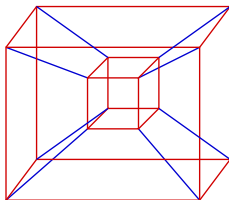


# Elliptic K3s, $T^4/\mathbb{Z}_2$ and Enriques involutions.

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The diagram shows a commutative square. On the top left is  $T^2$ , on the top right is  $K3$ , and on the bottom right is  $S^2$ . A horizontal arrow points from  $T^2$  to  $K3$ . A vertical arrow labeled  $\pi$  points from  $K3$  down to  $S^2$ . A curved arrow labeled  $\sigma$  points from  $S^2$  back up to  $K3$ .

- In particular: a  $K3$  described by a Weierstrass model.

# Weierstrass model

$$y^2 = x^3 + f_8(a, b)xz^4 + g_{12}(a, b)z^6$$

- $(y, x, z)$  are homogeneous coordinates of  $\mathbb{P}_{1,2,3}$ .

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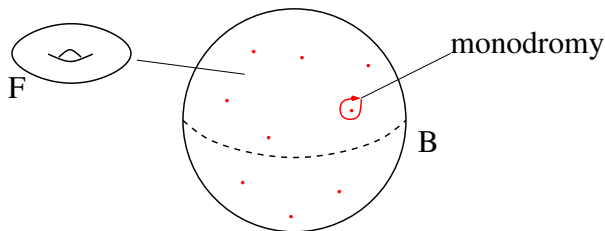
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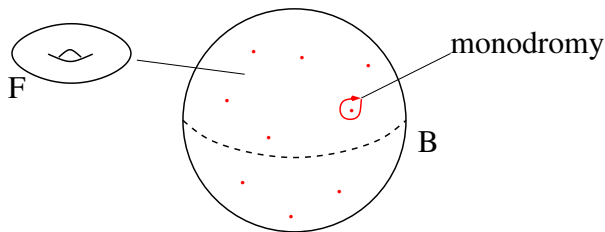
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$$B, F \in H^{1,1}(K3, \mathbb{Z}) \rightarrow \Omega_{2,0} \cdot B = \Omega_{2,0} \cdot F = 0.$$

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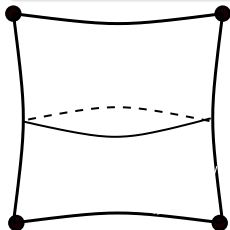
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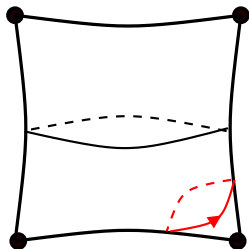
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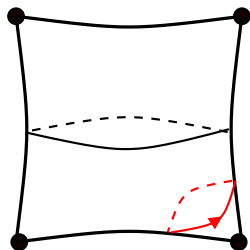


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This is like  $T^4/\mathbb{Z}_2$  !

The action of the Enriques Involution on  $T^4/\mathbb{Z}_2$  is known:  
(It is also known on the lattice of integral cycles of  $K3$ .)

$$(x_1, x_2, x_3, x_4) \mapsto (x_1, x_2 + \frac{1}{2}, -x_3, -x_4 + \frac{1}{2})$$



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**Let us use this to learn something about possible Enriques involutions on elliptic  $K3$  surfaces described by a Weierstrass model !**

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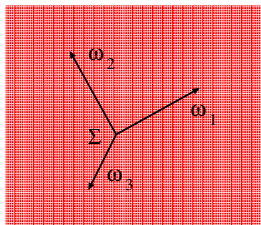
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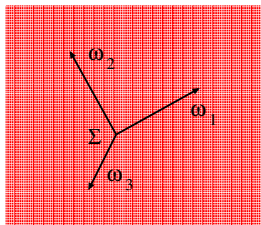
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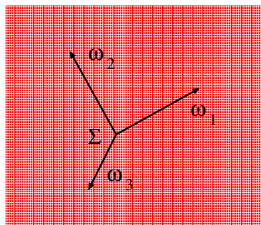


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$$J \mapsto J, \quad \Omega_{2,0} \mapsto -\Omega_{2,0} \quad \text{under the Enriques involution.}$$



