# From Fuzzy Physics to ncQFT

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#### Introduction

- Fuzzy Physics John Madore
- RG flows
- Scalar fields, modified action H G + R Wulkenhaar
- Taming the Landau ghost H G + R Wulkenhaar, V Rivasseau et al
- Fermions: spectral triple H G + R Wulkenhaar
- Further transl. inv. renormalizable models H G + F Vignes-Tourneret
- Sketch of proof
- Conclusions







# to John Madore Congratulations to your 70th birthday



## **Fuzzy Physics**

QFT needs regularization, example D = 2, euclidean John Madore: Fuzzy Sphere matrix algebra  $A_N$  generators:  $x_n$ 

$$[x_n^N, x_m^N] = i \frac{R}{\sqrt{N^2 - 1}} \epsilon_{nmp} x_p^N$$

gives embedding of algebra  $A_N$  into  $A_{N+1}$ ..... $A_{comm}$  regularization of QFT

$$\langle F \rangle = \frac{1}{N} \int [d\Phi]_N e^{-S_N[\Phi]} F(\Phi)$$

$$S_N[\Phi] = \frac{1}{N} Tr([x_m^N, \Phi][x_m^N, \Phi] + V(\Phi))$$

cutoff  $\frac{R}{N}$ , ex:  $CP^2$ ....,shows IR/UV mixing



#### **RG Flow**

#### **Project**

merge general relativity with quantum field theory through noncommutative geometry, use RG flow

- 50.... success of ren. pert. th. BPHZ....Connes-Kreimer
- 56.... ghost and triviality Landau....
- 74 RG flow Wilson, Polchinski... summability, safety
- use multi scale analysis V Rivasseau et al
- cure problems of ren. pert. exp. (IR,UV,convergence)
- require (Borel) summability
- take into account qu. gravity effects





#### **Formulation**

 $\phi^4$  on nc  $\mathbb{R}^4$ ,  $[\mathbf{x}^{\mu}, \mathbf{x}^{\nu}] = i\theta^{\mu\nu}$  antisymmetric, or equivalently star product

$$(a*b)(x) = \int dy \int dka(x + \frac{\theta k}{2})b(x + y)e^{iky}$$

 $\phi^4$  action

$$S = \int dp(p^2 + m^2)\phi_p \phi_{-p} + \lambda \int \prod_{j=1}^4 \left(dp_j \phi_{p_j}\right) \delta(\sum_{j=1}^4 p_j) e^{-i\sum_{i < j} p_i \wedge p_j}$$

#### Feynman rules

cyclic order of momenta leads to ribbon graphs Model is not renormalizable

One possible solution: modify action



#### Theorem

H. G. and R. Wulkenhaar  $\phi^4$  model modified, IR/UV mixing: short and long distances related Theorem: Action

$$S = \int d^4x \Big( \frac{1}{2} \partial_{\mu} \phi \star \partial^{\mu} \phi + \frac{\Omega^2}{2} (\tilde{\mathbf{x}}_{\mu} \phi) \star (\tilde{\mathbf{x}}^{\mu} \phi) + \frac{\mu_0^2}{2} \phi \star \phi + \lambda \phi \star \phi \star \phi \star \phi \Big) (\mathbf{x})$$

for 
$$ilde{ ilde{x}}_{\mu}:=2( heta^{-1})_{\mu
u}\, ilde{x}^{
u}$$

is perturbativly renormalizable to all orders in  $\lambda$ , 3 proofs, Rivasseau et al: Multiscale analysis in matrix base and in position space

Action has Langmann-Szabo position-momentum duality  $S[\phi; \mu_0, \lambda, \Omega] \mapsto \Omega^2 S[\phi; \frac{\mu_0}{\Omega}, \frac{\lambda}{\Omega^2}, \frac{1}{\Omega}]$ 

## $\beta$ function

evaluate  $\beta$  function, H. G. and R. Wulkenhaar,

$$\beta_{\lambda} = \frac{\lambda_{\text{phys}}^2}{48\pi^2} \frac{(1-\Omega_{\text{phys}}^2)^3}{(1+\Omega_{\text{phys}}^2)^3} + \mathcal{O}(\lambda_{\text{phys}}^3)$$
 flow bounded, L. ghost killed! 
$$\begin{array}{c} \lambda_{[\Lambda]} \\ 0.8 \\ \end{array}$$
 Due to wave fct. renormalization 
$$\begin{array}{c} \lambda_{[\Lambda]} \\ \Omega = 1 \text{ betafunction vanishes} \\ \Omega^2[\Lambda] \leq 1 \\ (\lambda[\Lambda] \text{ diverges in comm. case}) \\ \end{array}$$

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- $\Omega^2[\Lambda] < 1$  $(\lambda[\Lambda])$  diverges in comm. case)
  - perturbation theory remains valid at all scales!
  - non-perturbative construction of the model seems possible!

