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Monopoles and instantons on quantum projective spaces

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Part 1:

A spectral triple on the quantum projective space $\mathbb{C}P_q^2$
holomorphic calculus and spectral geometry

F. D'Andrea, L. Dabrowski, G. L.

The Noncommutative geometry of the quantum projective plane
Rev. Math. Phys. 20 (2008) 979–1006

Recently generalized to $\mathbb{C}P_q^N$

F. D'Andrea, L. Dabrowski

Dirac operators on quantum projective spaces
arXiv:0901.4735 [math.QA]

Part 2:

Monopoles connections on $\mathbb{C}P_q^2$

F. D'Andrea, G. L.

Anti-selfdual Connections on the Quantum Projective Plane:
Monopoles

arXiv:0903.3551v1 [math.QA]

F. D'Andrea, G. L.

Bounded and unbounded Fredholm modules for quantum projec-
tive spaces

arXiv:0903.3553v1 [math.QA]

A spectral triple for the quantum projective space $\mathbb{C}P_q^2$

$$\mathcal{A}(\mathbb{C}P_q^2), \mathcal{H}, D$$

a q-deformation of the classical principal fibrations:

$$U(2) \hookrightarrow SU(3) \rightarrow \mathbb{C}P^2$$

$$U(1) \hookrightarrow S^5 \rightarrow \mathbb{C}P^2$$

and q-deformation of associated bundles

With $0 < q < 1$, the function algebra

$$\mathcal{O} := \mathcal{A}(SU_q(3))$$

has generators u_j^i and relations

$$R_{kl}^{ij}(q) u_m^k u_n^l = u_m^k u_n^l R_{kl}^{ij}(q), \quad \sum_{p \in S_3} (-q)^{\|p\|} u_{p(1)}^1 u_{p(2)}^2 u_{p(3)}^3 = 1$$

$$(u_j^i)^* = (-q)^{j-i} (u_{l_1}^{k_1} u_{l_2}^{k_2} - q u_{l_2}^{k_1} u_{l_1}^{k_2})$$

$R_{kl}^{ij}(q)$ is the R -matrix of the $SU_q(n)$ series

Coproduct, counit and antipode:

$$\Delta(u_j^i) = \sum_k u_k^i \otimes u_j^k ,$$

$$\epsilon(u_j^i) = \delta_j^i ,$$

$$S(u_j^i) = (u_i^j)^*$$

The symmetry algebra

Symmetries are via the $*$ -Hopf algebra $\mathcal{U} := U_q(su(3))$

generated by $K_i, K_i^{-1}, E_i, F_i, i = 1, 2$ with $K_i = K_i^*, F_i = E_i^*$,
and relations ($a_{ij} = \text{Cartan matrix}$)

$$[K_i, K_j] = 0, \quad K_i E_j K_i^{-1} = q^{a_{ij}/2} E_j, \quad [E_i, F_j] = \delta_{ij} \frac{K_i^2 - K_i^{-2}}{q - q^{-1}}$$

$$E_i E_j^2 - (q + q^{-1}) E_i E_j E_i + E_j E_i^2 = 0 \quad \forall i \neq j.$$

Coproduct, counit and antipode are given by (with $i = 1, 2$)

$$\Delta(K_i) = K_i \otimes K_i, \quad \Delta(E_i) = E_i \otimes K_i + K_i^{-1} \otimes E_i,$$

$$\epsilon(K_i) = 1, \quad \epsilon(E_i) = 0,$$

$$S(K_i) = K_i^{-1}, \quad S(E_i) = -qE_i$$

Via a non-degenerate dual pairing $\langle , \rangle : \mathcal{U} \times \mathcal{O} \rightarrow \mathbb{C}$

define commuting left and right \mathcal{U} -actions on \mathcal{O} :

$$h \triangleright a = a_{(1)} \langle h, a_{(2)} \rangle , \quad a \triangleleft h = \langle h, a_{(1)} \rangle a_{(2)} ,$$

forall $h \in U_q(\mathfrak{su}(3))$, $a \in \mathcal{A}(SU_q(3))$

notation $\Delta(a) = a_{(1)} \otimes a_{(2)}$

The deformation of the quadratic Casimir of $U(\mathfrak{su}(3))$;
a central element in \mathcal{U} given by

$$\begin{aligned} \mathcal{C}_q := & (q - q^{-1})^{-2} \left((H + H^{-1}) \left\{ (qK_1K_2)^2 + (qK_1K_2)^{-2} \right\} + H^2 + H^{-2} - 6 \right) \\ & + (qHK_2^2 + q^{-1}H^{-1}K_2^{-2})F_1E_1 + (qH^{-1}K_1^2 + q^{-1}HK_1^{-2})F_2E_2 \\ & + qH[F_2, F_1]_q[E_1, E_2]_q + qH^{-1}[F_1, F_2]_q[E_2, E_1]_q \end{aligned}$$

with $H := (K_1K_2^{-1})^{2/3}$ $[a, b]_q := ab - q^{-1}ba$

The restriction of \mathcal{C}_q to the irrep (n_1, n_2) is

$$\mathcal{C}_q \Big|_{(n_1, n_2)} = \left[\frac{1}{3}(n_1 - n_2) \right]^2 + \left[\frac{1}{3}(2n_1 + n_2) + 1 \right]^2 + \left[\frac{1}{3}(n_1 + 2n_2) + 1 \right]^2$$

