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Emergent Locality in the AdS/CFT Correspondence

arXiv:0903.4437 (MG, Giddings, Penedones)

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work in progress (MG, Giddings)

Acknowledgements

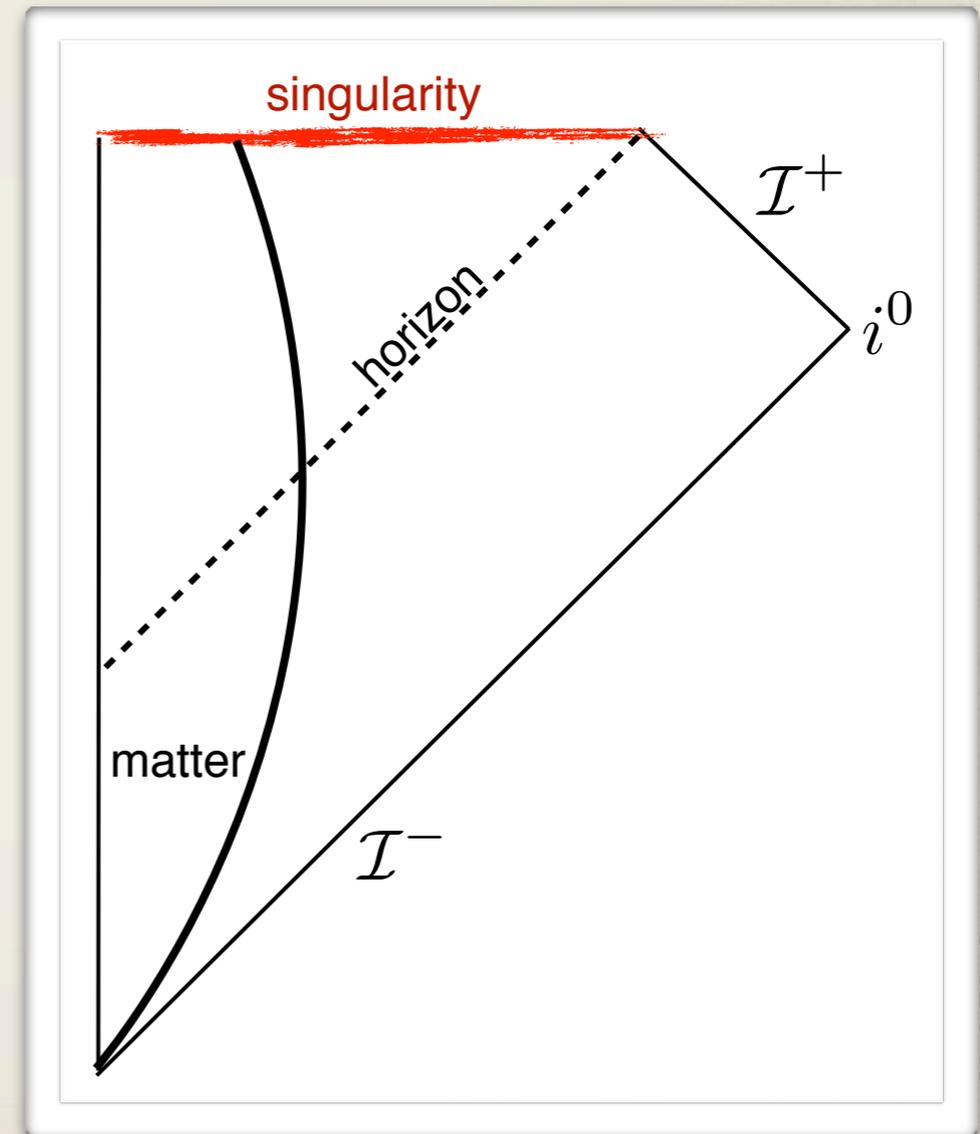
- * My committee: Joe Polchinski, Harry Nelson
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- * My family

Overview

- * Black hole evaporation
- * The AdS/CFT correspondence
- * Scattering in the large R limit
 - * Necessity of wavepackets
 - * Construction
 - * Singularity structure
- * Limits on locality
 - * Can we do better?
- * Understanding Bulk Unitarity
- * Conclusions & Open Questions

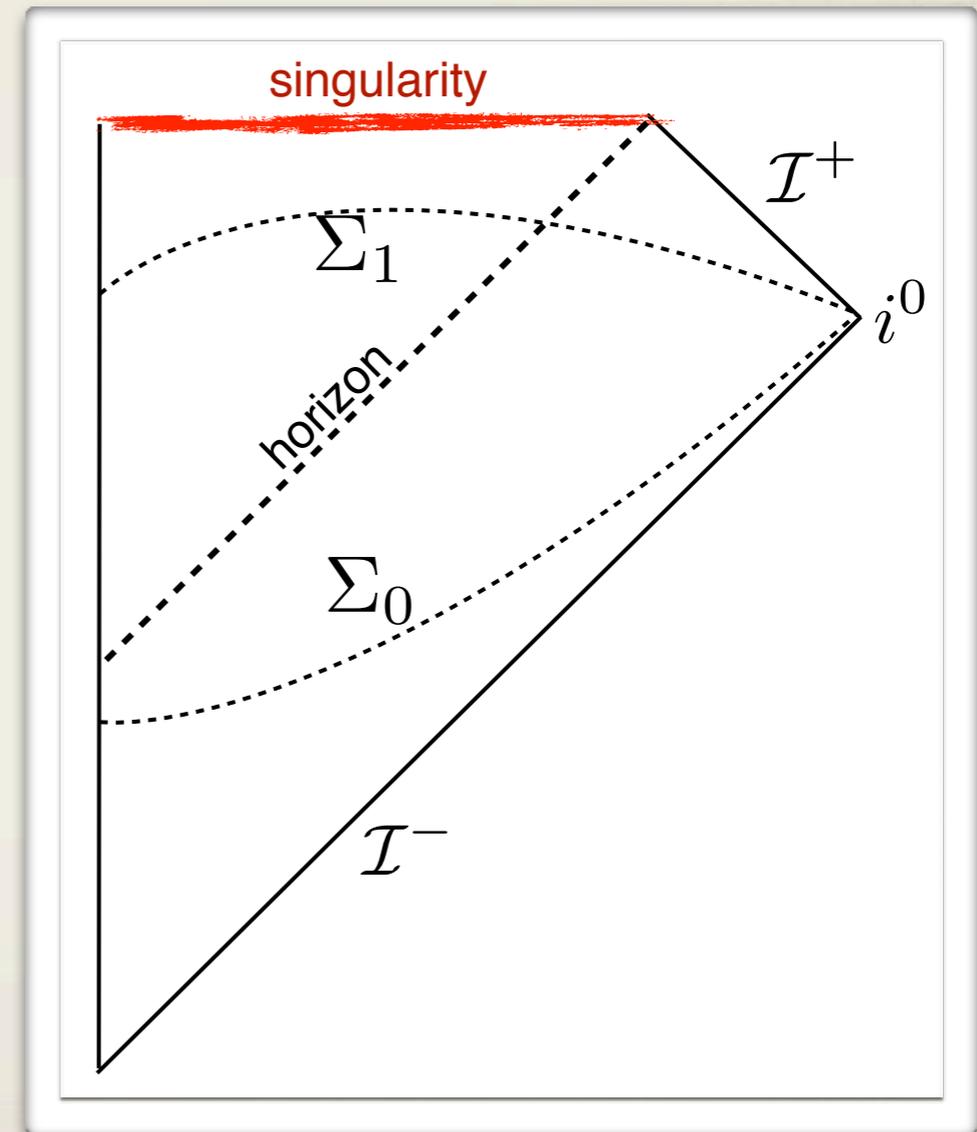
Evaporating Black Holes

- * Classically, everything that falls in is lost forever.
- * Hawking: Black holes aren't completely black.
- * Hawking Radiation \Rightarrow black holes evaporate.