Gurau, Magnen, Rivasseau, Tanasa new ren.m: add  $\int \Phi^2 \frac{\alpha}{n^2}$ 



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## A spectral triple

H. G. and Raimar Wulkenhaar, Take Dirac operator on Hilbert space  $L^2(R^4)\otimes C^{16}$ 

$$D_8 = (i\Gamma^{\mu}\partial_{\mu} + \Omega\Gamma^{\mu+4}\tilde{\mathbf{x}}_{\mu})$$

 $\mu=$  1, ...4,  $\Gamma_k$  generate 8-dim Clifford algebra  $\{\Gamma_k\Gamma_l\}=2\delta_{kl}$ 

$$D_8^2 = (-\Delta + \Omega^2 ||\tilde{x}||^2) 1 - i\Omega\Theta_{\mu\nu}^{-1}[\Gamma^{\mu}, \Gamma^{\nu+4}]$$

compute action of Dirac operator on sections of spinor bundle

$$[D_8, f] * \psi = i[\Gamma^{\mu} + \Omega \Gamma^{\mu+4}](\partial \mu f) * \psi$$

only 4 dim. differential appears leads to spectral triple configuration space dimension 4 phase space dim. 8 Clifford alg. dim., KO dim.,...

nc Gross-Neveu model ren. by F Vignes-Tourneret

## Translation invariant rncQFT

H G + F Vignes-Tourneret, action:

$$S[\Phi] = \int d^4x \Phi(x) (-\Delta + m^2) \Phi(x) + \frac{\mu}{\Theta^3} (\int d^4x \Phi(x))^2 + \int d^4x \lambda \Phi^{*4}(x)$$

Main result: This QFT is renormalizable to all orders of pertubation theory

Topology of Ribbon graphs: V vertices, I internal lines, F faces,

B broken faces: 2-2g=V-I+F, 2L+N=4n

g = 0 planar graph, g > 0 nonplanar,

B = 1 regular graph, B > 1 nonregular

Filk: phases for Rosette graphs:  $e^{-i\sum_{l < m} I_{l,m}p_l \wedge p_m}$ 

use multiscale analysis



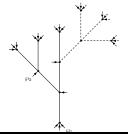
### Sketch of Proof

Slice propagator

$$C^k(p,q) = \delta(p+q) \int_{M^{-2k}}^{M^{-2(k-1)}} dt e^{-t(p^2+m^2)}$$

Derive power counting lemma:

The degree of convergence of a Feynman graph is given by N(G) - 4 if g=0, and N(G) + 4 for g>0, where N is nr of external points, choose tree



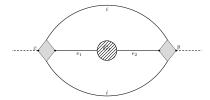
For non planar graphs use trick:

Go over a tree, there exist oscillations between momenta of lines crossing line I:

$$e^{-ip_l \wedge P_l} = \frac{1 - M^{2l} \partial_{p_l^2}^2}{1 + M^{2l} P_l^2} e^{-ip_l \wedge P_l}$$

do partial integration and use bound, gain factor  $M^{-8k}$  Irregular 4 pt fcts may diverge for B up to 4 finite number of counter terms

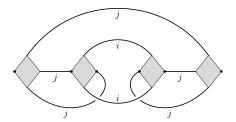
Ex: Irregular log. divergent 4-point subgraph



renormalization of 2 pt graph regulates 4 pt subdivergence,



If momenta of subgraph are neither external nor connected, it leads to nonplanar graphs, e g:



it has oscillations and is convergent

Planar regular and irregular 2 point fcts have to be renormalized Four point fct: Planar regular and some planar irregular graphs are divergent (appear only as subgraphs of planar two point graphs).



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## Renormalization of two and four point function

B = 1 perform Taylor expansion of propagator around zero external momenta

$$C^{i_k}(P+p) = C^{i_k}(p) + P\partial C^{i_k}(p) + \frac{1}{2}(P\partial)^2 C^{i_k}(p) + \dots$$

 $1^{st}$  term: quadratic mass ren.,  $2^{nd}$  term vanishes,  $3^{rd}$  term: wave fct ren. Two broken faces: Similar expansion of

$$A_G = \int d^4p \Phi(p) \Phi(-p) \int Propagators. Phases$$

Expand fields  $\Phi(p) = \Phi(0) + \int_0^1 ds p \nabla \Phi(sp)$  amplitude splits into 3 parts 1<sup>st</sup> term can be absorbed into nonlocal  $\mu$  term next terms treated by the trick...



### Conclusions

- Graphs couple internal and external momenta IR divergence not renormalizable
- modified actions for bosonic fields yield renorm. models
   Might give a nontrivial scalar Higgs model?
- RG flows save
- fermions give spectral triple
- further models: add nonlocal term
- generalization to gauge models?
- gravity?





## to John Many Happy

**Recurrencies** 