Some relevant subalgebras:

$U_q(su(2))$

the Hopf $*$ -subalgebra of $U_q(su(3))$ generated by $\{K_1, K_1^{-1}, E_1, F_1\}$

$U_q(u(2))$

the Hopf $*$ -subalgebra generated by $U_q(su(2))$ and $K_1K_2^2, (K_1K_2^2)^{-1}$

$K_1K_2^2$ commutes with all elements of $U_q(su(2))$

a class $\{\sigma_{\ell, N}\}$ of irreducible representation of $U_q(u(2))$ coming from a spin ℓ representation of $U_q(su(2))$, $\ell \in \frac{1}{2}\mathbb{N}$, and a representation of charge N on $K_1K_2^2$,

$$\sigma_{\ell, N}(K_1K_2^2) = q^N, \quad N \in \frac{1}{2}\mathbb{Z}$$

a constraint on the labels: $\ell + N \in \mathbb{Z}$

The quantum complex projective plane $\mathbb{C}P_q^2$

The $*$ -algebra $\mathcal{A}(SU_q(3))$ is an $U_q(\mathfrak{su}(3))$ -bimodule for the canonical actions \triangleright and \triangleleft

Call $\mathcal{A} := \mathcal{A}(\mathbb{C}P_q^2)$ the fixed point subalgebra for the **right** action of the $*$ -Hopf subalgebra $U_q(\mathfrak{u}(2)) \subset U_q(\mathfrak{su}(3))$:

$$\begin{aligned}\mathcal{A}(\mathbb{C}P_q^2) &= \mathcal{A}(SU_q(3))^{U_q(\mathfrak{u}(2))} \\ &= \left\{ a \in \mathcal{A}(SU_q(3)) \mid a \triangleleft h = \epsilon(h)a, \forall h \in U_q(\mathfrak{u}(2)) \right\}\end{aligned}$$

\mathcal{A} is a left $U_q(\mathfrak{su}(3))$ -module algebra;

\mathcal{A} generated by elements $p_{ij} := (u_i^3)^* u_j^3$, of a projection

$$p^2 = p = p^*$$

this projection is a line bundle over $\mathbb{C}P_q^2$ more later on

there are commutation rules (here $\text{sgn}(0) := 0$)

$$\begin{aligned}
 p_{ii}p_{jk} &= q^{\text{sgn}(i-j)+\text{sgn}(k-i)} p_{jk}p_{ii} && i, j, k \text{ distinct ,} \\
 p_{ii}p_{ij} &= q^{\text{sgn}(j-i)+1} p_{ij}p_{ii} - (1 - q^2) \sum_{k < i} q^{6-2k} p_{kk}p_{ij} && i \neq j , \\
 p_{ij}p_{ik} &= q^{\text{sgn}(k-j)} p_{ik}p_{ij} && i \notin \{j, k\} , \\
 p_{ij}p_{jk} &= q^{\text{sgn}(i-j)+\text{sgn}(k-j)+1} p_{jk}p_{ij} - (1 - q^2) \sum_{l < j} p_{il}p_{lk} && i, j, k \text{ distinct ,} \\
 p_{ij}p_{ji} &= (1 - q^2) \left(\sum_{l < i} p_{jl}p_{lj} - \sum_{l < j} p_{il}p_{li} \right) && i \neq j ,
 \end{aligned}$$

and 'projective plane' conditions

$$\sum_k p_{jk} p_{kl} = p_{jl} , \quad \text{Tr}_q(p) := q^4 p_{11} + q^2 p_{22} + p_{33} = 1$$

also: *-structure $(p_{ij})^* = p_{ji}$

There is a quantum sphere S_q^5

Call $\mathcal{B} := \mathcal{A}(S_q^5)$ the fixed point subalgebra for the **right** action of the $*$ -Hopf subalgebra $U_q(\mathfrak{su}(2)) \subset U_q(\mathfrak{su}(3))$:

$$\begin{aligned}\mathcal{A}(S_q^5) &= \mathcal{A}(SU_q(3))^{U_q(\mathfrak{u}(2))} \\ &= \left\{ a \in \mathcal{A}(SU_q(3)) \mid a \triangleleft h = \epsilon(h)a, \forall h \in U_q(\mathfrak{su}(2)) \right\}\end{aligned}$$