Evaporating Black Holes

- * State on Σ_0 can be pure.
- * Locality \Rightarrow trace over sub-sector of Σ_I inside horizon \Rightarrow mixed state outside.
- * Pure state evolving into mixed state violates unitarity.



Evaporating Black Holes

* 3 Possible Resolutions:

* Non-Unitary evolution

$$\rho \mapsto \mathcal{S}\rho$$

* Remnants

Black hole only evaporates to Planck scale

* Non-locality

Non-Unitary Evolution?

$$\rho \mapsto \mathcal{S}\rho$$

- * Information transfer requires energy.
- * Information loss \leftrightarrow energy loss.
- * Virtual effects \Rightarrow Planck-scale energy non-conservation.
- * Banks, Peskin, Susskind (1984):
Nonunitarity \Rightarrow thermal bath at $T \sim M_P$

Remnants?

- * Long lived Planck-scale remnant: if remnant decays and information gets out, takes

$$t_{\text{decay}} = \frac{S^2}{M_P}$$

- * Form black holes from arbitrarily many initial states, so arbitrarily many remnant species \Rightarrow arbitrarily large production cross-section.

Non-Locality?

- * Physical observables must be gauge invariant.
- * In gravity, this means observables must be diffeomorphism invariant.
- * There are no diffeomorphism invariant local observables in gravity. (Torre 1993)

$$\delta\mathcal{O}(x) = \epsilon^\mu \nabla_\mu \mathcal{O}(x)$$

Holographic Principle

- * Hints that locality should be given up.
- * BH Entropy grows with bounding area, not volume.
- * Bousso (1999): trying to access too many states in a fixed volume \rightarrow black hole formation.
- * Look for a non-local formulation of Quantum Gravity.
- * Area law \Rightarrow should look for a theory in one fewer dimensions.

AdS/CFT

- * Quantum Gravity in asymptotically Anti-de Sitter space (AdS) is conjectured to be dual to a Conformal Field Theory (CFT) living on the boundary. (Maldacena 1997)

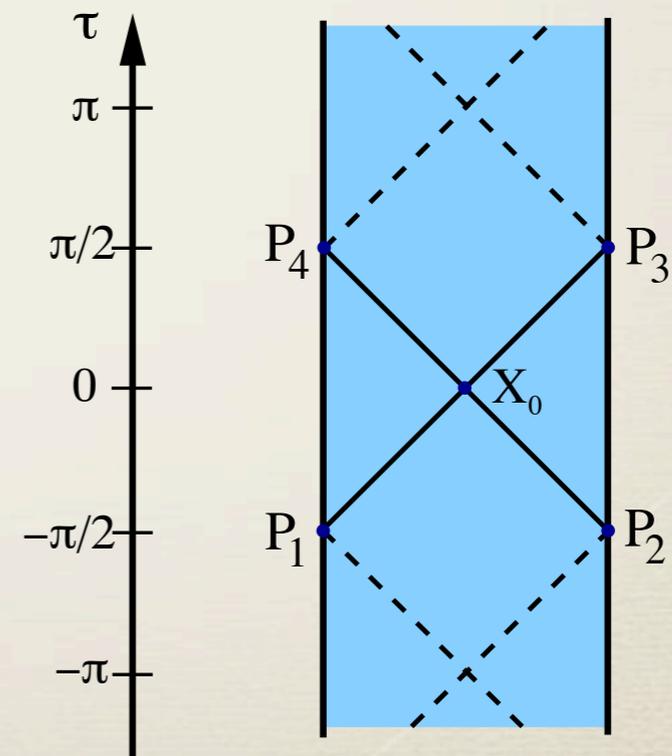
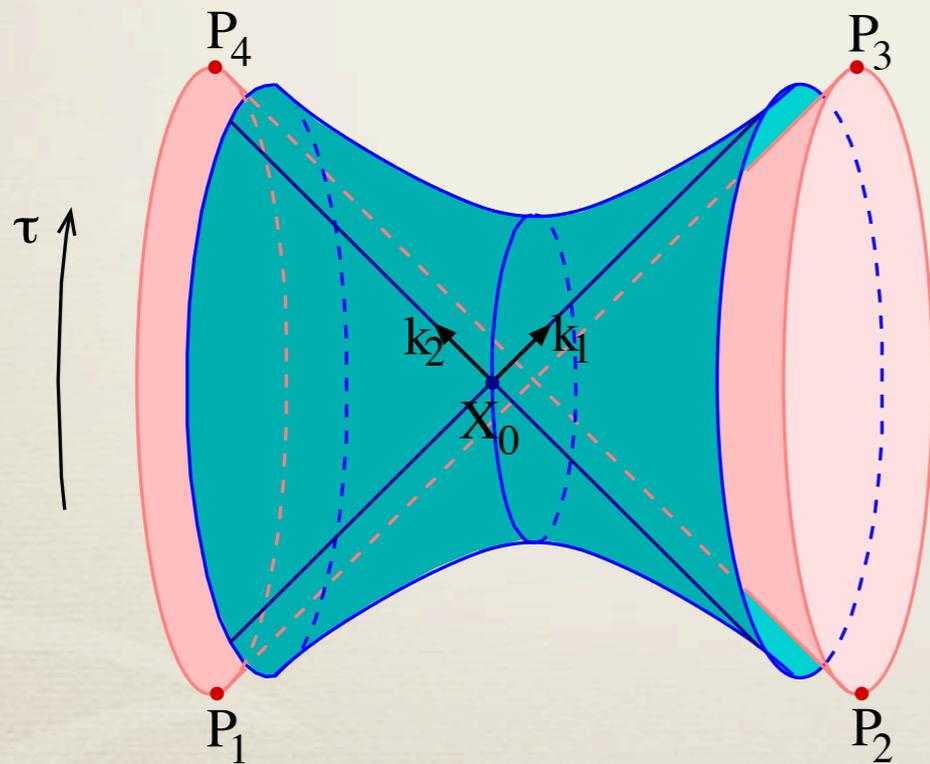
$$\left\langle \exp \left(i \int_{\partial \text{AdS}} \alpha_{\phi} \mathcal{O}_{\phi} \right) \right\rangle_{\text{CFT}} = Z_S [\alpha_{\phi}]$$

- * Operator insertions in the CFT \leftrightarrow boundary conditions for fields in AdS.