# K3 moduli space

$\gamma_i \in H^2(K3, \mathbb{Z})$  with  $\gamma_i \cdot \gamma_i = -2$  (a “root”) orthogonal to  $\Sigma \rightarrow$   
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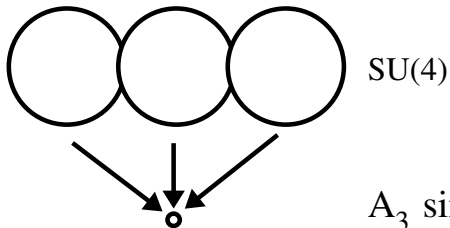
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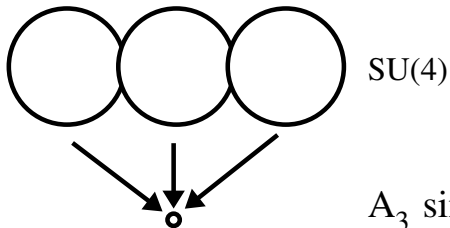


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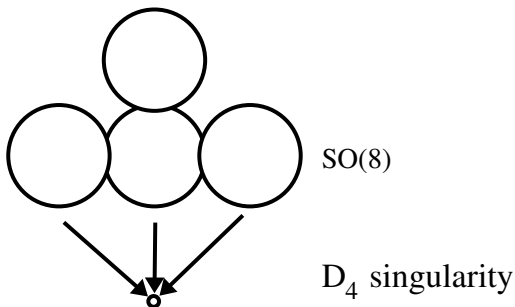
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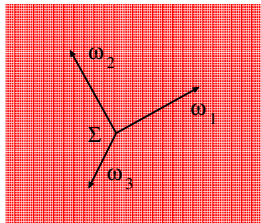
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# Wilson Lines

Put  $\Sigma$  is entirely in the  $U^{\oplus 3}$  lattice:

$$\omega_i = e_i + s_i e^i, \quad e_i \cdot e^j = \delta_i^j$$

As the whole  $-E_8 \times -E_8$  lattice is orthogonal to  $\Sigma$ , the singularity type/gauge enhancement is  $E_8 \times E_8$ .



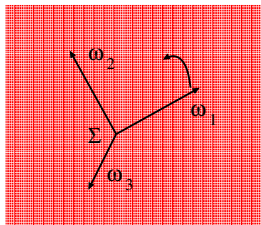
M-theory on this space is dual to  $\text{het}_{E_8 \times E_8}$  on  $T^3$ .

# Wilson Lines

Now rotate  $\omega_1$  into the  $E_8$  lattices:

$$\omega_1 = e_1 + s(e^1 + W),$$

All roots of the  $E_8$  lattice for which  $W \cdot \gamma_k = n \in \mathbb{N}$  still lead to shrunk cycles:  $\tilde{\gamma}_k = \gamma_k - ne_i$  now satisfies  $\gamma_k \cdot \omega_i = 0$ . This is the same condition as for Wilson-line breaking.



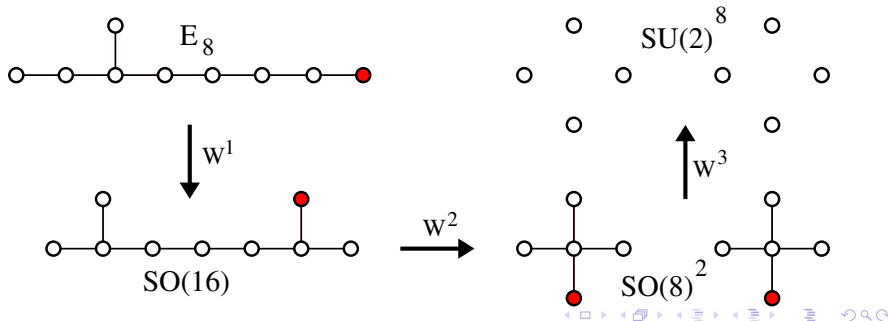
The resolution of singularities works like Wilson-line breaking !