\mathcal{B} is a left \mathcal{U} -module algebra, and is generated by $z_i := u_i^3$

Also:
$$\mathcal{A}(\mathbb{C}P_q^2) \simeq \left\{ a \in \mathcal{A}(S_q^5) \mid a \triangleleft K_1 K_2^2 = a \right\}$$

and its generators are written as $p_{ij} := z_i^* z_j$

Relations for the sphere generators:

$$z_i z_j = q z_j z_i \quad \forall i < j, \quad z_i^* z_j = q z_j z_i^* \quad \forall i \neq j,$$

$$[z_1^*, z_1] = 0, \quad [z_2^*, z_2] = (1 - q^2) z_1 z_1^*,$$

$$[z_3^*, z_3] = (1 - q^2)(z_1 z_1^* + z_2 z_2^*),$$

$$z_1 z_1^* + z_2 z_2^* + z_3 z_3^* = 1.$$

The classical spin^c -structure

$\mathbb{C}P^2$ does not admit a spin structure; only spin^c structures

it is a Kähler manifold \Rightarrow

the bundle of antiholomorphic forms $\Omega^{0,\bullet}$ with a natural \mathbb{Z}_2 -grading is a canonical spin^c -bundle and a spin^c -Dirac operator is given by the Dolbeault-Dirac operator

$$D = \partial + \bar{\partial}$$

tensoring with line bundles one gets other spin^c -bundles

Associated bundles: line bundles in particular

$\sigma : U_q(u(2)) \rightarrow \text{End}(\mathbb{C}^n)$ an n -dimensional $*$ -representation

$\mathcal{A}(\mathbb{CP}_q^2)$ -bimodule of equivariant elements associated to σ :

$$\mathfrak{M}(\sigma) := \left\{ v \in \mathcal{A}(SU_q(3))^n \mid \sigma(S(h_{(1)})) \cdot (v \triangleleft h_{(2)}) = \epsilon(h)v, \forall h \in U_q(u(2)) \right\}$$

These are finitely generated projective right- $\mathcal{A}(\mathbb{CP}_q^2)$ modules

$\sigma = \sigma_{\ell, N}$, with $\ell \in \frac{1}{2}\mathbb{N}$, $N \in \frac{1}{2}\mathbb{Z}$, and condition $\ell + N \in \mathbb{Z}$

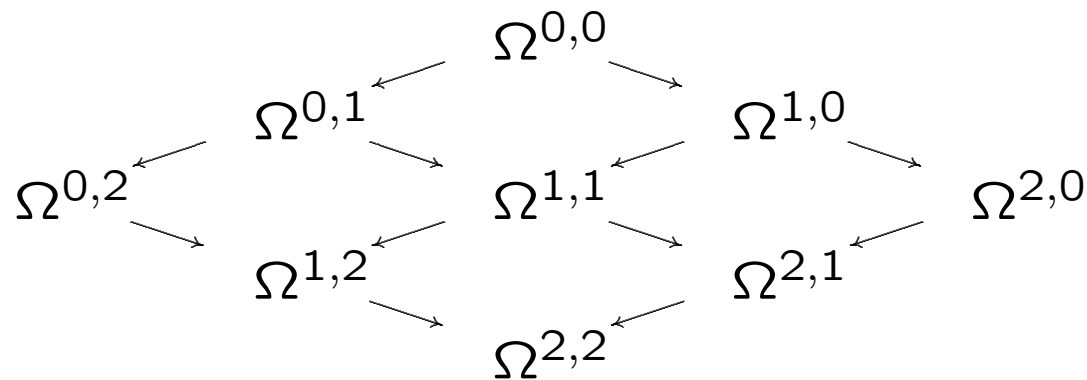
denote $\Sigma_{\ell, N} := \mathfrak{M}(\sigma_{\ell, N})$; in particular $\Sigma_{0,0} = \mathcal{A}$

Antiholomorphic forms again

$$\Omega^{0,0} := \mathcal{A} = \Sigma_{0,0}, \quad \Omega^{0,1} := \Sigma_{\frac{1}{2},\frac{3}{2}}, \quad \Omega^{0,2} := \Sigma_{0,3},$$

The full differential calculus

$$\Omega^{\bullet,\bullet} = \bigoplus \Omega^{i,j}$$



$\Omega^{i,j} \simeq \Sigma_{\ell,N}$, with suitable values of (ℓ, N)