AdS Geometry

*
$$ds^2 = \frac{R^2}{\cos^2 \rho} (-d\tau^2 + d\rho^2 + \sin^2 \rho d\Omega_{d-1}^2)$$

* Universal cover of hyperboloid of radius R in $M_{2,d}$



AdS/CFT Dictionary

- * Gubser, Klebanov, Polyakov; Witten (1998): Boundary conditions on fields in AdS \leftrightarrow operator insertions & VEVs in dual CFT.
- * Fields in AdS have normalizable & non-normalizable modes:
$$\phi \sim \cos^{2h_-} \rho \alpha(\tau, \Omega) + \dots + \cos^{2h_+} \rho \beta(\tau, \Omega)$$
- * Non-normalizable mode \leftrightarrow operator insertion:
$$\mathcal{L}_{\text{CFT}} \mapsto \mathcal{L}_{\text{CFT}} + \alpha_\phi \mathcal{O}_\phi$$
- * Normalizable mode \leftrightarrow operator expectation value:
$$\langle \mathcal{O}_\phi \rangle = \beta_\phi$$

Large R Limit

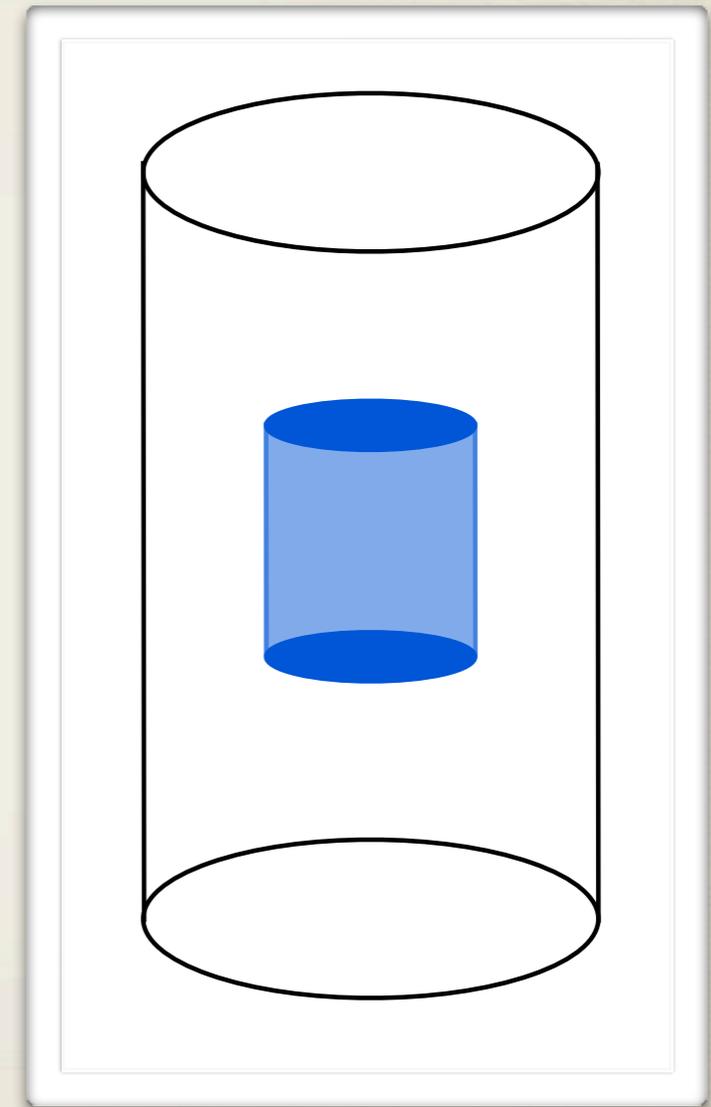
* $t = R\tau \quad r = R\rho \quad R \rightarrow \infty$

