will explode!

Langevin dynamics and dissipation effects [ESF & G. Krein (2005)]

Microscopic QFT -> noise and dissipation terms from thermal & quantum fluctuations - self-interaction or coupling to other fields [e.g., Gleiser & Ramos (1994); Rischke (1998)]

$$\partial^2 \phi + \Gamma \frac{\partial \phi}{\partial t} + U'_{eff}(\phi) = \xi(\vec{x}, t)$$

 $\Gamma = \Gamma(T): \text{ response coefficient (``intensity of dissipation'')}$ $\phi(\mathbf{x}, t): \text{ non-conserved order parameter (} \sigma \text{ field for } \mathbf{x} PT)$ Stochastic (noise) force [assumed Gaussian and white (*)]: $\langle \xi(\vec{x}, t) \rangle = 0 \qquad \langle \xi(\vec{x}, t) \xi(\vec{x'}, t') \rangle = 2\Gamma T \delta(\vec{x} - \vec{x'}) \delta(t - t')$

Effective theory

• Assumptions: 1^{st} order χ phase trans. (+ exp. system + finite size)

 Framework: coarse-grained Landau-Ginzburg effective potential [from linear σ model + quarks , N_f=2] [Csernai & Mishustin (1995); Scavenius & Dumitru (1999); ...]

$$L = \overline{q} \left[i \gamma^{\mu} \partial_{\mu} - g \sigma + i \gamma_{5} \tau \cdot \pi \right] q + \frac{1}{2} \left(\partial_{\mu} \sigma \partial^{\mu} \sigma + \partial_{\mu} \pi \cdot \partial^{\mu} \pi \right) - U(\sigma, \pi)$$

$$U(\sigma,\pi) = \frac{\lambda^2}{4} (\sigma^2 + \pi^2 - v^2)^2 - h_q \sigma$$

Parameters such that $SU(2) \otimes SU(2)$ broken in the vacuum, $\langle \sigma \rangle = f_{\pi}, \langle \pi \rangle = 0$, h_{q} from PCAC, masses, etc

Integrate over fermions (heat bath for the chiral fields)

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effective potential for $\phi = (\sigma, \pi)$



Some questions:

 Time scales: expansion rate, decay rate, domain coarsening – nucleation vs. spinodal decomposition (explosion?)



 Length scales: radius of the critical bubble, correlation length, size of the system
N.B.: finite size can affect decay rate and late-stage growth

• Inhomogeneities: chiral field inhomogeneities vs. fermionic det [ESF, A. Mócsy & B.G. Taketani, work in progress]

Explosive scenario:

Time scales (for nucleation):



[Scavenius, Dumitru, ESF, Lenaghan & Jackson (2001)]

$$\Gamma_{nucl} \sim e^{-F_B/T}$$

Phase conversion is likely to proceed via "explosive" spinodal decomposition!

+ similar results for instabilities from:

Rafelski & Letessier (2000) Dumitru & Pisarski (2001) Scavenius, Dumitru & Jackson (2001) Paech, Stoecker & Dumitru (2003) Randrup (2004)

> System would go directly to spinodal and explode, mostly bypassing bubble nucleation... -> Very short time scales

However, shouldn't dissipation effects be relevant?

Solving the Langevin dynamics on a lattice:

- Cubic space-like lattice with 64³ sites under p.b.c.'s
- Semi-implicit finite-difference scheme for time evolution (now using leap-frog)
- FFT for spatial dependence
- Temperature fixed to spinodal
- Lattice spacing: **a = 0.91 fm**
- Average over several realizations

Results:

For t < 5 fm/c the solution for the linearized eqn is very close to the complete one, then is "delayed" -> O(2) potential much shallower than complete V_{eff}





Even for conservative dissipation Γ =2T, the retardation effect for exponential growth is ~ 100% !!

For expansion times ~ 5 fm/c (~ RHIC times scales) there might be not enough time for the onset of the spinodal explosion !

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Still missing:

Finite-size effects: important right from the start in heavy-ion collisions [ESF and R. Venugopalan (2004)]

Expansion of the fluid (more "dissipation") [Scavenius, Dumitru & Jackson (2001)]

Effects from "Polyakov loop sector" [ESF, G. Krein & A.J. Mizher, work in progress]

Non-Markovian (memory) effects [ESF, T. Kodama, G. Krein & L. Palhares, to appear soon]

More complicated noise and dissipation from QFT [ESF, G. Krein & R.O. Ramos, to appear soon]

Improved Langevin description [work in progress]

Multiplicative noise:

> A more complete QFT description of nonequilibrium dissipative dynamics predicts the existence of additional terms in the Langevin equation.

- > In fact, dissipation effects should depend on the local density and, accordingly, the noise should contain a multiplicative term (fluctuation-dissipation theorem).
- Effects of multiplicative noise seem to be rather significant (e.g., in the Kibble-Zurek scenario of defect formation in one spatial dimension) [Antunes, Gandra & Rivers (2005)]

$$\partial^2 \phi + [\Gamma_1 + \Gamma_2 \phi^2] \frac{\partial \phi}{\partial t} + V'(\phi) = \xi_1(\vec{x}, t) + \xi_2(\vec{x}, t) \phi$$

$$\langle \xi_a(\vec{x},t)\xi_a(\vec{x'},t')\rangle = 2\Gamma_a T\delta(\vec{x}-\vec{x'})\delta(t-t')$$

Preliminary results:

- For degenerate double-well potential at given T
- Still in arbitrary units
- Thermalization is clearly retarded!



Non-Markovian corrections:

Linear response nonequilibrium QFT provides complicated memory kernels associated with dissipation instead of a simple Markovian term. The latter corresponds to a drastic simplification of the dissipation kernel.

> We obtain a systematic expansion in which higher-order terms correspond to more non-local contributions.

We start with the much simpler case of dissipative metastable quantum mechanics, where the memory kernel has its origin in the Feynman influence functional of the heat bath, and derive a (derivative) expansion for non-local corrections.

(approximations and important scales under control)

In the case of field theory, we phenomenologically generalize our results for quantum mechanics, obtaining a Langevin equation that incorporates the first (radiation) non-local correction.

<u>Final remarks</u>

Effects of dissipation seem to be important in the process of hadronization of the QGP. In particular, it could dramatically modify the explosive behavior.

- Multiplicative noise (and density-dependent dissipation) do play a role, and retard even more thermalization.
- Non-Markovian corrections can, in principle, be incorporated systematically in the Langevin evolution (simpler structure).

Still to include: chiral + Polyakov loop effective theory; expansion; finite-size effects; more realistic microscopic treatment of dissipation; connection to hydro evolution.