## Convergence of Star Products and the Nuclear Weyl Algebra

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## Bayrischzell 2013

- Beiser, S., Waldmann, S.: Fréchet algebraic deformation quantization of the Poincaré disk. Preprint (August 2011), 58 pages. To appear in Crelle's J. reine angew. Math.
- Waldmann, S.: A Nuclear Weyl Algebra. Preprint 1209.5551 (September 2012), 48 pages.


## Plan of the talk

Introduction

The topology on $\mathrm{S}^{\bullet}(V)$
Continuity of $\star_{z \wedge}$
Further properties of the Weyl algebra
An example: the Peierls bracket and Free QFT
Outlook

## Introduction: Convergence of star products?

Main objective: Find examples and a framework where formal star products actually converge.
Needed for: reasonable theory of star exponentials, phase space reduction ...

The set-up (most simplest example):

- A vector space $V$ together with a bilinear form

$$
\Lambda: V \times V \longrightarrow \mathbb{R}
$$

determines Poisson structure on the symmetric algebra $S^{\bullet}(V)$.

- Formal Weyl star product gives formal deformation quantization of this.
- Alternatively: C*-algebraic Weyl algebra.

Idea: Want something in between: $\hbar$ not formal anymore, but not yet $C^{*}$, since interesting topology on $V$ should be remembered.

More precise set-up:

- V a locally convex topological vector space (Hausdorff), possibly continuously $\mathbb{Z}_{2}$-graded.
- $\Lambda$ a continuous bilinear form (not separately continuous). No symmetry properties assumed, but only pairing between even-even and odd-odd parts of $V$. Continuity means there is a continuous seminorm p with

$$
|\Lambda(v, w)| \leq \mathrm{p}(v) \mathrm{p}(w)
$$

- Define $P_{\wedge} \in \operatorname{End}\left(S^{\bullet}(V) \otimes S^{\bullet}(V)\right)$ by

$$
\begin{aligned}
& P_{\wedge}\left(v_{1} \cdots v_{n} \otimes w_{1} \cdots w_{m}\right) \\
& =\sum_{k, l} \Lambda\left(v_{k}, w_{l}\right) v_{1} \cdots \stackrel{k}{\wedge} \cdots v_{n} \otimes w_{1} \cdots{ }_{\wedge}^{\prime} \cdots w_{m}
\end{aligned}
$$

plus some signs if in graded case.

- The Poisson bracket on $S^{\bullet}(V)$

$$
\{a, b\}=\mu \circ\left(P_{\wedge}-\tau \circ P_{\wedge} \circ \tau\right)(a \otimes b)
$$

with $\mu$ the symmetric tensor product.

- The (formal) Weyl product

$$
a \star_{z \Lambda} b=\mu \circ \mathrm{e}^{z P_{\wedge}}(a \otimes b)
$$

for $z$ formal parameter or $z \in \mathbb{C}$. Here everything converges since we are on the symmetric algebra!

The topology on $\mathrm{S}^{\bullet}(V)$
Need a topology on $\mathrm{S}^{\bullet}(V)$ making the Poisson bracket and the star product continuous and allows for an interestingly large completion.

Work on the tensor algebra (easier) and induce the topology for $S^{\bullet}(V) \subseteq T^{\bullet}(V)$ from there.
On each tensor power $V^{\otimes n}$ use the $\pi$-topology: for each continuous seminorm p on $V$ define $\mathrm{p}^{n}=\mathrm{p} \otimes \cdots \otimes \mathrm{p}$ on $V^{\otimes n}$. Then all of these seminorms define the $\pi$-topology.
For $a=\sum_{n} a_{n} \in T^{\bullet}(V)$ define new seminorm

$$
\mathrm{p}_{R}(a)=\sum_{n=0}^{\infty} n!^{R} \mathrm{p}^{n}\left(a_{n}\right)
$$

and equip $T^{\bullet}(V)$ with all these seminorms, with parameter

$$
R \geq 0
$$

Gives a locally convex space denoted by $T_{R}^{\bullet}(V)$, inducing a locally convex topology on the symmetric algebra, denoted by $S_{R}^{\circ}\left(V_{)}\right)$.

- Completion of $T_{R}^{\bullet}(V)$ and $S_{R}^{\bullet}(V)$ can be described explicitly:

$$
\widehat{T}_{R}^{\bullet}(V)=\left\{a \in \prod_{n=0}^{\infty} V^{\hat{\otimes}_{\pi} n} \mid \forall \mathrm{p}: \mathrm{p}_{R}(a)<\infty\right\}
$$

- Elements in $\hat{\mathrm{S}}_{R}^{\bullet}(V)$ can be viewed as "analytic functions" on dual $V^{\prime}$ with certain growth conditions on Taylor coefficients, depending on the parameter $R$.
- Alternative versions of $\mathrm{p}_{R}$ using other $\ell^{p}$-summations or even a sup-version does not yield anything new.
- For $R=0$ the tensor algebra $\hat{T}_{R=0}^{\bullet}(V)$ is the free complete locally multiplicatively convex algebra generated by $V$.
- Unfortunately, there seems to be nothing like a free locally convex algebra generated by $V$ : otherwise we could just take a quotient by the canonical commutation relations and we would be done.
Note that we necessarily have to go beyond the case of sub-multiplicative seminorms!
- Simple examples show that we can not even expect an entire calculus for the canonical commutation relations: There are entire functions $f$ such that $f(q) f(p)$ does not make any sense for $q$ and $p$ satisfying $[q, p]=\mathrm{i} \hbar$.