* Approximately flat
 $ds^2 \rightarrow -dt^2 + dr^2 + r^2 d\Omega^2$

* Normalizable frequencies
 $\omega_{nl} \rightarrow \omega R$

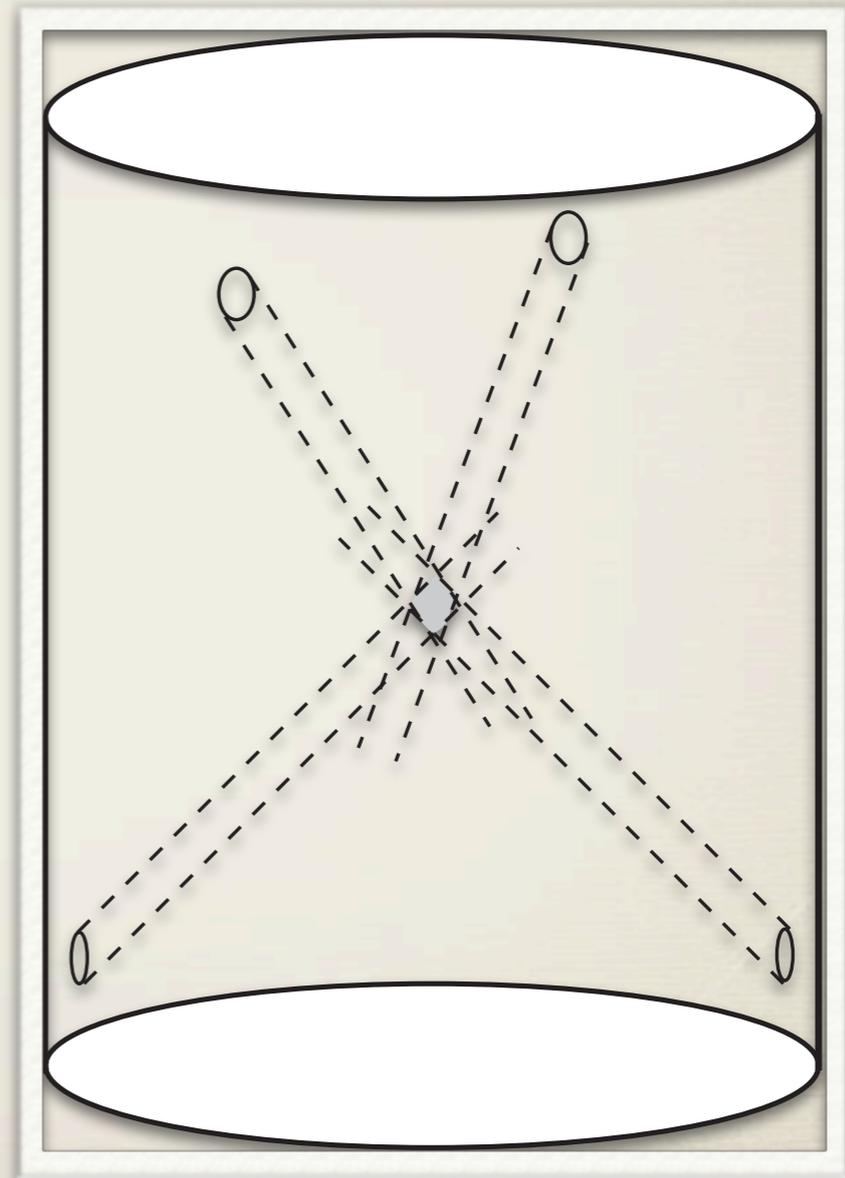
* Normalizable wavefunctions

$$\phi_{nl\vec{m}} \rightarrow \frac{\sqrt{2\omega^{d-1}}}{(\omega r)^{\frac{d}{2}-1}} J_{l+\frac{d}{2}-1}(\omega r) Y_{l\vec{m}}(\Omega)$$



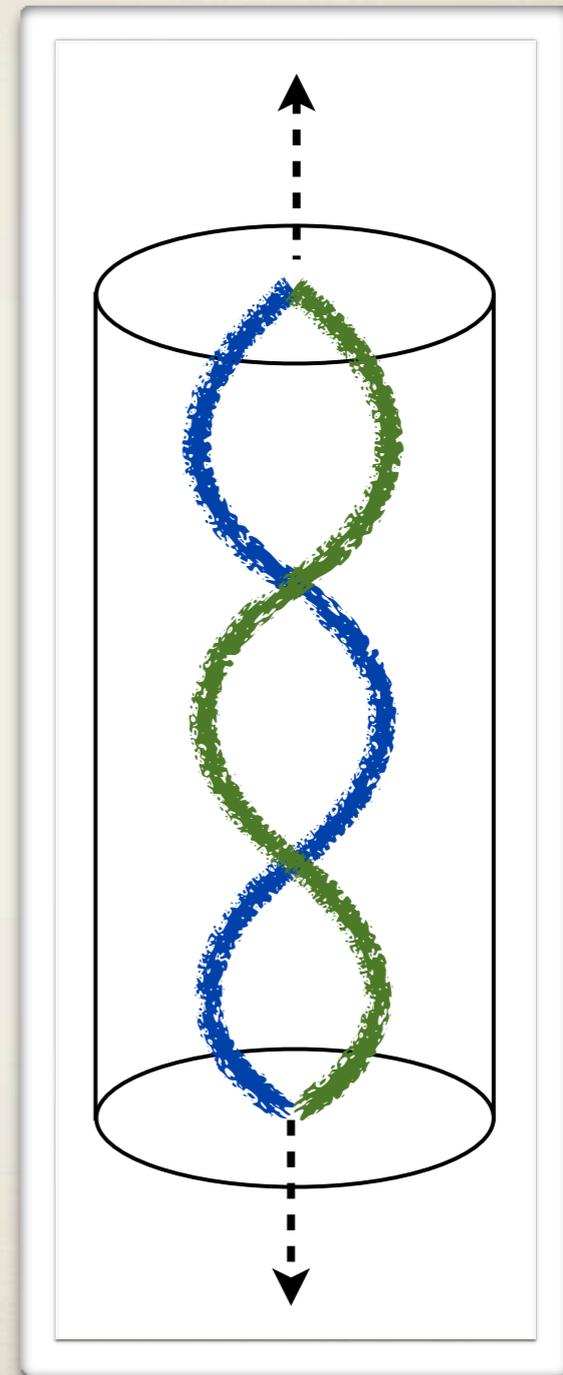
Scattering in the Flat Region

- * Scattering in the flat region should approximate local physics in our universe.
- * Need to construct wavepackets to localize scattering to a single flat region of AdS.



Multiple Scattering

- * Free fields in AdS are periodic.
- * Purely normalizable states will interact infinitely many times.
- * Can't isolate contribution from one scattering experiment.



Interactions Near the Boundary

- * Boundary sources \rightarrow infinite particle production near the boundary.
- * Single particle states not well defined when boundary sources turned on.

$$N = \int_{t_0} d^d \vec{x} \sqrt{-g} (\phi^* \overleftrightarrow{\partial}_t \phi) = \infty$$

- * Difficult to isolate scattering in flat region from scattering near the boundary.
- * Sources should be compact and non-overlapping to avoid infinite interactions near the boundary and normalize states.

Boundary-Compact Wavepackets

- * Construct using Bulk-Boundary Propagator.

$$\phi_f(x) = \int db f(b) G_{B\partial}(b, x)$$

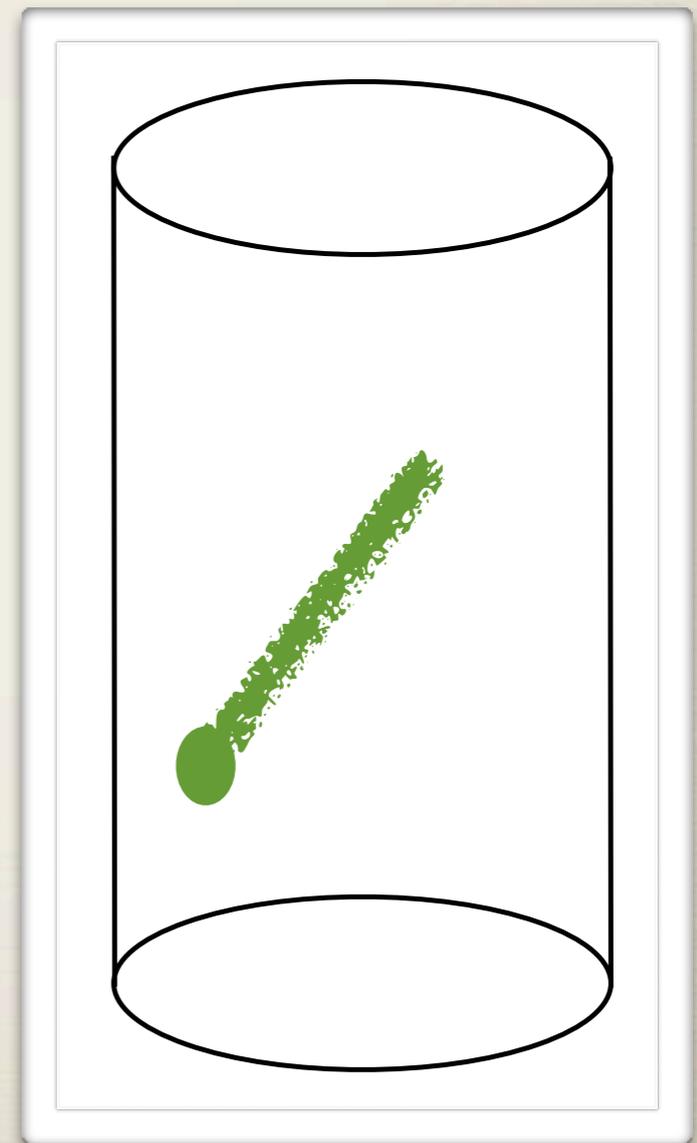
- * Compactly supported sources

$$f(b) = L \left(\frac{\tau - \tau_0}{\Delta\tau} \right) L \left(\frac{\theta}{\Delta\theta} \right) e^{-i\omega R(\tau - \tau_0)}$$

of size $\Delta\tau$, $\Delta\theta$.

