

# SYMMETRIC SPACES

How can we use some supposed symmetry of a metric space to gain information about the metric?

↳ Forminvariance of the metric :  $g'_{\mu\nu}(x') = g_{\mu\nu}(x) \quad \forall x'$  (1)

transformed metric :  $g'_{\mu\nu}(x') = \frac{\partial x^\rho}{\partial x'^\mu} \frac{\partial x^\sigma}{\partial x'^\nu} g_{\rho\sigma}(x) \stackrel{!}{=} g_{\mu\nu}(x')$

$$\Downarrow$$

$$g_{\mu\nu}(x) = \frac{\partial x'^\rho}{\partial x^\mu} \frac{\partial x'^\sigma}{\partial x^\nu} g_{\rho\sigma}(x') \quad (2)$$

Any transformation  $x^\mu \rightarrow x'^\mu$ , that satisfies (2) is called an isometry  
infinitesimal coordinate transformations  $x'^\mu = x^\mu + \epsilon \xi^\mu(x) \quad |\epsilon| \ll 1$  (3)

$$\begin{aligned} \mathcal{O}(\epsilon) \quad g_{\mu\nu}(x) &= (\delta_\mu^\rho + \epsilon \xi^{\rho, \mu}) (\delta_\nu^\sigma + \epsilon \xi^{\sigma, \nu}) (g_{\rho\sigma}(x) + \epsilon \xi^\lambda g_{\rho\sigma, \lambda} + \dots) \\ &= g_{\mu\nu}(x) + \epsilon g_{\rho\nu} \xi^{\rho, \mu}(x) + \epsilon g_{\mu\sigma} \xi^{\sigma, \nu} + \epsilon \xi^\lambda g_{\mu\nu, \lambda} \end{aligned}$$

$$\Downarrow$$

$$0 = \xi^{\rho, \mu} g_{\rho\nu} + \xi^{\sigma, \nu} g_{\mu\sigma} + \xi^\lambda g_{\mu\nu, \lambda} \quad (4)$$

this can be rewritten in terms of  $\xi_\mu = g_{\mu\nu} \xi^\nu$

$$0 = \frac{\partial}{\partial x^\mu} (\underbrace{\xi^\rho g_{\rho\nu}}_{\xi_\nu}) + \frac{\partial}{\partial x^\nu} (\underbrace{g_{\mu\rho} \xi^\rho}_{\xi_\mu}) - \underbrace{\xi^\rho g_{\rho\nu, \mu} - \xi^\rho g_{\mu\rho, \nu}}_{+ \xi^\lambda (g_{\mu\nu, \lambda} - g_{\lambda\nu, \mu} - g_{\lambda\mu, \nu})}$$

$$(-2) \Gamma_{\lambda\mu\nu} = (-2) \Gamma_{\lambda\nu\mu}$$

$$\Rightarrow 0 = \xi_{\mu; \nu} - \xi^\lambda \Gamma_{\lambda\mu\nu} + \xi_{\nu; \mu} - \xi^\lambda \Gamma_{\lambda\nu\mu} = \xi_{\mu; \nu} + \xi_{\nu; \mu}$$

$$\underline{0 = \xi_{\mu; \nu} + \xi_{\nu; \mu}} \quad (5)$$

Any vector, that satisfies (4) or (5) is said to form a Killing vector of the metric  $g_{\mu\nu}(x)$

↳ Killing condition (5) much more restrictive than it looks :  
 allows to determine the whole function  $\xi_\mu(x)$  from any given values  $\xi_\mu$  and  $\xi_{\mu; \nu}$  at some point  $X$

commutator of 2 covariant derivatives :  $\nabla_{\sigma} \nabla_{\rho} \xi_{\mu} - \nabla_{\rho} \nabla_{\sigma} \xi_{\mu} = -R^{\lambda}{}_{\sigma\rho\mu} \xi_{\lambda}$  (6)

cyclic sum rule for curvature tensor :  $R^{\lambda}{}_{\sigma\rho\mu} + R^{\lambda}{}_{\mu\sigma\rho} + R^{\lambda}{}_{\rho\mu\sigma} = 0$  (7)  $\big|_{\xi_{\lambda}}$

$$\Rightarrow \nabla_{\sigma} \nabla_{\rho} \xi_{\mu} - \nabla_{\rho} \nabla_{\sigma} \xi_{\mu} + \nabla_{\mu} \nabla_{\rho} \xi_{\sigma} - \nabla_{\rho} \nabla_{\mu} \xi_{\sigma} + \nabla_{\rho} \nabla_{\sigma} \xi_{\mu} - \nabla_{\sigma} \nabla_{\rho} \xi_{\mu} = 0$$