The double complex: $\partial : \Omega^{p,q} \rightarrow \Omega^{p+1,q}$ and $\bar{\partial} : \Omega^{p,q} \rightarrow \Omega^{p,q+1}$

denote $\partial_{i,j} := \partial|_{\Omega^{i,j}}$ and $\bar{\partial}_{i,j} := \bar{\partial}|_{\Omega^{i,j}}$ and $\mathcal{R}_h a := a \triangleleft h$

$$\partial_{0,0} := \begin{pmatrix} \mathcal{R}_{E_2} \\ \mathcal{R}_Y \end{pmatrix},$$

$$\bar{\partial}_{0,0} := \begin{pmatrix} \mathcal{R}_X \\ \mathcal{R}_{F_2} \end{pmatrix},$$

$$\partial_{1,0} := - \begin{pmatrix} \mathcal{R}_{X^*} & \mathcal{R}_{E_2} \end{pmatrix},$$

$$\bar{\partial}_{0,1} := \begin{pmatrix} \mathcal{R}_{F_2} & \mathcal{R}_{Y^*} \end{pmatrix},$$

$$\partial_{0,1} := - \begin{pmatrix} -\mathcal{R}_Y & \mathcal{R}_{E_2} \\ \mathcal{R}_{E_2} & 0 \\ \mathcal{R}_Y & \mathcal{R}_{E_2} \\ 0 & \mathcal{R}_Y \end{pmatrix},$$

$$\bar{\partial}_{1,0} := \begin{pmatrix} \mathcal{R}_{F_2} & -\mathcal{R}_X \\ \mathcal{R}_X & 0 \\ \mathcal{R}_{F_2} & \mathcal{R}_X \\ 0 & \mathcal{R}_{F_2} \end{pmatrix},$$

$$\partial_{1,1} := \begin{pmatrix} -\mathcal{R}_{E_2} & \mathcal{R}_{X^*} & \mathcal{R}_{E_2} & 0 \\ \mathcal{R}_{X^*} & 0 & \mathcal{R}_{X^*} & \mathcal{R}_{E_2} \end{pmatrix}, \quad \partial_{0,2} := \begin{pmatrix} \mathcal{R}_{E_2} \\ \mathcal{R}_Y \end{pmatrix},$$

$$\bar{\partial}_{1,1} := \begin{pmatrix} \mathcal{R}_{Y^*} & \mathcal{R}_{F_2} & \mathcal{R}_{Y^*} & 0 \\ -\mathcal{R}_{F_2} & 0 & \mathcal{R}_{F_2} & \mathcal{R}_{Y^*} \end{pmatrix}, \quad \bar{\partial}_{2,0} := \begin{pmatrix} \mathcal{R}_X \\ \mathcal{R}_{F_2} \end{pmatrix},$$

$$\partial_{1,2} := - \begin{pmatrix} \mathcal{R}_{X^*} & \mathcal{R}_{E_2} \end{pmatrix},$$

$$\bar{\partial}_{2,1} := \begin{pmatrix} \mathcal{R}_{F_2} & \mathcal{R}_{Y^*} \end{pmatrix}$$

One finds: $\partial^2 = \bar{\partial}^2 = \partial\bar{\partial} + \bar{\partial}\partial = 0$; $d = \partial + \bar{\partial}$; $d^2 = 0$

An inner product on $\Omega^{\bullet,\bullet}$: $\langle \omega, \eta \rangle := \sum_{p,q} \varphi(\omega_{p,q}^\dagger \eta_{p,q})$

with φ the Haar functional of $\mathcal{A}(SU_q(3))$

$\Omega^{2,2}$ is a rank one free module with basis a central element vol

use it to define an integral

$$\int \omega := \langle \text{vol}, \omega \rangle = \varphi(\omega_{2,2}), \quad \omega \in \Omega^{\bullet,\bullet}$$

since ∂ and $\bar{\partial}$ are constructed with the right action of elements in $\ker \epsilon$ and the Haar functional is invariant ($\varphi(a \triangleleft x) = \epsilon(x)\varphi(a)$), the integral is closed

$$\int \bar{\partial}\omega = \int \partial\omega = 0$$

A spectral triple over $\mathbb{C}P_q^2$ from antiholomorphic forms

$$\Omega^{(0,0)} \xrightarrow{\bar{\partial}} \Omega^{(0,1)} \xrightarrow{\bar{\partial}} \Omega^{(0,2)} \rightarrow 0 .$$

\mathcal{H}_+ the completion of $\Omega^{(0,0)} \oplus \Omega^{(0,2)}$, \mathcal{H}_- of $\Omega^{(0,1)}$,

$$\mathcal{H} := \mathcal{H}_+ \oplus \mathcal{H}_-$$

Dolbeault-Dirac operator:

$$D\omega := (\bar{\partial}^\dagger v, \bar{\partial}a + \bar{\partial}^\dagger b, \bar{\partial}v) , \quad \omega = (a, v, b) \in \Omega^{(0,\bullet)}$$

$\bar{\partial}^\dagger$ the Hermitian conjugate of $\bar{\partial}$

$$[D, f] \in \mathcal{B}(\mathcal{H}) \text{ for all } f \in \mathcal{A}$$

A self-adjoint extension of D defined once it is diagonalized
 Compactness of $(D + i)^{-1}$ from asymptotic behaviour of $\text{Sp}(D)$