- * $\frac{1}{\omega R} \ll \Delta\tau, \Delta\theta \ll 1$

- * Scatter when sources turned off.



Wavepackets in the Scattering Region

- * Near the center of AdS,

$$\phi_f(x) \approx \phi_f(0) \frac{\tilde{L}_{d-1}(x_{\perp} \omega \Delta \theta)}{\tilde{L}_{d-1}(0)} L\left(\frac{u}{\Delta t}\right) e^{-i\omega u}.$$

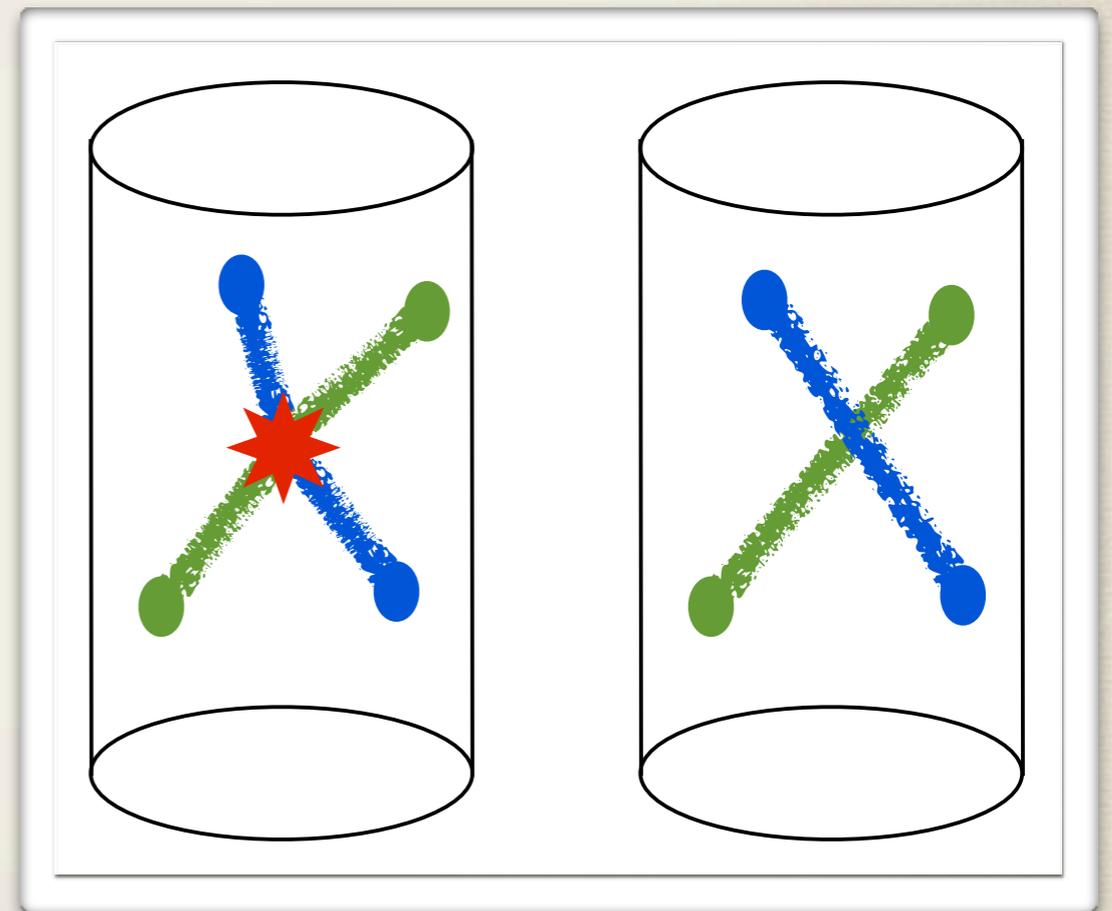
- * Longitudinal width $\Delta t \sim R\Delta\tau$.

- * Transverse width $\Delta x_{\perp} \sim 1/(\omega\Delta\theta)$.

- * Well localized for $1/\omega \ll \Delta t, \Delta x_{\perp} \ll R$, equivalent to earlier requirement.

Singularity Structure

- * Signal of interaction from intersecting wavepackets:
local bulk physics!