$\downarrow$  ( $\nabla_{\rho} \xi_{\sigma} + \nabla_{\sigma} \xi_{\rho} = 0$ !)  
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$$2 [ \nabla_{\sigma} \nabla_{\rho} \xi_{\mu} - \nabla_{\rho} \nabla_{\sigma} \xi_{\mu} - \nabla_{\sigma} \nabla_{\mu} \xi_{\rho} ] = 0$$

$$\downarrow$$

$$-R^{\lambda}{}_{\sigma\rho\mu} \xi_{\lambda}$$

$$\downarrow$$

$$\nabla_{\mu} \xi_{\rho} = -R^{\lambda}{}_{\sigma\rho\mu} \xi_{\lambda} \quad (8)$$

If  $\xi_{\mu}$  and  $\xi_{\mu;\nu}$  are known at a typical point  $X$ , all higher derivatives of  $\xi_{\mu}$  at  $X$  can be determined and hence the entire function  $\xi_{\mu}(x)$  in a neighbourhood of  $X$  by Taylor expansion:

$$\xi_{\lambda}(x) = A_{\lambda}{}^{\rho} (x, X, g) \xi_{\rho}(X) + B_{\lambda}{}^{\rho\sigma} (x, X, g) \xi_{\rho;\sigma}(X) \quad (9)$$

$$\Downarrow$$

$$B_{\lambda}{}^{\rho\sigma} = -B_{\lambda}{}^{\sigma\rho}$$

$\rightarrow$  In a spacetime of  $N$  dimensions there are up to  $N$  quantities  $\xi_{\lambda}(X)$  and up to  $\frac{1}{2}N(N-1)$  quantities  $\xi_{\mu;\nu}(X) = -\xi_{\nu;\mu}(X)$

$$\Downarrow$$

maximally  $N + \frac{1}{2}N(N-1) = \frac{1}{2}N(N+1)$  Killing vectors in  $N$  dimensions

homogeneity

A spacetime is said to be homogeneous, if there are infinitesimal isometries, that carry any given point  $X$  into any other point in its neighbourhood.

$\Rightarrow$  the metric must admit Killing vectors, that at any given point take on all possible values for  $\xi_{\mu}(X) \hat{=} N$  Killing-vectors in  $N$  dimensions

$$\xi_{\mu}^{(n)}(X) \quad n=1,2,\dots,N$$

isotropy

A space time is said to be isotropic at a given point  $X$ , if there are Killingvectors  $\xi_{\mu}$  for which  $\xi_{\mu}(X) = 0$  and  $\xi_{\mu;\nu}(X)$  take on all possible values  $\hat{=} \frac{1}{2}N(N-1)$

→ not all metrics admit maximum number of Killing vectors

Whether  $\nabla_{\mu} \nabla_{\nu} g_{\alpha\beta} = -R^{\sigma}{}_{\rho\gamma\mu} g_{\sigma\alpha}$  is soluble for a given set of initial data  $g_{\alpha}(X), g_{\alpha\beta}(X)$  depends on the integrability of this equation

