

# Fermions II

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# Review

- In lecture 30 we calculated the fermionic path integral in QM

$$Z_F = \int \mathcal{D}c^* \mathcal{D}c e^{-\int_0^\beta d\tau [c^* \frac{dc}{d\tau} + H(c^*, c)]}. \quad (32.1)$$

- Here  $\hat{H}(\hat{a}^\dagger, \hat{a})$  was the Hamiltonian built out of anti-commuting ladder operators, e.g.

$$\hat{H} = \omega \left( \hat{a}^\dagger \hat{a} - \frac{1}{2} \right), \quad \{ \hat{a}, \hat{a}^\dagger \} = 1. \quad (32.2)$$

- In lecture 31, we considered the Dirac field equations

$$(i\cancel{\partial} - m) \psi = 0. \quad (32.3)$$

- In this lecture, we construct the Hamiltonian for the Dirac field, defining the path integral for half-integer spin fields

# Lagrangian for Dirac Field

- In lecture 22, we constructed a Lagrangian density  $\mathcal{L}$  for a complex scalar field  $\phi$  based on symmetry (e.g. Lorentz transformations)
- In lecture 23, we found the equations of motion for  $\phi, \phi^*$ , which were treated as independent components
- For the Dirac field, we *have* the equations of motion: the Dirac equation (32.3)
- Treating  $\psi$  and its Dirac adjoint  $\bar{\psi}$  as independent components, the equations of motion

$$\frac{\partial \mathcal{L}}{\partial \bar{\psi}} - \partial_\mu \frac{\partial \mathcal{L}}{\partial \partial_\mu \bar{\psi}} = 0, \quad (32.4)$$

therefore must correspond to (32.3)

## Lagrangian for Dirac Field

- The Lagrangian density that gives rise to the Dirac equation therefore is given by

$$\mathcal{L} = \bar{\psi} (i\not{\partial} - m) \psi \quad (32.5)$$

- From the Lagrangian density, we can construct the Hamiltonian density  $\mathcal{H}$  as

$$\mathcal{H} = \pi \partial_0 \psi - \mathcal{L}. \quad (32.6)$$

- Here  $\pi$  is the conjugate momentum to the field  $\psi$ ,

$$\pi = \frac{\partial \mathcal{L}}{\partial \partial_0 \psi} = \bar{\psi} i \gamma^0 = i \psi^\dagger. \quad (32.7)$$

- We therefore find the Hamiltonian density given by

$$\mathcal{H} = \bar{\psi} (-i\gamma^i \partial_i + m) \psi \quad (32.8)$$

## Hamiltonian for the Dirac Field

- From the Hamiltonian density, the Hamiltonian is given by

$$H = \int_x \mathcal{H} = \int_x \psi^\dagger (-i\gamma^0 \gamma^i \partial_i + \gamma^0 m) \psi \quad (32.9)$$

- Here  $\psi, \psi^\dagger$  are classical (but anti-commuting) fields
- If we were to *quantize* the theory, the operators corresponding to  $\psi, \psi^\dagger$  would play the same role as the fermionic ladder operators  $\hat{a}, \hat{a}^\dagger$
- Also, the form of the Hamiltonian (32.9) corresponds to (32.2), except for the constant term
- We therefore identify the classical (but anti-commuting) fields  $\psi, \psi^\dagger$  with the classical (but anti-commuting) variables  $c, c^*$  in (32.1)

# Path Integral for Fermions

- Using this identification, we can write down a path integral for the quantum field theory of a Dirac fermion
- We have

$$\mathcal{Z}_F = \int \mathcal{D}\psi^\dagger \mathcal{D}\psi e^{-S_E}, \quad (32.10)$$

where

$$S_E \equiv \int_0^\beta d\tau \int_{\mathbf{x}} \psi^\dagger [\partial_\tau - i\gamma^0 \gamma^i \partial_i + \gamma^0 m] \psi. \quad (32.11)$$

- As in lecture 5, we refer to  $S_E$  as the “Euclidean” action
- Note that the fermionic fields are *anti-periodic* in Euclidean time, e.g.  
 $\psi(\beta, \mathbf{x}) = -\psi(0, \mathbf{x})$

# Path Integral for Fermions

- We can define “Euclidean” Dirac matrices

$$\gamma_0^E = \gamma_0, \quad \gamma_i^E = -i\gamma^i, \quad (32.12)$$

which obey

$$\{\gamma_a^E, \gamma_b^E\} = 2\delta_{ab}\mathbf{1}, \quad \gamma_a^{E\dagger} = \gamma_a^E \quad (32.13)$$

- We can then simplify the Euclidean Lagrangian density as

$$\mathcal{L}_E = \bar{\psi} \left[ \gamma_a^E \partial_a + m \right] \psi \quad (32.14)$$

## Path Integral for Fermions

- Finally, since  $\psi^\dagger$  and  $\bar{\psi}$  only differ by a (matrix)-constant, the Jacobian for changing  $\mathcal{D}\psi^\dagger$  to  $\mathcal{D}\bar{\psi}$  is just a constant
- Ignoring the constant Jacobian, we have

$$Z_F = \int \mathcal{D}\bar{\psi} \mathcal{D}\psi e^{-\int_x \bar{\psi} [\gamma_a^E \partial_a + m] \psi}, \quad (32.15)$$

as the path integral for a quantum field theory for a free Dirac fermion

- We will evaluate the partition function in the next lecture