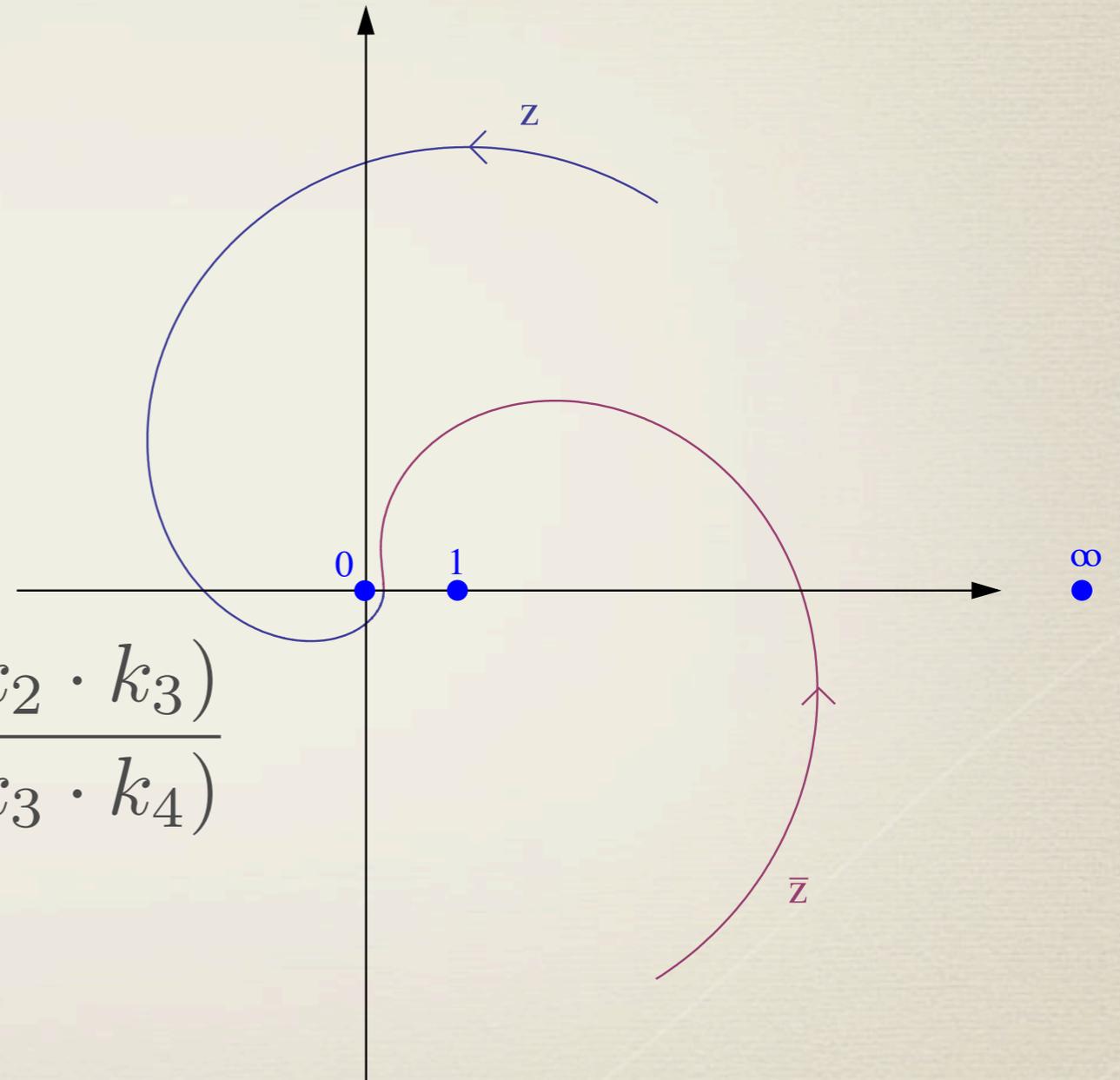
Analytic Continuation

* Singularity only present in physical Lorentzian continuation.

$$* z\bar{z} = \frac{(k_1 \cdot k_3)(k_2 \cdot k_4)}{(k_1 \cdot k_2)(k_3 \cdot k_4)}$$

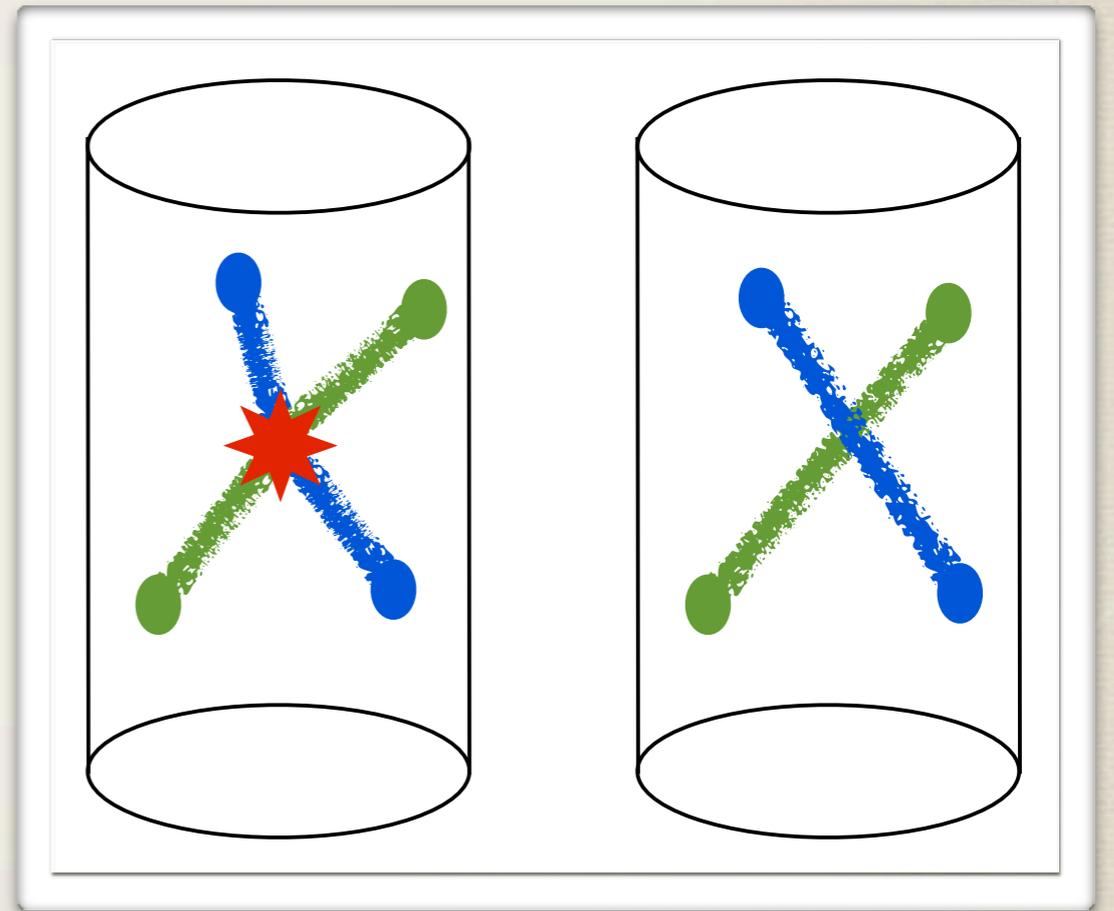
$$(1 - z)(1 - \bar{z}) = \frac{(k_1 \cdot k_4)(k_2 \cdot k_3)}{(k_1 \cdot k_2)(k_3 \cdot k_4)}$$

$$* z \rightarrow \bar{z}$$



Singularity Structure

- * Signal of interaction from intersecting wavepackets:
local bulk physics!
- * $z = \sigma e^{-\rho}$, $\bar{z} = \sigma e^{\rho}$
- * $\mathcal{A}(z, \bar{z}) \approx g^2 R^{5-d-2j} \frac{F(\sigma)}{(-\rho^2)^\beta}$
- * $\beta = \Delta_1 + \Delta_2 + j - 5/2$



Momentum Conservation

- * Flat space S-Matrix conserves momentum:

$$S = 1 + i(2\pi)^D \delta^D \left(\sum k_i \right) \mathcal{T}(s, t)$$

- * Momentum conserving δ -function must emerge from CFT amplitude in appropriate limit.

- * Delta function does emerge from form of boundary compact sources and singularity structure:

$$\lim_{R \rightarrow \infty} \int d\nu \frac{R^n e^{-i\nu}}{(R^2 \kappa^2 - (\nu + i\epsilon)^2)^\beta} \propto \delta^n(\vec{\kappa})$$

The S-Matrix

- * Determine flat space scattering amplitude from residue of CFT singularity:

$$\mathcal{A}(z, \bar{z}) \rightarrow g^2 R^{5-d-2j} \frac{\mathcal{F}(\sigma)}{(-\rho^2)^\beta}$$

$$iT(s, t) = \mathcal{K} g^2 s^{j-1} \left(\frac{-t}{s} \right)^{j-2} \left(\frac{-u}{s} \right)^{3-j-\Delta_1-\Delta_2} \mathcal{F} \left(\frac{-s}{t} \right)$$

Examples

* Compute CFT correlators using AdS/CFT dictionary (D'Hoker *et al.*, 1999), read off scattering amplitudes.