→ commutator of covariant derivatives of tensors:

$$T_{\alpha\mu\nu;\kappa} - T_{\alpha\mu\kappa;\nu} = -T_{\sigma\mu} R^{\sigma}{}_{\alpha\nu\kappa} - T_{\alpha\sigma} R^{\sigma}{}_{\mu\nu\kappa}$$

$$\circ) \nabla_{\rho} \nabla_{\mu} \nabla_{\nu} \kappa - \nabla_{\rho} \nabla_{\nu} \nabla_{\mu} \kappa = -\nabla_{\lambda} g_{\mu}{}^{\lambda} R^{\lambda}{}_{\rho\sigma\nu} - \nabla_{\rho} g_{\lambda}{}^{\lambda} R^{\lambda}{}_{\mu\sigma\nu}$$

$$\circ) \nabla_{\rho} \nabla_{\mu} g_{\alpha\beta} = -R^{\lambda}{}_{\sigma\mu\rho} g_{\lambda\alpha}$$

$$\begin{aligned} \underline{\underline{[-R^{\lambda}{}_{\rho\sigma\nu} \partial_{\mu}^{\kappa} + R^{\lambda}{}_{\mu\sigma\nu} \partial_{\rho}^{\kappa} + R^{\lambda}{}_{\nu\sigma\mu} \partial_{\sigma}^{\kappa} - R^{\lambda}{}_{\sigma\mu\nu} \partial_{\nu}^{\kappa}] g_{\lambda\kappa} =}} \\ = \underline{\underline{(R^{\lambda}{}_{\sigma\mu\nu} - R^{\lambda}{}_{\nu\sigma\mu}) g_{\lambda\alpha}} \quad (10)} \end{aligned}$$

maximally symmetric space : can find Killing vectors, for which  $\nabla_{\mu} g_{\alpha\beta} = 0$  &  $g_{\mu\nu}$  arbitrary antisymm. matrix

⇓

coefficient in (10) of  $g_{\lambda\kappa}$  must have a vanishing antisymmetric part

$$[\ ]_{\lambda\kappa} - [\ ]_{\kappa\lambda} = 0$$

⇓

$$\begin{aligned} -R^{\lambda}{}_{\rho\nu\sigma} \partial_{\mu}^{\kappa} + R^{\lambda}{}_{\mu\sigma\nu} \partial_{\rho}^{\kappa} + R^{\lambda}{}_{\nu\sigma\mu} \partial_{\sigma}^{\kappa} - R^{\lambda}{}_{\sigma\mu\nu} \partial_{\nu}^{\kappa} = \\ = -R^{\kappa}{}_{\rho\nu\sigma} \partial_{\mu}^{\lambda} + R^{\kappa}{}_{\mu\sigma\nu} \partial_{\rho}^{\lambda} + R^{\kappa}{}_{\nu\sigma\mu} \partial_{\sigma}^{\lambda} - R^{\kappa}{}_{\sigma\mu\nu} \partial_{\nu}^{\lambda} \quad | \cdot \partial_{\mu}^{\kappa} \end{aligned}$$

$$\begin{aligned} -N R^{\lambda}{}_{\rho\nu\sigma} - R^{\lambda}{}_{\rho\nu\sigma} + \underbrace{R^{\lambda}{}_{\nu\rho\sigma}}_{-R^{\lambda}{}_{\nu\sigma\rho}} - R^{\lambda}{}_{\sigma\rho\nu} = -R^{\lambda}{}_{\rho\nu\sigma} + 0 + \underbrace{R^{\mu}{}_{\nu\rho\mu} \partial_{\sigma}^{\lambda} - R^{\mu}{}_{\sigma\rho\mu} \partial_{\nu}^{\lambda}}_{-R_{\nu\rho} \partial_{\sigma}^{\lambda} + R_{\sigma\rho} \partial_{\nu}^{\lambda}} \end{aligned}$$

(with  $R^{\mu}{}_{\mu\sigma\nu} = 0$        $R_{\sigma\rho} = R^{\mu}{}_{\sigma\mu\rho} = -R^{\mu}{}_{\rho\mu\sigma}$ )

$$(N-1) R^{\lambda}{}_{\rho\sigma\nu} = R_{\nu\rho} \partial_{\sigma}^{\lambda} - R_{\sigma\rho} \partial_{\nu}^{\lambda} \quad | \cdot g_{\alpha\lambda}$$

$$(N-1) R_{\alpha\rho\sigma\nu} = g_{\alpha\sigma} R_{\nu\rho} - g_{\alpha\nu} R_{\sigma\rho} \quad (13)$$

must be antisymmetric in  $(\alpha\rho)$

$$(N-1) R_{\alpha\rho\sigma\nu} = -g_{\rho\sigma} R_{\nu\alpha} + g_{\rho\nu} R_{\sigma\alpha} \quad | \cdot g^{\alpha\nu}$$

$$(N-1) (-R_{\rho\sigma}) = -g_{\rho\sigma} R^{\lambda}{}_{\lambda} + R_{\sigma\rho}$$

$$\Rightarrow \underline{\underline{R_{\sigma\rho} = \frac{1}{N} g_{\sigma\rho} R^{\lambda}{}_{\lambda}}} \quad (14)$$