$\ker D = \mathbb{C}$ are the constant 0-forms;
 non-zero eigenvalues of D are ($n \geq 1$)

$$\pm \sqrt{\frac{2}{[2]}[n][n+2]} \quad \text{with multiplicity } (n+1)^3,$$

$$\pm \sqrt{[n+1][n+2]} \quad \text{with multiplicity } \frac{1}{2}n(2n+3)(n+3)$$

The spectrum of D is a q -deformation of the spectrum of the
 Dolbeault-Dirac operator of $\mathbb{C}P^2$

$\text{Sp}(D)$ grows exponentially: a 0^+ -dimensional spectral triple

The spectrum of the Dirac operator from

$$D^2 \omega = [2]^{-1} \omega \triangleleft (\mathcal{C}_q - 2) \quad \mathcal{C}_q \quad \text{is the quadratic Casimir}$$

The Hodge star operator

the linear operator $*_H : \Omega^{i,j} \rightarrow \Omega^{2-j,2-i}$

$$\int \omega^* \wedge_q \omega' = \langle *_H \omega, \omega' \rangle$$

Since $\Omega^{2,2} \simeq \mathcal{A}$ and the Haar state is faithful, equivalently

$$\omega^* \wedge_q \omega' = \langle *_H \omega, \omega' \rangle \text{vol}$$

The calculus defined so that $*_H^2 \omega = (-1)^{\text{dg}(\omega)} \omega$

$$d^\dagger = *_H d *_H .$$

$\epsilon(\omega)$ the left 'exterior product' : $\epsilon(\omega)\omega' := \omega \wedge_q \omega'$

$i(\omega) = \epsilon(\omega)^\dagger$ the 'contraction' by ω

$$*_H \omega = i(\omega^*) \text{vol}$$

Line bundles

$\Sigma_{0,N}$ are line bundles of 'degree' N

As right \mathcal{A} -modules:

$$\Sigma_{0,N} \simeq P_N \mathcal{A}^{rN}, \quad P_N = \Psi_N \Psi_N^\dagger$$

Ψ_N is the column vector with components $\psi_{j,k,l}^N$:

$$\begin{aligned} (\psi_{j,k,l}^N)^* &:= \sqrt{[j, k, l]!} z_1^j z_2^k z_3^l, & \text{if } N \geq 0 ; j + k + l = N, \\ (\psi_{j,k,l}^N)^* &:= q^{-N+j-l} \sqrt{[j, k, l]!} (z_1^j z_2^k z_3^l)^*, & \text{if } N < 0 ; j + k + l = -N \end{aligned}$$

$$\Psi_N^\dagger \Psi_N = \mathbf{1} \quad \Rightarrow \quad (P_N)^2 = P_N$$

The size is $r_N := \frac{1}{2}(N + 1)(N + 2)$; think of Ψ_N as a column vector of size r_N , and of Ψ_N^\dagger as a row vector of the same size.

$[j, k, l]!$ are the q -trinomial coefficients:

$$[j, k, l]! = q^{-(jk+kl+lj)} \frac{[j + k + l]!}{[j]![k]![l]!}$$

q -factorial is

$$[n]! := [n][n - 1] \dots [2][1], \quad n > 0; \quad [0]! := 1$$

The Grassmann connection:

$$\nabla : \Sigma_{0,N} \otimes_{\mathcal{A}} \Omega \rightarrow \Sigma_{0,N} \otimes_{\mathcal{A}} \Omega \quad \nabla := P_N \circ d$$

has curvature which is constant:

$$\nabla_N^2 = q^{N-1} [N] \nabla_1^2 \in \Omega^{(1,1)}$$

and anti-self-dual:

$$* \nabla_N^2 = -\nabla_N^2$$

the bundles are of rank 1

'first Chern number' N

'first Chern number' $\frac{1}{2}N(N+1)$

Gauged Laplacian operator

$$\square_N = \nabla_N^\dagger \nabla_N$$

Related to the Casimir \mathcal{C}_q

$$\square_N = q^{-\frac{3}{2}} \frac{q^{\frac{3}{2}} + q^{-\frac{3}{2}}}{q^{\frac{N}{3}} + q^{-\frac{N}{3}}} \left(\mathcal{C}_q - \left[\frac{1}{3}N \right]_q^2 - \left[\frac{1}{3}N + 1 \right]_q^2 - \left[\frac{2}{3}N + 1 \right]_q^2 \right) + [2]_q [N]_q ,$$

The spectrum $\{\lambda_{n,N}\}_{n \in \mathbb{N}}$ of \square_N

$$\lambda_{n,N} = (1 + q^{-3})[n]_q[n + N + 2]_q + [2]_q[N]_q \quad \text{if } N \geq 0 ,$$

$$\lambda_{n,N} = (1 + q^{-3})[n + 2]_q[n - N]_q + [2]_q[N]_q \quad \text{if } N \leq 0 .$$

with $n \in \mathbb{N}$.