## Continuity of $\star_{z \Lambda}$

Depending on the value of the parameter $R$, the product $\star_{z \Lambda}$ becomes continuous.

## Lemma

Let $R \geq \frac{1}{2}$. Then for all p satisfying $|\Lambda(v, w)| \leq \mathrm{p}(v) \mathrm{p}(w)$ there exists a constant $c>0$ with

$$
\mathrm{p}_{R}\left(a \star_{z \wedge} b\right) \leq c(c \mathrm{p})_{R}(a)(c \mathrm{p})_{R}(b)
$$

Proof: this is a longer but straightforward estimate. The constant $c$ depends on $R$ and on $z$.

## Definition

Let $R \geq \frac{1}{2}$. The symmetric algebra $\mathrm{S}_{R}^{\circ}(V)$ equipped with the Weyl product $\star_{z \Lambda}$ is called the locally convex Weyl algebra $\mathcal{W}_{R}\left(V, \star_{z \Lambda}\right)$.

Theorem
$\mathcal{W}_{R}\left(V, \star_{z \Lambda}\right)$ is a locally convex algebra. It is first countable, if $V$ is first countable.

Some first features of the construction:

- The star product $a \star_{z \Lambda} b$ for $a, b \in \mathcal{W}_{R}\left(V, \star_{z \Lambda}\right)$ converges absolutely and gives a entire deformation.
- For a real vector space $V$ and a complex-valued the complexified Weyl algebra $\mathcal{W}_{R}\left(V, \star_{\frac{1}{\hbar} 2 \Lambda}\right)$ with real $\hbar$ becomes a *-algebra with respect to complex conjugation if the symmetric part of $\Lambda$ is imaginary and the antisymmetric part is real.
- For $R<1$ the exponential series $\exp (v)$ for $v \in V$ belongs to the completion.

Further properties of the Weyl algebra
The construction of the locally convex Weyl algebra preserves many properties of the locally convex space $V$, quite like the $C^{*}$-algebraic construction.

## Theorem

If $V$ has an absolute Schauder basis, then $\mathcal{W}_{R}\left(V, \star_{z \Lambda}\right)$ has an absolute Schauder basis, too.
Recall that $\left\{e_{i}\right\}_{i \in I} \subseteq V$ with $\left\{\varphi^{i}\right\}_{i \in I} \subseteq V^{\prime}$ is called absolute Schauder basis if

$$
v=\sum_{i \in I} \varphi^{i}(v) e_{i}
$$

converges and for all continuous seminorms $p$

$$
v \mapsto \sum_{i \in I}\left|\varphi^{i}(v)\right| \mathrm{p}\left(e_{i}\right)
$$

is a continuous seminorm, too.

Theorem
The Weyl algebra $\mathcal{W}_{R}\left(V, \star_{z \Lambda}\right)$ is nuclear iff $V$ is nuclear.
Nuclearity has many equivalent definitions (technical) but the upshot is that nuclear spaces behave in many ways much nicer than general locally convex space, in particular concerning tensor products.

## Corollary

If $V$ is finite-dimensional, then the Weyl algebra $\mathcal{W}_{R}\left(V, \star_{z \Lambda}\right)$ is nuclear and it has an absolute Schauder basis.

- This gives nice features of the Weyl algebra concerning deformation quantization in finite dimensions, i.e. for the usual star products.
- The absolute Schauder basis can e.g. be obtained from the symmetric tensor powers of a basis of $V$, i.e. the "monomials".
- In this case the construction reproduces results of Omori, Maeda, Miyasaki and Yoshioka. Moreover, a slight variant reproduces the results of Beiser, Römer, W.

Theorem
The construction of $\mathcal{W}_{R}\left(V, \star_{z \Lambda}\right)$ depends functorially on the pair $(V, \Lambda)$.
In particular, the continuous linear Poisson maps $(V, \Lambda) \longrightarrow(\tilde{V}, \tilde{\Lambda})$ lift to continuous algebra homomorphisms.

