

The Geodesic Equation

paul.romatschke@colorado.edu

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From Coordinate Transformations to Gravity

- Last lecture, we introduced **general covariance**, the notion that physics does not depend on (arbitrarily chosen) coordinates
- By itself, this is just an exercise in math
- In this lecture, we'll take the first step towards considering **physical consequences** of taking the math seriously
- The procedure is akin to our treatment of **special relativity**
- In SR, we considered linear coordinate transformations (math); this led to physics consequences (e.g. time-delay)

Particle in Minkowski spacetime

Consider a massive particle in Minkowski spacetime:

- No forces
- Particle moves on a “straight line”
- Proper time for particle:

$$d\tau'^2 = -ds^2 = -g'_{\mu\nu} dx'^{\mu} dx'^{\nu} \quad (10.1)$$

- Absence of force means particle moves with constant velocity:

$$\frac{dx'^{\mu}}{d\tau'} = \text{const.} \quad (10.2)$$

- Take another derivative wrt to proper time:

$$\frac{d^2 x'^{\mu}}{d\tau'^2} = 0. \quad (10.3)$$

Einstein and the Elevator

Einstein:

- Enclosed in an elevator, you can't tell if it's accelerating or if some alien has messed with gravity
- Locally, gravity is **indistinguishable** from acceleration
- But there is a non-linear coordinate transformation from Minkowski space x' to accelerated coordinate system x :

$$x'^{\mu} = x'(x), \quad \text{where } x'^{\mu} = \begin{pmatrix} t \\ \vec{x} \end{pmatrix}. \quad (10.4)$$

- In Minkowski space

$$\frac{d^2 x'^{\mu}}{d\tau'^2} = 0, \quad (10.5)$$

so this equation must hold true

- Let's transform (10.5) to coordinates x^{μ} :

General Coordinate Transformations

- We perform (non-linear) transform from Minkowski coordinate x'^{μ} to new coordinates x^{μ}
- The transformation matrix is

$$R^{\mu}_{\nu} \equiv \frac{\partial x'^{\mu}}{\partial x^{\nu}} . \quad (10.6)$$

- The new metric $g_{\mu\nu}$ will (in general) be non-Minkowski
- But we can use that scalars are invariant under general coordinate transformations, e.g.

$$d\tau'^2 = -ds'^2 = -ds^2 = d\tau^2 , \quad (10.7)$$

hence $d\tau' = d\tau$.

General Coordinate Transformations

We thus have

$$\begin{aligned} 0 &= \frac{d^2 x'^{\mu}}{d\tau'^2} = \frac{d^2 x'^{\mu}}{d\tau^2} = \frac{d}{d\tau} \left(\frac{\partial x'^{\mu}}{\partial x^{\nu}} \frac{dx^{\nu}}{d\tau} \right), \\ &= \frac{\partial x'^{\mu}}{\partial x^{\nu}} \frac{d^2 x^{\nu}}{d\tau^2} + \frac{dx^{\nu}}{d\tau} \frac{d}{d\tau} \left(\frac{\partial x'^{\mu}}{\partial x^{\nu}} \right), \end{aligned}$$

General Coordinate Transformations

We thus have

$$\begin{aligned} 0 &= \frac{d^2 x'^{\mu}}{d\tau'^2} = \frac{d^2 x'^{\mu}}{d\tau^2} = \frac{d}{d\tau} \left(\frac{\partial x'^{\mu}}{\partial x^{\nu}} \frac{dx^{\nu}}{d\tau} \right), \\ &= \frac{\partial x'^{\mu}}{\partial x^{\nu}} \frac{d^2 x^{\nu}}{d\tau^2} + \frac{dx^{\alpha}}{d\tau} \frac{d}{d\tau} \left(\frac{\partial x'^{\mu}}{\partial x^{\alpha}} \right), \\ &= \frac{\partial x'^{\mu}}{\partial x^{\nu}} \frac{d^2 x^{\nu}}{d\tau^2} + \frac{dx^{\alpha}}{d\tau} \frac{dx^{\beta}}{d\tau} \frac{\partial^2 x'^{\mu}}{\partial x^{\alpha} \partial x^{\beta}}, \\ &= \frac{\partial x'^{\mu}}{\partial x^{\nu}} \left[\frac{d^2 x^{\nu}}{d\tau^2} + \frac{\partial x^{\nu}}{\partial x'^{\mu}} \frac{dx^{\alpha}}{d\tau} \frac{dx^{\beta}}{d\tau} \frac{\partial^2 x'^{\mu}}{\partial x^{\alpha} \partial x^{\beta}} \right], \\ &= R^{\mu}_{\nu} \left[\frac{d^2 x^{\nu}}{d\tau^2} + \frac{\partial x^{\nu}}{\partial x'^{\mu}} \frac{dx^{\alpha}}{d\tau} \frac{dx^{\beta}}{d\tau} \frac{\partial^2 x'^{\mu}}{\partial x^{\alpha} \partial x^{\beta}} \right], \end{aligned} \tag{10.8}$$

The Geodesic Equation

- The transformation matrix is non-singular, so (10.5) in coordinates x^μ becomes

$$\frac{d^2 x^\nu}{d\tau^2} + \frac{\partial x^\nu}{\partial x'^\mu} \frac{\partial^2 x'^\mu}{\partial x^\alpha \partial x^\beta} \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau} = 0. \quad (10.9)$$

- Looks complicated, let's introduce the new notation

$$\frac{\partial x^\nu}{\partial x'^\mu} \frac{\partial^2 x'^\mu}{\partial x^\alpha \partial x^\beta} \equiv \Gamma_{\alpha\beta}^\nu. \quad (10.10)$$

- With these new symbols, we have

$$\frac{d^2 x^\nu}{d\tau^2} + \Gamma_{\alpha\beta}^\nu \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau} = 0 \quad (10.11)$$

- This is the **geodesic equation**

Vectors and such

- Recall my warning about things that look like but are not vectors in lecture 4
- Take $\frac{d^2 x'^{\mu}}{d\tau'^2}$ as an example
- It carries one index, so one might be tempted to assume it is a vector under general coordinate transformations
- If it **was a vector**, then it would have to transform as

$$\frac{d^2 x'^{\mu}}{d\tau'^2} \rightarrow R^{\mu}_{\nu} \frac{d^2 x^{\nu}}{d\tau^2} \quad (10.12)$$

- However, we found by direct calculation in (10.11) that there is an additional piece proportional to $\Gamma^{\nu}_{\alpha\beta}$
- **This implies that $\frac{d^2 x^{\mu}}{d\tau^2}$ is not a vector under diffeomorphisms!**