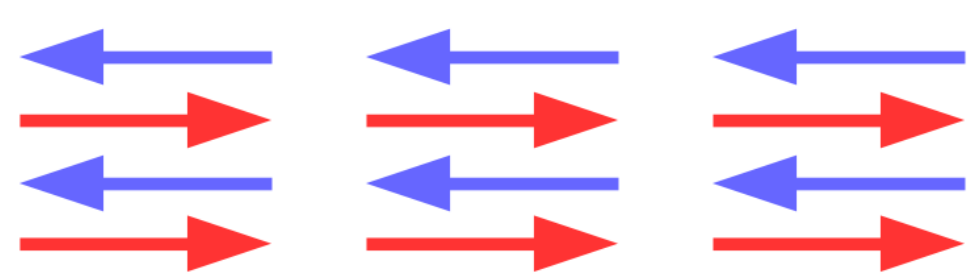


1 INTRODUCTION

- **Two-component superfluids** can be realized by mixtures of two different species at sufficiently low temperatures. Examples are
 - ³He-⁴He mixtures
 - Superfluid Bose-Fermi mixtures of ultra-cold atomic gases
 - Mixture of a neutron superfluid with a proton superconductor in the interior of **compact stars**
- A **counterflow** between fluids can be created experimentally (lab) or occurs necessarily (**compact stars**).



- Instabilities like **Landau's critical velocity** or the **two-stream instability** might be important for astrophysical phenomena like **pulsar glitches**. These instabilities can serve as a **trigger** for the collective unpinning of the neutron-superfluid vortices from the neutron lattice, which can be viewed as a second (normal) fluid [1].

2 MICROSCOPIC MODEL

- Two coupled, complex scalar fields with selfinteraction and two types of interspecies coupling: **derivative/entrainment coupling** and **non-entrainment coupling** [2, 3].

$$\begin{aligned}\mathcal{L} &= \mathcal{L}_1 + \mathcal{L}_2 + \mathcal{L}_I, \\ \mathcal{L}_i &= \partial_\mu \varphi_i \partial^\mu \varphi_i^* - m^2 |\varphi_i|^2 - \lambda |\varphi_i|^4, \\ \mathcal{L}_I &= 2h |\varphi_1|^2 |\varphi_2|^2 - g \varphi_1 \varphi_2 \partial_\mu \varphi_1^* \partial^\mu \varphi_2^*.\end{aligned}$$

- Superflows \vec{v}_i and the chemical potentials μ_i are both related to the phases of the condensates (expectation values of the fields, $\langle \varphi_i \rangle = \rho_i e^{i\psi_i}$) [4]:

$$\mu_i = \partial_0 \psi_i, \quad \vec{v}_i = \frac{\vec{\nabla} \psi_i}{\mu_i}.$$

- **Two-component superfluid:** spontaneous symmetry breaking with pattern $U(1) \times U(1) \rightarrow 1$ leads to **two Goldstone modes**