not invariant under the exchange $N \leftrightarrow -N$,

not even when sending $q \leftrightarrow q^{-1}$

K-theory and K-homology

the C^* -algebra

$$0 \rightarrow \mathcal{K} \rightarrow C(\mathbb{C}P_q^2) \rightarrow C(\mathbb{C}P_q^1) \rightarrow 0$$

$$C(\mathbb{C}P_q^1) \simeq C(S_q^2) = \mathcal{K} \oplus \mathbb{C}1$$

$$K_0(C(\mathbb{C}P_q^2)) = \mathbb{Z}^3, \quad K_1(C(\mathbb{C}P_q^2)) = 0$$

$$K^0(C(\mathbb{C}P_q^2)) = \mathbb{Z}^3, \quad K^1(C(\mathbb{C}P_q^2)) = 0$$

There are interesting twisted cocycles;
in particular a twisted volume form

$K_0(\mathcal{A})$: the a. group of finitely generated projective \mathcal{A} -module

$K^0(\mathcal{A})$: the a. group of even Fredholm modules

their pairing is via Chern characters; these are cyclic (co)-cycle,

a cyclic $2n$ -cocycle: $\tau_n : \mathcal{A}^{2n+1} \rightarrow \mathbb{C}$ which is cyclic:

$$\tau_n(a_0, a_1, \dots, a_{2n}) = \tau_n(a_{2n}, a_0, \dots, a_{n-1}),$$

and Hochschild boundary closed: $b\tau_n = 0$

$$b\tau_n(a_0, \dots, a_{2n+1}) := \sum_{j=0}^{2n} (-1)^j \tau_n(a_0, \dots, a_j a_{j+1}, \dots, a_{2n+1}) \\ - \tau_n(a_{2n+1} a_0, a_1, \dots, a_{2n})$$

Even cyclic cocycles $\text{ch}_n^{(\pi, \mathcal{H}, F)}$, $2n \geq k$,
 canonically associated to a $k + 1$ -summable Fredholm module:

a triple (π, \mathcal{H}, F) :

$\mathcal{H} = \mathcal{H}_+ \oplus \mathcal{H}_-$ a \mathbb{Z}_2 -graded Hilbert space with grading γ

a graded representation $\pi : \mathcal{A} \rightarrow \mathcal{B}(\mathcal{H}_+) \oplus \mathcal{B}(\mathcal{H}_-)$

an odd operator F such that

$$[F, a_0][F, a_1] \dots [F, a_k] \quad \text{is traceclass}$$

then

$$\text{ch}_n^{(\pi, \mathcal{H}, F)}(a_0, \dots, a_{2n+1}) := \frac{1}{2}(-1)^n \text{Tr}_{\mathcal{H}}(\gamma F [F, a_0][F, a_1] \dots [F, a_{2n}])$$

The coupling with idempotents:

$$\langle , \rangle : K^0(\mathcal{A}) \times K_0(\mathcal{A}) \rightarrow \mathbb{Z} ,$$

$$\langle [(\pi, \mathcal{H}, F)], [e] \rangle = \frac{1}{2}(-1)^n \text{Tr}_{\mathcal{H} \otimes \mathbb{C}^m} (\gamma F [F, e]^{2n+1})$$

it is an integer ; the index of a Fredholm operator

Classical Invariants

Back to line bundles over $\mathbb{C}P_q^2$

The rank

The algebra $\mathcal{A} = \mathcal{A}(\mathbb{C}P_q^2)$ with generators p_{ij} has a character (1-dimensional representation):

$$\tau_0 : \mathcal{A} \rightarrow \mathbb{C}, \quad \tau_0(p_{ij}) := \delta_{i3}\delta_{j3}$$

then

$$\begin{aligned} \langle [\tau_0] | [P_N] \rangle &:= \tau_0(\text{ch}_0(P_N)) = \tau_0(\text{Tr}_{\mathbb{C}^{r_N}} P_N) \\ &= 1 \end{aligned}$$

The monopole charge via a Fredholm module on $\mathcal{A}(\mathbb{C}P_q^2)$

$\ell^2(\mathbb{N})$, with orthonormal basis $|n\rangle$;
the Hilbert space is $\mathcal{H}_1 := \ell^2(\mathbb{N}) \otimes \mathbb{C}^2$, and

$$F := \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \gamma := \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix},$$

representation $\pi := \pi_+ \oplus \pi_-$ and

$$\pi_{\pm}(p_{11}) = \pi_{\pm}(p_{12}) = \pi_{\pm}(p_{13}),$$

$$\pi_-(p_{22}) = \pi_-(p_{23}) = 0$$

$$\pi_+(p_{22})|n\rangle = q^{2n}|n\rangle, \quad \pi_+(p_{23})|n\rangle = q^{n+1}\sqrt{1 - q^{2(n+1)}}|n+1\rangle,$$