Inserting (14) into (13)

$$(N-1) R_{\lambda\rho\sigma\tau} = \frac{1}{N} R^{\alpha}{}_{\alpha} (g_{\lambda\sigma} g_{\rho\tau} - g_{\lambda\rho} g_{\sigma\tau})$$

$$\underline{R_{\lambda\rho\sigma\tau} = \frac{R^{\alpha}{}_{\alpha}}{N(N-1)} (g_{\lambda\sigma} g_{\rho\tau} - g_{\lambda\rho} g_{\sigma\tau})} \quad (15)$$

Use Bianchi identities to say something about the position dependence of  $R^{\alpha}{}_{\alpha}$

$$\text{BI} \quad R_{\lambda\mu\nu\kappa;\tau} + R_{\lambda\mu\tau\nu\kappa} + R_{\lambda\mu\kappa\tau\nu} = 0 \quad | \cdot g^{\lambda\nu}$$

$$R_{\mu\kappa;\tau} - R_{\mu\tau;\kappa} + R^{\nu}{}_{\mu\kappa\tau;\nu} = 0 \quad | \cdot g^{\mu\kappa}$$

$$R_{;\tau} - R^{\mu}{}_{\tau;\mu} - R^{\nu}{}_{\tau;\nu} = 0$$

$$(R^{\mu}{}_{;\tau} - \frac{1}{2} \delta^{\mu}_{\tau} R)_{;\mu} = 0 \quad \Rightarrow \quad \underline{(R^{\mu\nu} - \frac{1}{2} g^{\mu\nu} R)_{;\mu} = 0}$$

$$\Leftrightarrow \text{with } R^{\mu\nu} = \frac{1}{N} g^{\mu\nu} R$$

$$\left( \frac{1}{N} \delta^{\mu}_{\tau} R - \frac{1}{2} \delta^{\mu}_{\tau} R \right)_{;\mu} = \left( \frac{1}{N} - \frac{1}{2} \right) R_{;\tau} = \underline{\left( \frac{1}{N} - \frac{1}{2} \right) R_{;\tau} = 0} \quad R \text{ scalar!}$$

$$N \neq 2 \quad \underline{R_{;\tau} = 0}$$

$\Rightarrow$  maximally symmetric spaces have a constant curvature

Maximally symmetric spaces - construction

maximally symm. spaces  $\rightarrow$  essentially unique (specified by a curvature constant  $K$  & number of positive or negative eigenvalues of the metric)

$\Rightarrow$  consider a flat  $(N+1)$  dim space with metric given by

$$\underline{ds^2 = g_{AB} dx^A dx^B = C_{\mu\nu} dx^\mu dx^\nu + \frac{1}{K} dz^2} \quad (1) \quad \begin{matrix} A=1, \dots, N+1 \\ \mu, \nu=1, \dots, N \end{matrix}$$

where  $C_{\mu\nu}$  is a constant  $(N \times N)$  matrix symmetric  $K$  some constant

$\Rightarrow$  Embed a non-Euclidean  $N$ -dim space in a larger space by restricting the variables  $x^\mu$  &  $z$  to the surface of a (pseudo)sphere:

$$\underline{K C_{\mu\nu} x^\mu x^\nu + z^2 = 1} \quad (2) \Rightarrow 2z dz + 2K C_{\mu\nu} x^\mu dx^\nu = 0$$

$$\Rightarrow (dz)^2 = \dots$$

$$\underline{ds^2 = C_{\mu\nu} dx^\mu dx^\nu + \frac{K (C_{\mu\nu} dx^\mu dx^\nu)^2}{1 - K C_{\mu\nu} x^\mu x^\nu}}$$

metric 
$$\underline{g_{\mu\nu}(x) = C_{\mu\nu} + \frac{K}{1 - K C_{\mu\nu} x^\mu x^\nu} (C_{\mu\alpha} x^\alpha C_{\nu\beta} x^\beta)}$$