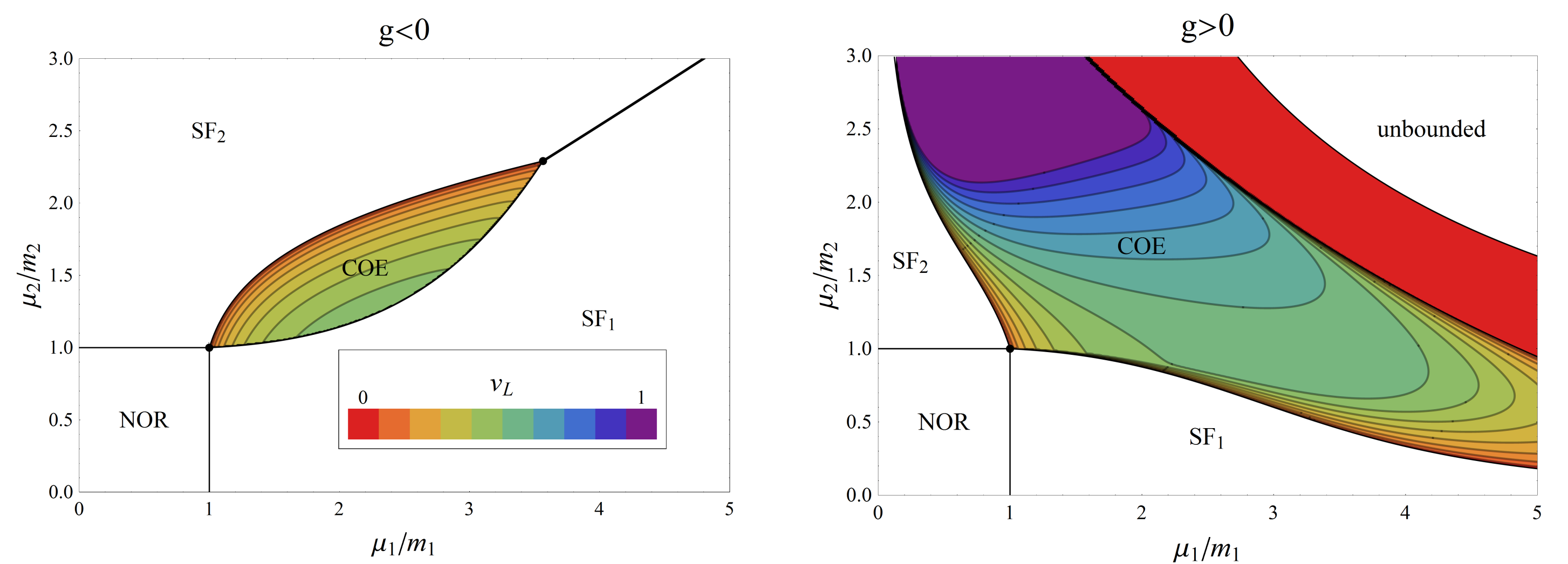
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3 PHASE DIAGRAMS & LANDAU'S CRITICAL VELOCITY

Landau's critical velocity v_L : **energy of Goldstone mode** becomes **negative**

Phase diagrams for positive and negative values of the entrainment coupling g in the plane of the chemical potentials. The colors represent Landau's critical velocity in units of $c = 1$. For simplicity, $h = 0$.



Four different phases, separated by second order phase transitions except for first order transition between SF₁/SF₂:

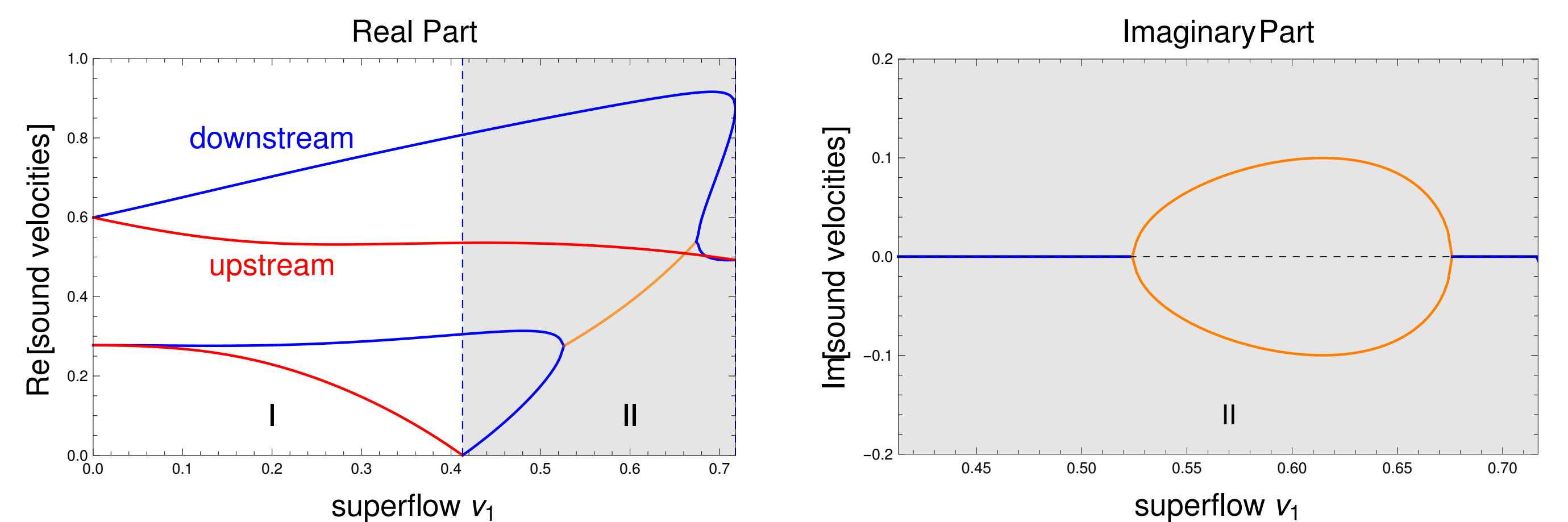
- NOR: **normal phase**
- SF_{1/2}: **superfluid_{1/2} phase**
- COE: **two-component superfluid** (coexistence phase)