$\pi_+(p_{ij}) - \pi_-(p_{ij})$ is trace class for all i, j

the Fredholm module is 1-summable; a cyclic 0-cocycle:

$$\tau_1 : \mathcal{A}(\mathbb{C}\mathbb{P}_q^2) \rightarrow \mathbb{C}, \quad \tau_1(a) := \frac{1}{2} \text{Tr}_{\mathcal{H}}(\gamma F[F, a])$$

then

$$\begin{aligned} \langle [\tau_1] | [P_N] \rangle &:= \tau_1(\text{ch}_0(P_N)) = \frac{1}{2} \text{Tr}_{\mathcal{H} \otimes \mathbb{C}^{r_N}}(\gamma F[F, P_N]) \\ &= \dots \\ &= N \end{aligned}$$

The number depends only on the restriction of the bundle to the subspace $\mathbb{C}\mathbb{P}_q^1 \simeq S_q^2$, (a reason to call it the *monopole charge*)

The instanton charge via a Fredholm module on $\mathcal{A}(\mathbb{C}P_q^2)$

the Hilbert space \mathcal{H}_2 (is two copies of) the linear span of orthonormal vectors $|\ell, m\rangle$, with $\ell \in \frac{1}{2}\mathbb{N}$ and $\ell + m \in \mathbb{N}$

the grading γ and the operator F are like before.

for the representation $\pi_2 = \pi_+ \oplus \pi_-$

$$\begin{aligned} \pi_+(p_{11}) &= \pi_+(p_{12}) = \pi_+(p_{13}) = 0, \\ \pi_+(p_{22})|\ell, m\rangle &= \begin{cases} q^{2(\ell+m)}|\ell, m\rangle & \text{if } m \leq \ell, \\ 0 & \text{if } m > \ell, \end{cases} \\ \pi_+(p_{23})|\ell, m\rangle &= \begin{cases} q^{\ell+m+1}\sqrt{1 - q^{2(\ell+m+1)}}|\ell, m+1\rangle & \text{if } m \leq \ell - 1, \\ 0 & \text{if } m \geq \ell. \end{cases} \end{aligned}$$

$$\begin{aligned}
\pi_-(p_{11}) &= \pi_-(p_{12}) = \pi_-(p_{13}) = 0, \\
\pi_-(p_{22})|\ell, m\rangle &= q^{2(\ell+m)}|\ell, m\rangle, \\
\pi_-(p_{23})|\ell, m\rangle &= q^{\ell+m+1}\sqrt{1 - q^{2(\ell+m+1)}}|\ell, m + 1\rangle.
\end{aligned}$$

$\pi_+(p_{ij}) - \pi_-(p_{ij})$ is trace class for all i, j

the Fredholm module is 1-summable; a cyclic 0-cocycle:

$$\tau_2 : \mathcal{A}(\mathbb{C}P_q^2) \rightarrow \mathbb{C}, \quad \tau_2(a) := \frac{1}{2} \text{Tr}_{\mathcal{H}}(\gamma^F[F, a])$$

then

$$\begin{aligned}
\langle [\tau_2] | [P_N] \rangle &:= \tau_2(\text{ch}_0(P_N)) = \text{Tr}_{\mathcal{H}_2 \otimes \mathbb{C}^m}(\gamma^F[F, P_N]) \\
&= \dots \\
&= \frac{1}{2}N(N + 1)
\end{aligned}$$

For any $N \in \mathbb{Z}$, the (right) module $\Sigma_{0,N}$ has

'rank' 1

'monopole charge' N

'instanton number' $\frac{1}{2}N(N+1)$

The three generators $\{e_1, e_2, e_3\}$ of $K_0(\mathcal{A}(\mathbb{C}P_q^2))$ are:

$e_1 = [1]$ is the class of the rank one free $\mathcal{A}(\mathbb{C}P_q^2)$ -module $\Sigma_{0,0}$

e_2 is the class of $\Sigma_{0,-1}$ (the dual of the tautological bundle)

and e_3 is the class of $\Sigma_{0,1}$ (the tautological bundle)

The three generators of $K^0(\mathcal{A}(\mathbb{C}P_q^2))$ are the classes of the Fredholm modules $(\pi_i, \mathcal{H}_i, F_i)$, $i = 0, 1, 2$, given before

Quantum invariants

Classically, invariants of vector bundles are computed by integrating powers of the curvature of a connection on the bundle, the result being independent of the particular chosen connection.