Let $\varphi \in V^{\prime}$ be even. Then there is a unique algebra automorphism $\tau_{\varphi}$ of $\mathcal{W}_{R}\left(V, \star_{z \Lambda}\right)$ with

$$
\tau_{\varphi}(v)=v+\varphi(v) \mathbb{1}
$$

Geometrically, this is the pull-back of a polynomial by the translation by $\varphi$, if we interpret $S^{\bullet}(V)$ as polynomials on $V^{\prime}$.
Theorem
The automorphism $\tau_{\varphi}$ is continuous. If $R<1$ and if $\varphi$ is in the image of the musical homomorphism

$$
\sharp: V \ni w \mapsto w^{\sharp}=\Lambda_{-}(w, \cdot) \in V^{\prime}
$$

then $\tau_{\varphi}$ is inner.
The inner element is a $\star_{z \wedge}$-exponential series of a $w$ with $w^{\sharp}=\varphi$.

## Theorem

Suppose $\Lambda, \Lambda^{\prime}$ have the same antisymmetric part. Then the Weyl algebras $\mathcal{W}_{R}\left(V, \star_{z \Lambda}\right)$ and $\mathcal{W}_{R}\left(V, \star_{z \Lambda}\right)$ are isomorphic via a continuous isomorphism.
For formal star products this is nothing new, the emphasize lies on the word "continuous" here.

## An example: the Peierls bracket and Free QFT

As a first application in infinite dimensions: (Q)FT on globally hyperbolic space-times.

The set-up:

- Take a globally hyperbolic space-times $M \cong \mathbb{R} \times \Sigma$ with a normally hyperbolic differential operator $D=\square^{\nabla}+B=D^{*}$ on a real vector bundle $E \longrightarrow M$ with fiber metric $h$.
- Let $F_{M}=F_{M}^{+}-F_{M}^{-}$be the propagator where

$$
F_{M}^{ \pm}: \Gamma_{0}^{\infty}\left(E^{*}\right) \longrightarrow \Gamma_{\mathrm{sc}}^{\infty}\left(E^{*}\right)
$$

are the causal Green operators.

Canonical Poisson algebra: modelled on initial data on $\Sigma$

- Phase space will be $\Gamma_{0}^{\infty}\left(E_{\Sigma}\right) \oplus \Gamma_{0}^{\infty}\left(E_{\Sigma}\right)$.
- Take $V_{\Sigma}=\Gamma_{0}^{\infty}\left(E_{\Sigma}^{*}\right) \oplus \Gamma_{0}^{\infty}\left(E_{\Sigma}^{*}\right)$ as linear functions on the phase space (very few!)
- Take $S^{\bullet}\left(V_{\Sigma}\right)$ as polynomial algebra on the phases space.
- Endow this with the constant (canonical) Poisson structure coming from

$$
\Lambda_{\Sigma}\left(\left(\varphi_{0}, \dot{\varphi}_{0}\right),\left(\psi_{0}, \dot{\psi}_{0}\right)\right)=\int_{\Sigma}\left(h_{\Sigma}^{-1}\left(\varphi_{0}, \dot{\psi}_{0}\right)-h_{\Sigma}^{-1}\left(\dot{\varphi}_{0}, \psi_{0}\right)\right) \mu_{\Sigma}
$$

- This is a continuous bilinear form, so we can apply the Weyl algebra construction.

Now the covariant set-up:

- As covariant phase space we take $\Gamma_{\text {sc }}^{\infty}(E)$.
- The observables will then be polynomials on this, modeled by the symmetric algebra over $\Gamma_{0}^{\infty}\left(E^{*}\right)$.
- On $\Gamma_{0}^{\infty}\left(E^{*}\right)$ we have the bilinear form

$$
\Lambda(\varphi, \psi)=\int_{M} h^{-1}\left(F_{M}(\varphi), \psi\right) \mu
$$

- Again continuous.
- The corresponding Poisson bracket on $S^{\bullet}\left(\Gamma_{0}^{\infty}\left(E^{*}\right)\right)$ is the Peierls bracket.
- We can again apply the Weyl algebra construction for this.

The covariant Poisson bracket is now very much degenerate.

- The ideal generated by the kernel of $F_{M}$ turns out to be a Poisson ideal.
- This ideal coincides with those polynomials which vanish when evaluated on solutions $u \in \Gamma_{\text {sc }}^{\infty}(E)$ to the wave equation defined by $D$.


## Theorem

The quotient of the covariant Poisson algebra modulo the ideal generated by ker $F_{M}$ is canonically isomorphic to the canonical Poisson algebra. The same holds true for the corresponding Weyl algebras.

The Weyl algebra for the covariant Poisson algebra yields a local net of Weyl algebras satisfying time-slice and the usual locality axioms as required in AQFT.

## Outlook

- Representation theory needs to be developed via usual techniques (GNS)
- Time evolution for quadratic elements more tricky: *-exponential will not be in the completion (does not exists in a meaningful way).
- Still: it might give a one-parameter group of outer automorphisms.
- Go beyond the constant case: in finite dimensions possible for several examples via phase space reduction.
- Go beyond continuous case: separately continuous or bornological setting?

Thank you very much and have a good trip back!