# Wilson Lines

We can reach  $T^4/\mathbb{Z}_2$  by three Wilson-lines that break  $E_8 \times E_8 \rightarrow SU(2)^{16}$

$$W^1 = (1, 0^7, \dots) \quad W^2 = (0^4, \frac{1}{2}, \frac{1}{2}, \dots)$$

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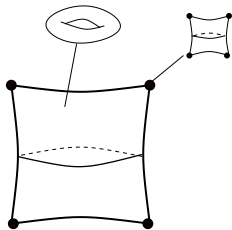
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This gives an (explicit) embedding of

$$A_1^{\oplus 16} \oplus U(2)^3 \subset U^{\oplus 3} \oplus -E_8^{\oplus 2}.$$





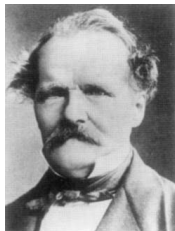
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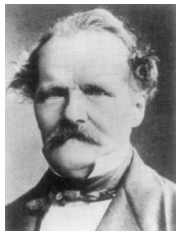
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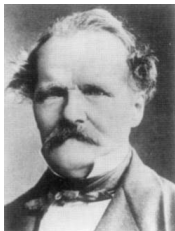
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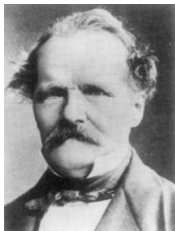
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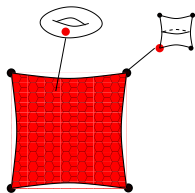


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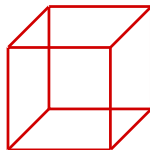
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- Besides the  $A_1^{\oplus 16} \oplus U(2)^3$  lattice we find naively, there are extra cycles that stem from the pillows.
- They complete  $A_1^{\oplus 16} \oplus U(2)^3$  to  $H^2(K3, \mathbb{Z}) = U^{\oplus 3} \oplus -E_8^{\oplus 2}$
- Over the reals, they can be expressed as

$$\sigma_{ij}^k = \frac{1}{2} \cdot (\pi_{ij} - \sum_i C_i).$$

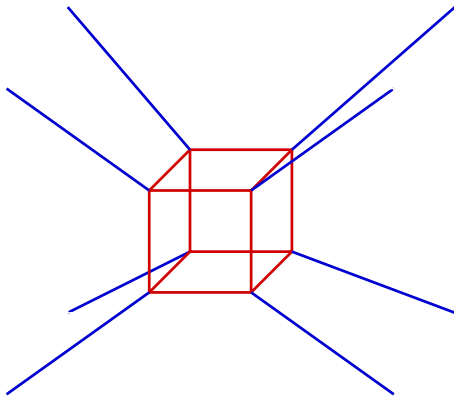


We can visualize these relations by drawing a 4D hypercube:



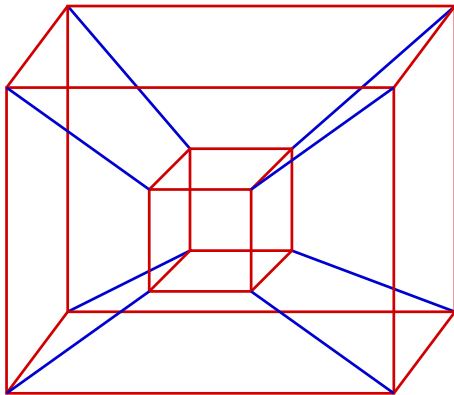
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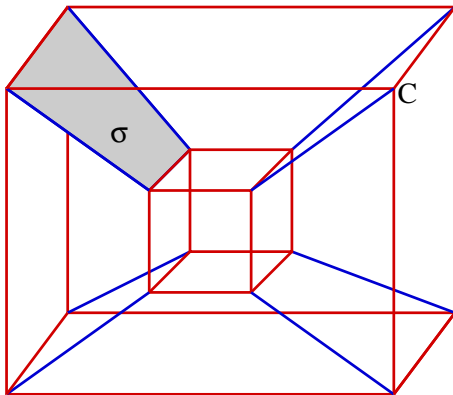
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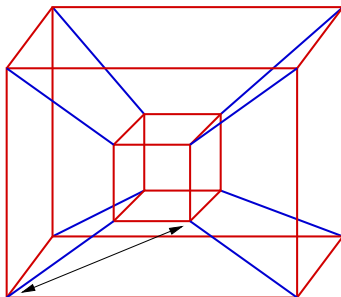
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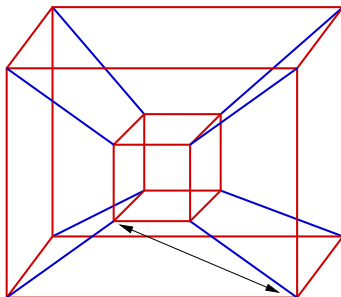
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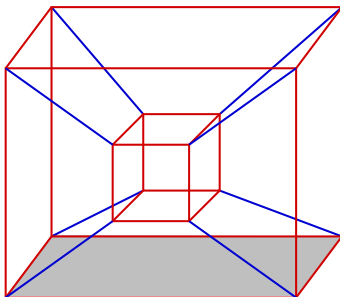