In order to integrate the curvature of a connection on the quantum projective space $\mathbb{C}P_q^2$ one needs 'twisted integrals'; the result, is no longer an integer but rather its q -analogue

Invariants for $\mathbb{C}P_q^2$

The Haar state of $\mathcal{A}(SU_q(3))$

$$\varphi(ab) = \varphi((K \triangleright b \triangleleft K)a) , \quad \text{for } a, b \in \mathcal{A}(SU_q(3)) ,$$

that when $a, b \in \mathcal{A}(\mathbb{C}P_q^2)$ means

$$\varphi(ab) = \varphi((K \triangleright b)a) = \varphi(\eta(b)a) .$$

η -twisted cyclic $2n$ -cocycles: $\tau_n : \mathcal{A}^{2n+1} \rightarrow \mathbb{C}$

$$\tau_n(a_0, a_1, \dots, a_{2n}) = \tau_n(\eta(a_{2n}), a_0, \dots, a_{n-1}),$$

and Hochschild boundary closed: $b_\eta \tau_n = 0$

$$\begin{aligned} b_\eta \tau_n(a_0, \dots, a_{2n+1}) &:= \sum_{j=0}^{2n} (-1)^j \tau_n(a_0, \dots, a_j a_{j+1}, \dots, a_{2n+1}) \\ &\quad - \tau(\eta(a_{2n+1}) a_0, a_1, \dots, a_{2n}) \end{aligned}$$

The restriction of the Haar state to $\mathcal{A}(\mathbb{C}\mathbb{P}_q^2)$ is the representative of a class in the cohomology $[\tau_0] \in HC_\eta^0(\mathcal{A}(\mathbb{C}\mathbb{P}_q^2))$

An element $[\tau_4] \in HC_\eta^4(\mathcal{A}(\mathbb{C}\mathbb{P}_q^2))$ is constructed as

$$\tau_4(a_0, \dots, a_4) := \int a_0 da_1 \wedge_q \dots \wedge_q da_4 .$$

A 2-cocycle can be defined in a similar way.

Elements of $\Omega^{1,1}(\mathbb{C}\mathbb{P}_q^2)$ are $\omega = (\alpha, \alpha_4)$, with $\alpha_4 \in \mathcal{A}(\mathbb{C}\mathbb{P}_q^2)$.

Let $\pi : \Omega^{1,1}(\mathbb{C}\mathbb{P}_q^2) \rightarrow \mathcal{A}(\mathbb{C}\mathbb{P}_q^2)$ be such that $\pi(\omega) = \alpha_4$;

extend it to a projection $\pi : \Omega^2(\mathbb{C}\mathbb{P}_q^2) \rightarrow \mathcal{A}(\mathbb{C}\mathbb{P}_q^2)$

by $\pi(\omega) = 0$ if $\omega \in \Omega^{0,2}$ or $\omega \in \Omega^{2,0}$.

The map $\tau_2(a_0, a_1, a_2) := \varphi \circ \pi(a_0 da_1 \wedge_q da_2)$

is the representative of a class $[\tau_2] \in HC_\eta^2(\mathcal{A}(\mathbb{C}\mathbb{P}_q^2))$.

Both classes $[\tau_4]$ and $[\tau_2]$ are proven to be not trivial by pairing them with the monopole projections $P_N = \Psi_N \Psi_N^\dagger$

The pairing of $[\tau_4]$ with P_N :

$$\begin{aligned} \langle [\tau_4] | [P_N] \rangle &= \int \text{Tr} \left(P_N (dP_N)^4 \sigma^N (K_1^{-4} K_2^{-4})^t \right) \\ &= q^{-2N} \int \nabla_N^2 \wedge_q \nabla_N^2 \\ &\sim [N]^2 \end{aligned}$$

For $q = 1$, the integral of the square of the curvature is the *instanton number* of the bundle.

The pairing of $[\tau_2]$ with P_N :

$$\begin{aligned}\langle [\tau_2] | [P_N] \rangle &= \varphi \circ \text{Tr} \left(\pi(P_N dP_N \wedge_q dP_N) \sigma^N (K_1^{-4} K_2^{-4})^t \right) \\ &= q^{-2N} \varphi \circ \pi(\nabla_N^2) \\ &\sim [N]\end{aligned}$$

At $q = 1$ the integral of the curvature is the *monopole number* of the bundle.

The pairing of $[\tau_0]$ with P_N

$$\begin{aligned}\langle [\tau_0] | [P_N] \rangle &= \varphi \circ \text{Tr} \left(P_N \sigma^N (K_1^{-4} K_2^{-4})^t \right) \\ &= q^{-2N}\end{aligned}$$

At $q = 1$ this is the rank of the bundle.

If q is trascendental, this means that all $[P_N]$ are independent, i.e. the equivariant K_0 -group is infinite dimensional

Indeed, were the classes $[P_N]$ not independent, there would exist a sequence $\{k_N\}$ of integers – all zero but for finitely many – such that $\sum_N k_N q^{-2N} = 0$, and q^{-1} would be the root of a non-zero polynomial with integer coefficients.

thank you !!