(1) and (2) invariant under  $x'^A = R^A_B x^B$   $g_{AB} R^A_C R^B_D = g_{CD}$   
 $\frac{1}{2} N(N+1)$  parameter  
 o)  $x'^\mu = R^\mu_\nu x^\nu$  with  $\sum_{\mu, \nu} R^\mu_\alpha R^\nu_\beta = C_{\alpha\beta}$   $\frac{1}{2} N(N-1)$  parameter (isotropy)

o)  $x'^\mu = R^\mu_A x^A = x^\mu + a^\mu \cdot [\sim]$  (homogeneity: origin is taken into any other point)  
 $N$  parameter

$\Rightarrow \frac{1}{2} N(N+1)$  parameter  $\Rightarrow$  max symmetric

Spaces with maximally symmetric subspaces

suitable coordinate system  $(N\text{-dim})$   $v^a$   $(N-M)$  dim  $u^i$   $M$ -dim max symm. subspace

$$\underline{ds^2 = g_{\mu\nu} dx^\mu dx^\nu = g_{ab}(v) dv^a dv^b + f(v) \tilde{g}_{ij}(u) du^i du^j}$$

$\tilde{g}_{ij} \hat{=}$  metric of max symm. subspace

Many cases of practical importance  $\rightarrow$  max. symmetric spaces are SPACES, as opposed to space-times, so all eigenvalues of the submatrix  $\tilde{g}_{ij}$  are positive!

Take for  $K \neq 0$   $C_{ij}$  as  $\frac{1}{|K|} \delta_{ij}$

$$\underline{d\bar{\ell}^2} = \left( \frac{1}{|K|} \delta_{ij} + \frac{K}{\left(1 - \frac{K}{|K|} \delta_{km} x^k x^m\right)} \cdot \frac{1}{|K|^2} x_i x_j \right) dx^i dx^j$$

$$\rightarrow \underline{K > 0} \quad \Downarrow \quad d\bar{\ell}^2 = \frac{1}{K} \left( d\bar{x}^2 + \frac{(\bar{x} \cdot d\bar{x})^2}{1 - \bar{x}^2} \right)$$

$$\underline{K < 0} \quad d\bar{\ell}^2 = \frac{1}{K} \left( d\bar{x}^2 - \frac{(\bar{x} \cdot d\bar{x})^2}{1 + \bar{x}^2} \right)$$

$$\underline{K = 0} \quad C_{ij} = \delta_{ij} \quad d\bar{\ell}^2 = d\bar{x}^2$$

$$\underline{ds^2} = \underline{g_{ab}(v) dv^a dv^b + f(v) \left\{ d\bar{u}^2 + \frac{k (\bar{u} d\bar{u})^2}{1 - k \bar{u}^2} \right\}}$$

$$\begin{aligned} k = +1 & \text{ for } K > 0 \\ k = -1 & \text{ for } K < 0 \\ k = 0 & \text{ for } K = 0 \end{aligned}$$

$\Rightarrow$  Spherically symmetric homogeneous space-time

$N=4$  & 3dim maximally symmetric subspace

$$ds^2 = g(v) dv^2 - f(v) \left\{ d\bar{u}^2 + \frac{k (\bar{u} d\bar{u})^2}{1 - k \bar{u}^2} \right\}$$

new coordinates  $d = \int \sqrt{g(v)} dv$

$$u^1 = r \sin \theta \cos \varphi$$

$$u^2 = r \sin \theta \sin \varphi$$

$$u^3 = r \cos \theta$$

$$\underline{\underline{ds^2 = dt^2 - R^2(t) \left\{ \frac{dx^2}{1 - kx^2} + x^2 d\theta^2 + x^2 \sin^2 \theta d\varphi^2 \right\}}}$$