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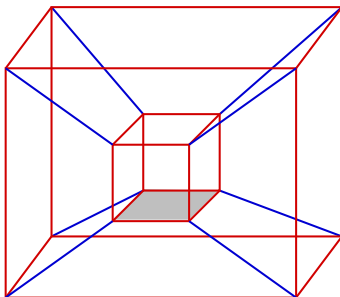
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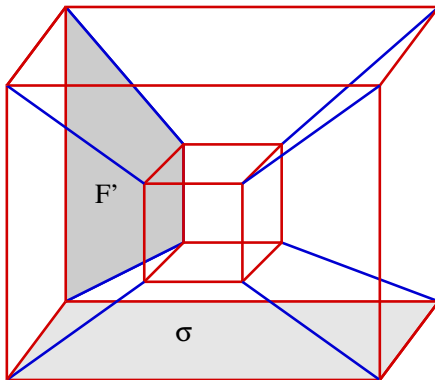
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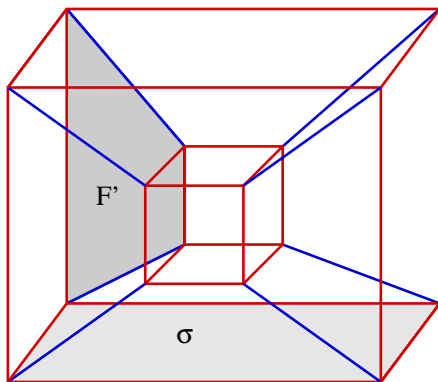
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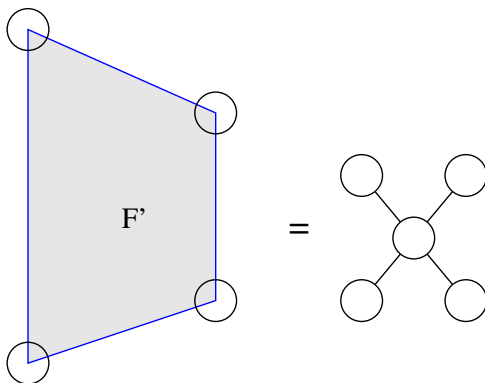
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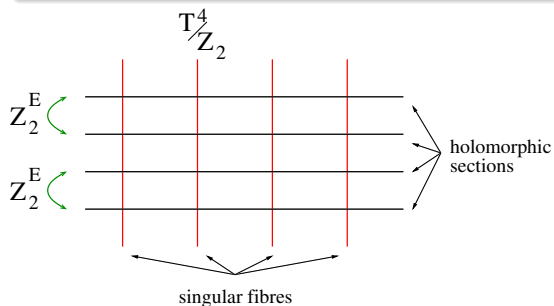
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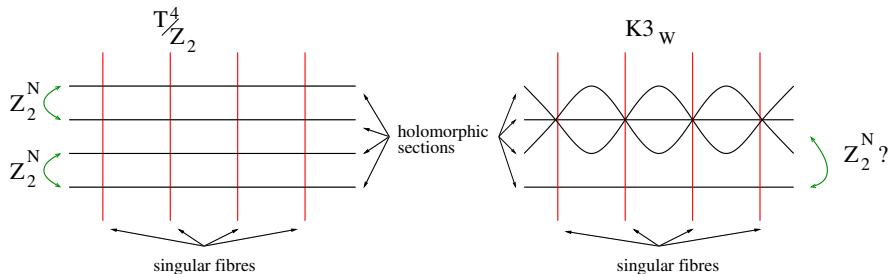
Going from  $T^4/\mathbb{Z}_2$  to the Weierstrass model means blowing-up the singularities hitting the section and collapsing the singular fibres to produce the  $D_4 \sim SO(8)$  singularities.