* Scalar exchange:

$$\mathcal{F}(\sigma) \propto \sigma(1 - \sigma)^{\Delta_1 + \Delta_2 - 3} \quad \rightarrow \quad \mathcal{T}(s, t) = \frac{g^2}{-t}$$

* Graviton exchange:

$$\mathcal{F}(\sigma) \propto \frac{(1 - \sigma)^8}{\sigma} \quad \rightarrow \quad \mathcal{T}(s, t) = 8\pi G_5 \frac{s^2 + ts}{-t}$$

Limits on Resolution

- * Well behaved wavepackets are critical for derivation of LSZ decomposition.
- * Bulk-Boundary propagator maps boundary-compact sources to non-compact wavepackets in the bulk.
- * Cannot construct regular wavepackets, typically used for formal derivations of LSZ decomposition.
- * What about Schwartz wavepackets?

Power-Law Tails

* *NO*: boundary-compact wavepackets have power-law tails in flat region.

* $x_{\perp} \gg \Delta t / \Delta \theta, u / \Delta \theta$:

$$\phi_f(x) \approx \phi_f(0) \frac{\omega \Delta t \tilde{L}(\omega \Delta t) \hat{L}}{(x_{\perp} \omega \Delta \theta)^{\Delta}}$$

* $u \gg \Delta t$

$$\phi_f(x) \approx \phi_f(0) \frac{\omega \Delta t \tilde{L}(\omega \Delta t) \Gamma(\Delta)}{2\pi (i\omega u)^{\Delta}}$$

Recovering the S-Matrix?

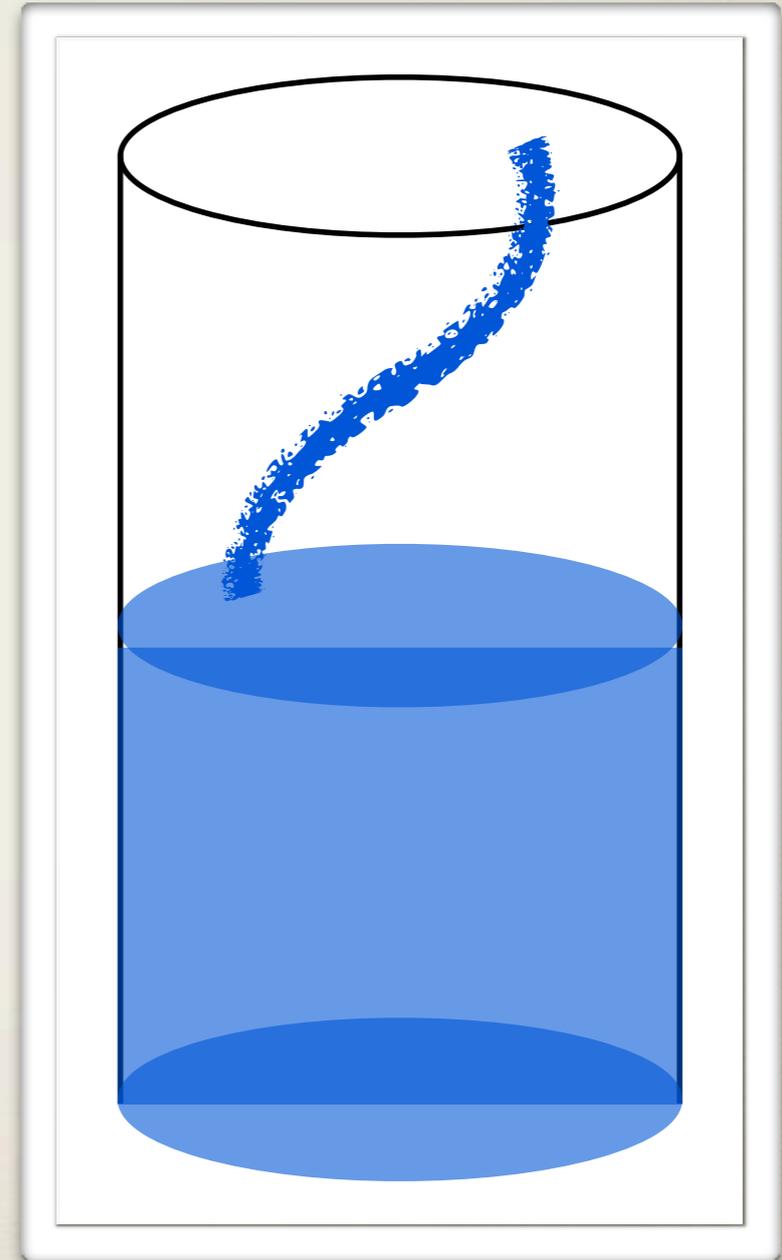
- * Tails from direct contribution interfere with scattered contribution to amplitude.
- * Scattered contribution doesn't always go in the correct direction: haze of order $\tilde{L}(\omega\Delta t)$.
- * We cannot recover the full flat space S-Matrix this way.
- * Can we build better wavepackets?

Resonant Wavepackets

- * Use resonant structure of AdS to build any normalizable wavefunction.

$$\phi_f(x) = \sum_{nl\vec{m}} c_{nl\vec{m}} \phi_{nl\vec{m}}(x)$$

- * Source compactly supported for $1/2$ AdS time.
- * Well-behaved in interaction region.
- * What about while the source is turned on?



Multiple Interactions and Power Law Tails

* While source is on, two terms contribute in the large R limit.

* First term: wavepacket builds/decays linearly in time,

$$\phi_f(x) \left(1 + \frac{\tau}{\pi}\right)$$

→secondary interactions while wavepackets built/decay.

