Scale Factor Evolution

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- In lecture 35, we found the Friedmann equations for cosmology
- Using the shorthand notation $H(t) \equiv \frac{\dot{a}(t)}{a(t)}$ these are

$$\partial_t \epsilon + 3(\epsilon + p)H = 0, \quad 3\left(H^2 + \frac{\kappa}{a^2}\right) = 8\pi G\epsilon.$$
 (37.1)

In lecture 36, we identified a(t) with the scale factor of the universe
Let's get solutions to (37.1) for a(t)!

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Dust

- Eqns. (37.1) are two equations for three unknowns (ϵ, p, a)
- In order to close the system, we need another equation
- For instance, we can use the equation of state (EoS) $p = p(\epsilon)$
- EoS depends on the matter content we describe
- Simplest case: "dust" (pressure-less):

$$p = 0.$$
 (37.2)

• For dust, the matter part of (37.1) becomes

$$\partial_t \epsilon = -3\epsilon H \tag{37.3}$$

Other EoS

- We can also consider forms of matter other than dust
- A very simplified parametrization is

$$p = w\epsilon, \quad w = \text{const.}$$
 (37.4)

- For dust w = 0, while for pure radiation $w = \frac{1}{3}$
- Can also describe pure cosmological constant with w = -1
- In general, however, w is not a constant for real matter!
- In terms of (37.4), the matter part of (37.1) becomes

$$\partial_t \epsilon = -3H\epsilon(1+w).$$
 (37.5)

- For the simplified class of EoS (37.4), we can solve the matter part of the Friedmann equations
- We find

$$\epsilon(t) = \epsilon(t_0) \left(\frac{a(t_0)}{a(t)}\right)^{3(1+w)} = \epsilon_0 \left(\frac{a_0}{a}\right)^{3(1+w)}.$$
 (37.6)

Critical Energy Density

• Plugging this into (37.1), the geometry equation reads

$$H^{2} = \frac{8\pi G}{3} \epsilon_{0} \left(\frac{a_{0}}{a}\right)^{3(1+w)} - \frac{\kappa}{a^{2}}.$$
(37.7)

• Dividing by H^2 we find

$$1 = \frac{8\pi G\epsilon(t)}{3H^2} - \frac{\kappa}{a^2 H^2}.$$
 (37.8)

• Define the critical energy density

$$\epsilon_c \equiv \frac{3H^2}{8\pi G} \tag{37.9}$$

• In terms of ϵ_c , we have

$$1 = \frac{\epsilon}{\epsilon_c} - \frac{\kappa}{a^2 H^2} \tag{37.10}$$

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Density Parameter

• In a proliferation of symbols, let's furthermore define the **density parameter**

$$\Omega \equiv \frac{\epsilon(t)}{\epsilon_c} \tag{37.11}$$

which is a time-dependent quantity

• In terms of Ω , we have

$$\frac{\kappa}{a^2 H^2} = \Omega - 1. \qquad (37.12)$$

Now a²H² > 0 so that the sign on both sides of this equation has to match

Density and Geometry

• Consider the implications of

$$\frac{\kappa}{a^2 H^2} = \Omega - 1. \qquad (37.13)$$

- For a flat universe (Euclidean geometry) we have $\kappa=0$
- For $\kappa = 0$, we must have

$$\Omega = 1 \tag{37.14}$$

or $\epsilon = \epsilon_c$

Using our solution for the energy density, this gives

$$\epsilon_0 \left(\frac{a_0}{a}\right)^{3(1+w)} = \frac{3H^2}{8\pi G} \tag{37.15}$$

• Since $H = \frac{\dot{a}}{a}$ we can solve this equation for a(t)

Flat Universe Scale Factor

• Let's solve this equation using a power-law ansatz:

$$a(t) = a_0 \left(\frac{t}{t_0}\right)^{\alpha}, \quad \alpha = \text{const}$$
 (37.16)

- We have $H = \frac{\dot{a}}{a} = \frac{\alpha}{t}$ (37.17)
- Plugging this into (37.15) we have

$$\epsilon_0 \left(\frac{t}{t_0}\right)^{-3(1+w)\alpha} = \frac{3}{8\pi G} \frac{\alpha^2}{t^2}$$
(37.18)

• The power of t on both sides has to match:

$$\alpha = \frac{2}{3(1+w)} \tag{37.19}$$

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• For the flat universe with $\kappa = 0$, we therefore find

$$a(t) \propto t^{rac{2}{3(1+w)}}$$
. (37.20)

- For w > -1, this means the scale factor is always growing
- We have an expanding flat universe

The Flat Universe

- The evolution of the scale factor depends on the matter content
- For dust, where w = 0, we have

$$a(t) \propto t^{\frac{2}{3}} \tag{37.21}$$

• By contrast, for radiation, we have $w = \frac{1}{3}$, so we have

$$a(t) \propto t^{rac{1}{2}}$$
 (37.22)

• Of course, in a realistic setting we have both, and *w* changes depending on which form of matter dominates

Age of the Universe

- Let's look at the age of the universe again
- The beginning of the universe is fixed at t = 0 where a(0) = 0
- Today's time is t₀, and today's Hubble constant for the flat universe is

$$H_0 = \left. \frac{\dot{a}}{a} \right|_{t=t_0} = \frac{\alpha}{t_0} \tag{37.23}$$

• This implies for the age of the universe

$$t_0 = \frac{\alpha}{H_0} \tag{37.24}$$

• Plugging in
$$H_0 \simeq 70 rac{km/s}{Mpc}$$
 this gives

$$t_0 \simeq 14 \alpha \, Gyr \tag{37.25}$$

Age Problem Again

• For dust where $\alpha = \frac{2}{3}$ this gives

$$t_0 \simeq 9.3 Gyr \tag{37.26}$$

• For radiation, where $\alpha = \frac{1}{2}$, it's

$$t_0 \simeq 7 \, Gyr \tag{37.27}$$

 Since for a flat universe w > 0, all of these are too short for observed stars age

$$t_{\rm star} \simeq 15.2 \pm 3.7 \, Gyr$$
 (37.28)

• The flat universe with $\kappa = 0$ cannot be our universe!