

Quantum Mechanics as a Path Integral (III)

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References

- Laine and Vuorinen, “Basics of Thermal Field Theory”, Chapter 1:
▶ <https://arxiv.org/pdf/1701.01554.pdf>

Partition Function

Recall from Lecture 4:

- Quantum Mechanical Partition Function after integrating momenta:

$$Z(T) = \lim_{N \rightarrow \infty} \int \left[\prod_{i=1}^N \frac{dx_i}{\sqrt{2\pi\epsilon/m}} \right] e^{-\epsilon \sum_{j=0}^N \left[\frac{m}{2} \left(\frac{x_{j+1} - x_j}{\epsilon} \right)^2 + V(x_j) \right]}, \quad (5.1)$$

together with $x_{N+1} = x_1$.

Partition Function

- We can again write $Z(T)$ in continuum form

$$Z = C \int_{x(0)=x(\beta)} \mathcal{D}x \exp \left[- \int_0^\beta d\tau \left(\frac{m}{2} \left(\frac{dx(\tau)}{d\tau} \right)^2 + V(x(\tau)) \right) \right] \quad (5.2)$$

where C is a constant factor that for fixed N is given by

$$C = \left(\frac{m}{2\pi\epsilon} \right)^{N/2}. \quad (5.3)$$

Partition Function

- The factor C is divergent for $N \rightarrow \infty$
- However, C is independent from the potential $V(x)$
- Therefore, even though C is divergent, it contains no dynamical information. We (physicists) will simply not worry about it.

Partition Function

- So we have

$$Z = C \int_{x(0)=x(\beta)} \mathcal{D}x e^{-S_E}, \quad (5.4)$$

- We call

$$S_E = \int_0^\beta d\tau L_E, \quad (5.5)$$

the “Euclidean” **action** and

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$$L_E = \frac{m}{2} \left(\frac{dx(\tau)}{d\tau} \right)^2 + V(x(\tau)), \quad (5.6)$$

the “Euclidean” **Lagrangian**

Partition Function

- For the moment being, these labels (“action”, “Lagrangian”) are rather mysterious
- Let’s try to give some meaning to them by transforming our “imaginary” time variable to “real time” t :

$$\tau = it \tag{5.7}$$

We find

$$L_E \rightarrow -\frac{m}{2} \left(\frac{dx(t)}{dt} \right)^2 + V(x), \tag{5.8}$$

which is nothing else but potential minus kinetic energy of a classical particle

Partition Function

- Let's turn the argument around
- Starting with the classical Lagrangian of a point particle in potential $V(x)$

$$L = \frac{m}{2} \left(\frac{dx(t)}{dt} \right)^2 - V(x) \quad (5.9)$$

we can do a transformation to imaginary time

$$t \rightarrow -i\tau \quad (5.10)$$

and get the Euclidean Lagrangian from

$$L_E = -L(t = -i\tau). \quad (5.11)$$

Partition Function

- There is more:
- The action for a classical point particle is defined as

$$S = \int dt L, \quad (5.12)$$

where L is the classical Lagrangian from before

- If we also transform the action to imaginary time, we have

$$e^{iS} \rightarrow e^{-\int d\tau L_E} = e^{-S_E}, \quad (5.13)$$

with S_E the Euclidean action (5.5)

- Note: we have to put the integration limits $\tau \in [0, \beta]$ in S_E by hand

Partition Function

- To summarize, starting with the classical Lagrangian of a point particle (5.9) we get the *quantum mechanical* partition function $Z(T)$ by writing

$$Z = C \int_{x(0)=x(\beta)} \mathcal{D}x e^{-S_E}, \quad (5.14)$$

where we get e^{-S_E} from

$$\lim_{t \rightarrow -i\tau} e^{iS} = e^{-S_E}, \quad (5.15)$$

and where C is a (divergent) constant