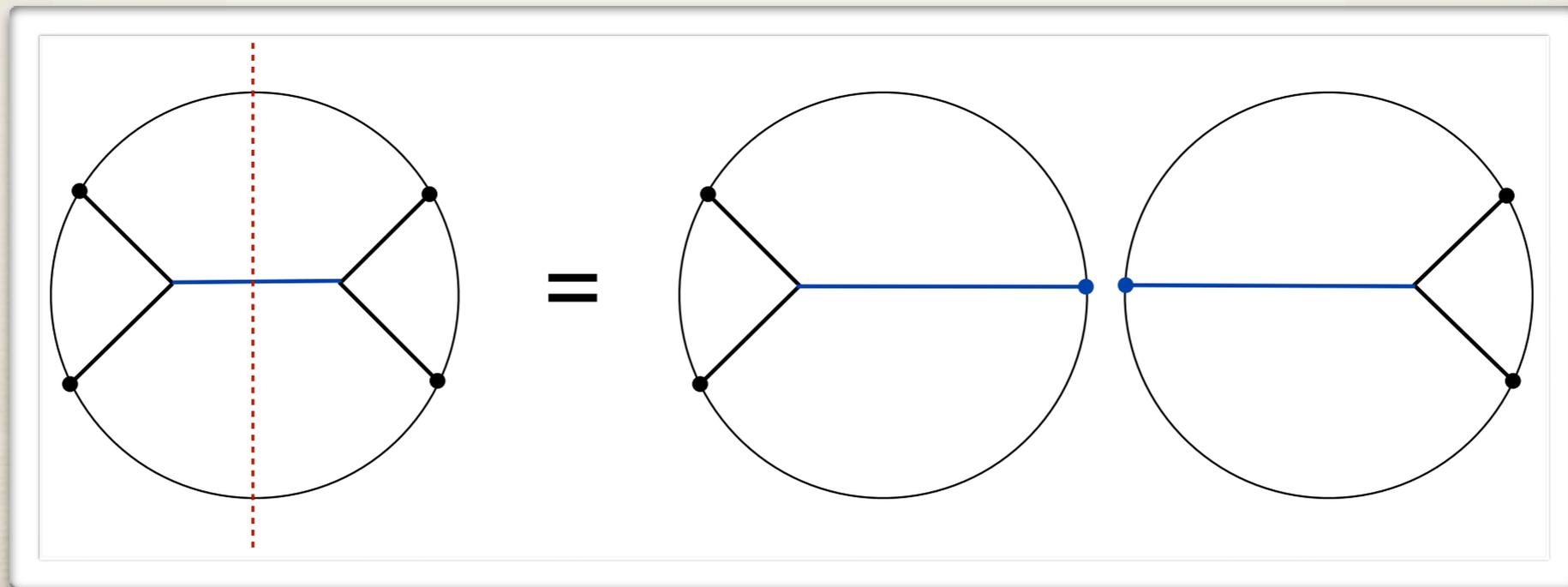
* Second term: power law tail in time.

$$\int d\omega \phi_f(\omega) \frac{\omega^{\frac{d-1}{2}} e^{-i\omega R\tau}}{(-i\omega R(\pi + \tau))^{\Delta}}$$

* Not quite a No-Go theorem, but hinders S-Matrix.

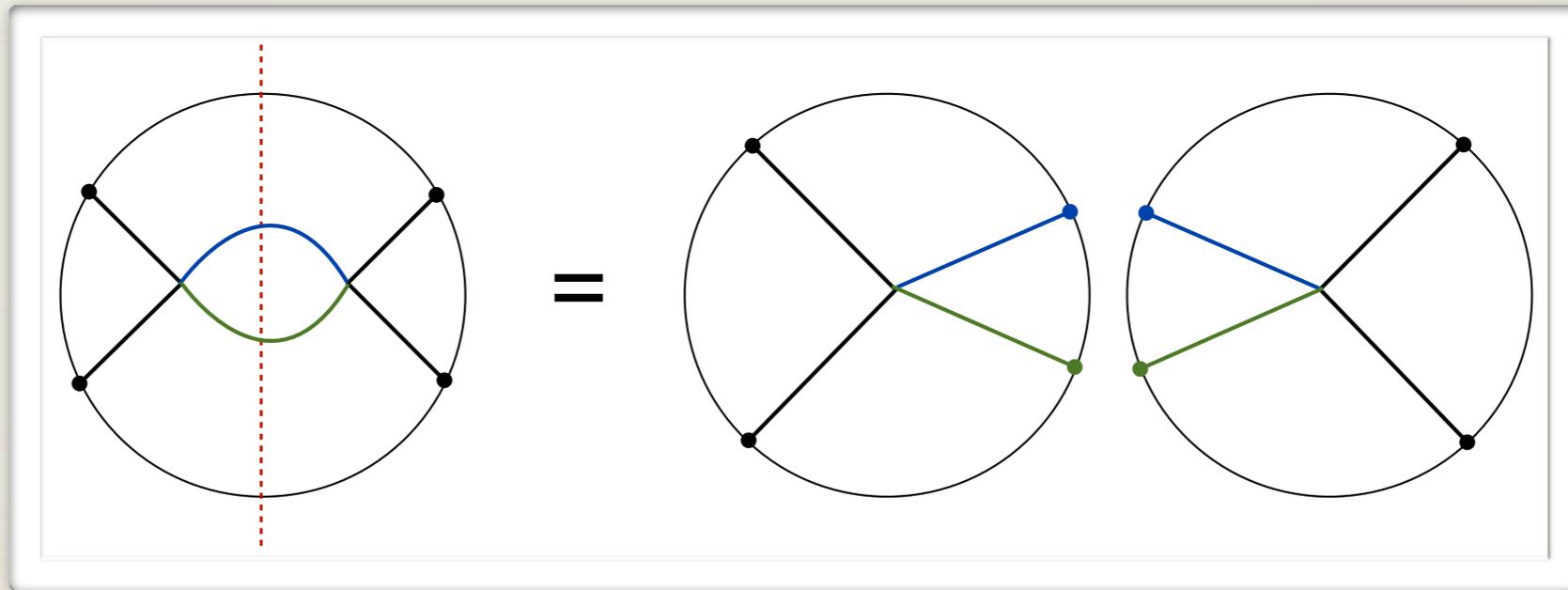
Bulk Unitarity

- * Often claimed: Unitary evolution in the CFT \Rightarrow unitary evolution in the bulk.
- * Can we see bulk unitarity from the CFT?
- * Look at perturbative bulk unitarity. At tree level, just the OPE:



Bulk Unitarity

- * What about loops?
- * Seem to need new relations:



$$\langle \mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_3 \mathcal{O}_4 \rangle_{1\text{loop}} = \sum_{A, B \neq 1} \langle \mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_A \mathcal{O}_B \rangle_{\text{tree}} \langle \mathcal{O}_A \mathcal{O}_B \mathcal{O}_3 \mathcal{O}_4 \rangle_{\text{tree}}$$

- * Are these relations present? What does this imply?

Conclusions

- * CFT singularity signature of fine-grained locality.
- * Boundary-compact wavepackets are not as well-localized as flat space wavepackets.
- * Power-law tails can hide important physics.
 - * Signature of black hole formation is exponentially suppressed $2 \rightarrow 2$ amplitude, swamped by tails.
- * Can build better wavepackets in AdS from compact sources on the boundary, but have multiple scattering.

Unresolved Questions

- * Is there a way to build nice wavepackets that don't have multiple scattering or power-law tails?
- * Are power-law tails a signature of inherent non-locality?
 - * Does the CFT really capture everything in the gravitational theory?
- * How is bulk unitarity encoded in the boundary CFT?
 - * Necessary for understanding black hole evaporation & resolving information problem.

Thank You