For $g < 0$, the entrainment term acts as an energy penalty \Rightarrow size of COE phase reduced. For $g > 0$: energy gain \Rightarrow COE phase enlarged.

4 SOUND MODES & INSTABILITIES

- At $T = 0$, the sound modes are given by the slope of the Goldstone modes or can equivalently be calculated using relativistic, linearized hydrodynamics.

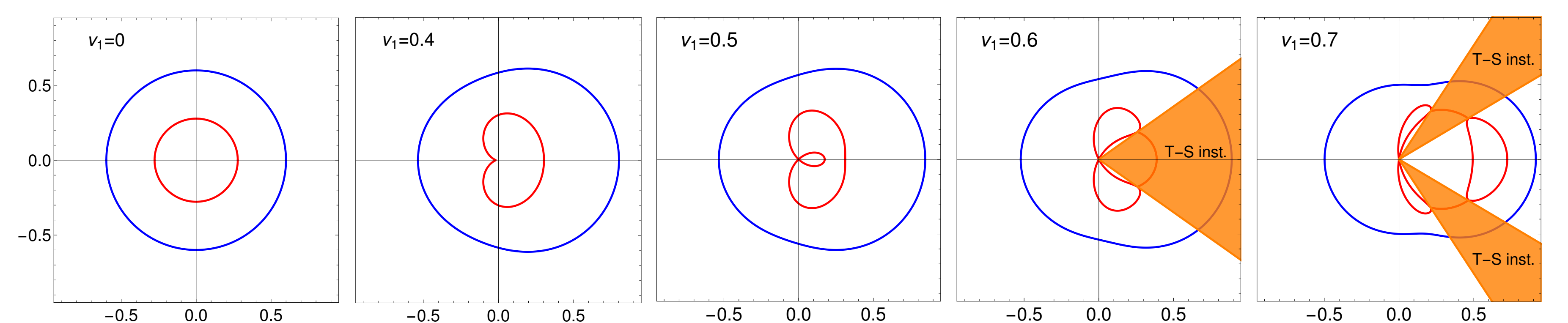
Real and imaginary parts of the sound velocities for $g < 0$ in up- and downstream direction.



- Two regions:

- I: Stable regime, two sound modes in **upstream** direction (superflow antiparallel to sound wave) and two in **downstream** direction
- II: Energetically unstable regime (shaded region), one **upstream** mode flips to **downstream** (analogy: r-mode instability in rotating star). **Orange line**: two solutions merge and become complex (see right panel) \Rightarrow onset of **two-stream instability**

- **Polar plots** show sound modes and development of instabilities for all angles between the superflow \vec{v}_1 and the wave vector \vec{k}



- **Two-stream instability** starts to develop in **downstream** direction "after" Landau's critical velocity

5 SUMMARY

- **Two-component superfluids** with non-zero superflow feature several instabilities, which can be calculated by analyzing the sound modes.
- They might serve as a **trigger for pulsar glitches**.
- It is found that the **two-stream instability** always takes place in an **energetically unstable** regime (after Landau's critical velocity). This is different for a two-component **normal fluid** (not shown here).

CONTACT & FUNDING

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