

Phases of Flat Space Higher Spin Gravity

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D Grumiller, M Riegler, J Rosseel



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Outline

- 1 Motivation
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- 4 Phases of Spin-3 Flat Space
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Motivation

- Holography in flat space
- Higher spins interesting as possible unbroken/Hagedorn phase of String Theory
- Interesting to evade No-Go theorems
- Phase structure of 3D HS gravity in flat space

Higher Spins in $D > 3$

- Higher spin particles with long-range interactions in $D > 3$ forbidden by soft-theorems, Coleman-Mandula, Weinberg-Witten
 - evade through long-range cutoff implemented by cosmological constant (dS or AdS) rather than mass
- No-Go theorems on forms of interaction vertices involving higher spins
 - evade through non-minimal interactions and infinite tower of spins

Higher Spins in AdS₃

- Generalization of first order formulation of gravity

$$S = \frac{k}{4\pi} \int \left\langle A \wedge dA + \frac{2}{3} A \wedge A \wedge A \right\rangle$$

where $\langle \cdot, \cdot \rangle$ is a non-degenerate bilinear form.

- For pure gravity, gauge group $G = \mathfrak{sl}(2) \oplus \mathfrak{sl}(2)$
- $G = \mathfrak{sl}(N) \oplus \mathfrak{sl}(N)$ gives spins $2, \dots, N$
- $G = \mathfrak{hs}(\lambda) \oplus \mathfrak{hs}(\lambda)$ gives infinite tower of spins, direct analog of higher spin theories in $D > 3$
- Natural \mathbb{Z}_2 grading on algebra, (zu-)Vielbein and Spin connection given by even and odd parts of A

Holography in 3D flat space

- The large ℓ limit of AdS is flat space
- Taking this limit yields a contraction of the gauge group

$$\begin{aligned}\mathfrak{sl}(2) \oplus \mathfrak{sl}(2) &\rightarrow \mathfrak{isl}(2) \\ \mathfrak{sl}(N) \oplus \mathfrak{sl}(N) &\rightarrow \mathfrak{isl}(N)\end{aligned}$$

- $\mathfrak{isl}(2)$ theory is pure gravity in flat space
- $\mathfrak{isl}(N)$ theory is gravity in flat space coupled to higher spins $3, \dots, N$

Flat Space Cosmologies

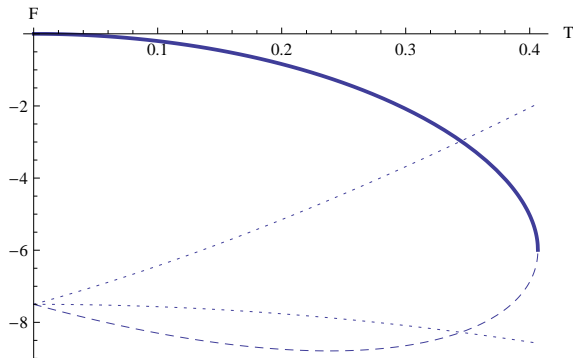
- $l \rightarrow \infty$ limit of BTZ black hole
- Outer horizon goes to infinity
- Inner horizon becomes cosmological horizon
- Have finite energy, entropy
- For pure gravity, described by 2 parameters (M, J) or (T, Ω)
- For each additional spin, 2 additional parameters

Higher Spin Flat Space Cosmologies

- 4 branches of solutions¹
- Branch 1 connects continuously to the pure gravity solution in appropriate limit
- First order phase transitions between branches 1 and 2
- Also Hawking-Page phase transition to hot flat space

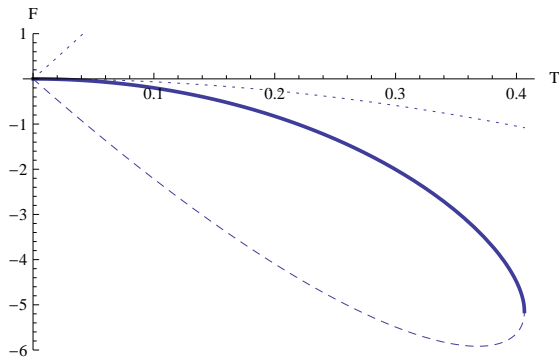
¹Assuming simplest solution of holonomy conditions

$$0 < \frac{\Omega \Omega_V}{\Omega_U} < \frac{3}{2}$$



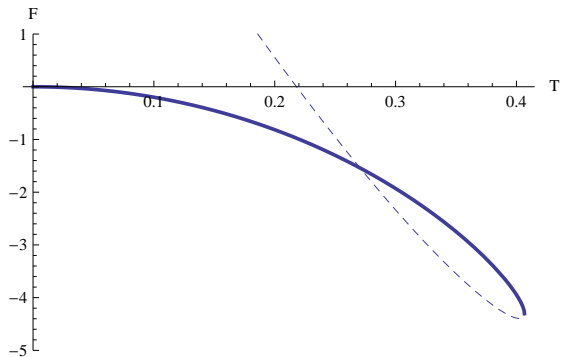
branch 1 is thermodynamically unstable for all temperatures

$$\frac{\Omega\Omega_V}{\Omega_U} = \frac{3}{2}$$



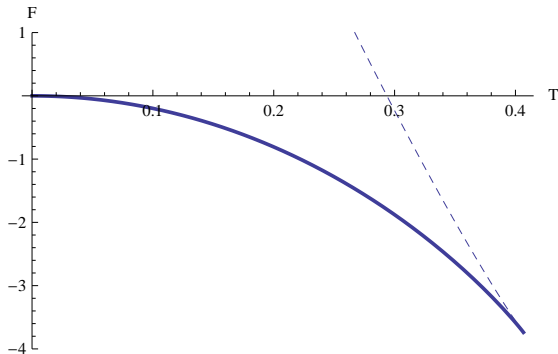
branches 1 and 2 degenerate at $T = 0$ and branch 1 is thermodynamically unstable at all $T > 0$

$$\frac{3}{2} < \frac{\Omega\Omega_V}{\Omega_U} < 2$$



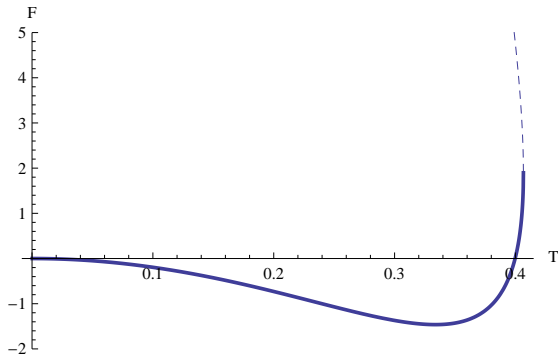
branches 1 and 2 degenerate at some $T_c > 0$. Below T_c branch 1 is stable (up to Hawking-Page), above T_c branch 1 is unstable

$$\frac{\Omega\Omega_V}{\Omega_U} = 2$$



branches 1 and 2 degenerate at the maximal temperature (when the discriminant vanishes), branch 1 is stable below the maximal temperature (up to Hawking-Page)

$$\frac{\Omega\Omega_V}{\Omega_U} > 2$$



branches 1 and 2 degenerate at the maximal temperature (when the discriminant vanishes), branch 1 is stable below the maximal temperature (up to Hawking-Page)

Future Work

- Understanding possible 0^{th} order phase transitions
- Understanding or eliminating $2\pi\mathbb{N}$ conical surplus solutions

Thank You