The structure of the holomorphic sections and the singular fibres for  $T^4/\mathbb{Z}_2$  is symmetric under the Enriques involution:



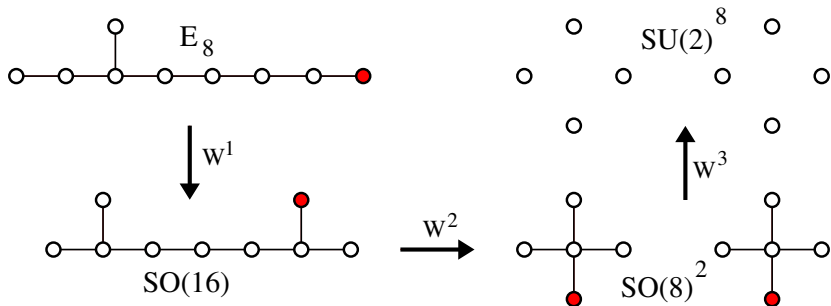
# Enriques revisited

$K3$  as described by the Weierstrass model does not allow an Enriques involution that keeps the holomorphic section:



# Enriques revisited

There is a further way to understand what goes wrong.  
Remember the Wilson lines:

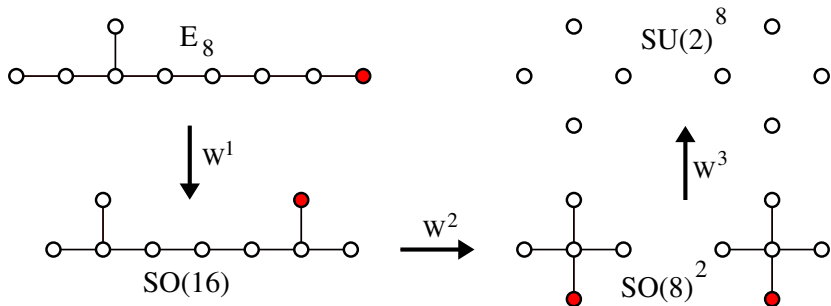


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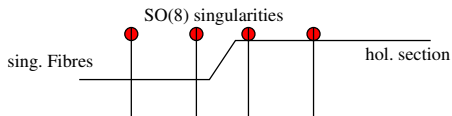
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We get to  $SO(8)$  by switching off  $W^3$ .

# Enriques revisited

- However, this is not the situation described by the Weierstrass model:

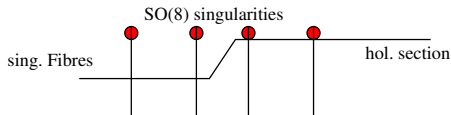
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- Each hol. sections meets two singularities.

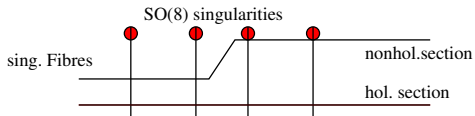


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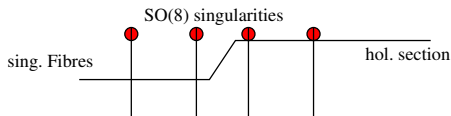


- We can get back to the Weierstrass model  $K3$  by a complex structure deformation.

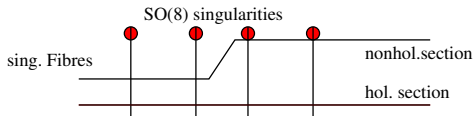


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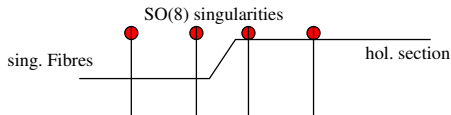
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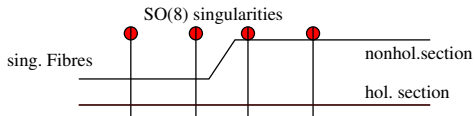
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 $\Omega \mapsto -\Omega$ .
- There is a clash  
hol. of the W. model section  $\leftrightarrow$  hol. of Enriques involution

# Conclusions

- Constructing the lattice of integral cycles of  $K3$  from a blow-up of  $T^4/\mathbb{Z}_2$  provides a nice picture for studying the action of Enriques involutions.

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# Bonus in progress

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- The Enriques involution acts as  $y \mapsto -y, z \mapsto -z$  and exchanges the two sections.
- This equation describes a ten-dimensional family of elliptic  $K3$  spaces, which agrees with the number of complex structure deformations of an Enriques surface.